Cosmic Axion Detection with an Amplifying B-field Ring Apparatus



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Two important facts to keep in mind in any dark matter talk (at least, today)

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Fact 1: we know a lot about dark matter

Fact 2: we know almost nothing about dark matter

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Fact 2: we know almost nothing about dark matter

No evidence for non-gravitational interactions



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No evidence for particular dark-matter mass

Over 20 orders of magnitude in DM mass!



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Dark-matter: BSM physics exists



- Clear evidence that dark-matter (BSM physics) exists
- Well motivated dark-matter models (WIMPs, axions, ...)

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Dark matter models

Name	What is it?	Motivation
Axion	$\left(\bar{\theta} + \frac{a}{f_a}\right) G_{\mu\nu} \tilde{G}^{\mu\nu}$	Strong CP
Neutralino (WIMP)	$ ilde{B}, ilde{W}_3, ilde{H}_u, ilde{H}_d$	Hierarchy Problem (why Higgs mass so light)

How can we probe axion dark matter?



- Astrophysics/cosmology: stellar cooling, CMB, BBN (Phys. Lett.
- B. 2014: K. Blum, R. D'Agnolo, M. Lisanti, B.S.), superradiance
- Laboratory experiments: ADMX (resonant cavity), CAST (axion helioscope)
- New proposal: *PRL* 117, Sept. 2016 (Y. Kahn, **B.S.**, J. Thaler): A broadband approach to axion dark matter detection

Outline

- Axion particle physics (review)
- Axion cosmology (review and work in progress)
- ABRACADABRA: Cosmic axion detection (theory)

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ABRACADABRA-10 cm at MIT (experiment)

Why axions and what are they?

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• Problem: CP-violating $\delta_{\text{CKM}} \sim O(1)$, but $|\bar{\theta}| < 10^{-10}$

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 - Axion can be dark matter

$$\mathcal{L}_{\text{QCD}}^{\mathcal{CP}} = -\frac{\theta g^2}{32\pi^2} G_{\mu\nu} \tilde{G}^{\mu\nu} - \sum_q \bar{q} \, m_q e^{-i\phi_q \gamma_5} q$$

►
$$U(1)_A$$
 anomaly: $q \to e^{-i\alpha_q \gamma_5} q$
 $\theta \to \theta + 2 \sum_q \alpha_q$

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- Calculation: $d_n \approx 2.4 \times 10^{-16} \,\overline{\theta} \, \mathrm{e} \cdot \mathrm{cm}$

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- Calculation: $d_n \approx 2.4 \times 10^{-16} \bar{\theta} \, \mathrm{e} \cdot \mathrm{cm}$
- Measurement: $|\bar{\theta}| < 10^{-10}$
- No anthropic argument for why $\bar{\theta}$ is so small!

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QCD generates axion mass:

$$\begin{split} V(a) &\approx \frac{1}{2} f_a^2 m_a^2 \left(\bar{\theta} + \frac{a}{f_a}\right)^2 \\ m_a &\approx \frac{f_\pi}{f_a} m_\pi \approx 10^{-9} \ \mathrm{eV}\left(\frac{10^{16} \ \mathrm{GeV}}{f_a}\right) \end{split}$$

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ABRACADABRA-10 cm at MIT (experiment)

Axion dark matter is a classical field



- ► de Broglie wavelength: $\lambda_{dB} = \frac{2\pi}{p} \approx \frac{2\pi}{mv}$
 - Axion ($m = 10^{-9}$ eV): $\lambda_{dB} \approx 8 \times 10^3$ km
 - WIMP (m = 100 GeV): $\lambda_{dB} \approx 8 \times 10^{-17} \text{ km}$

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 - WIMP (m = 100 GeV): $\lambda_{dB} \approx 8 \times 10^{-17} \text{ km}$
- ► Local DM energy density: $\rho_{DM} \approx 0.4 \text{ GeV/cm}^3$
- Local occupancy number: $\mathcal{N} \approx (\rho_{DM}/m) \times \lambda_{db}^3$
 - $\mathcal{N}_{\text{axion}} \approx 10^{44}$
 - $\mathcal{N}_{\text{WIMP}} \approx 10^{-36}$

The axion as dark matter $(f_a > H_1/2\pi)$ $\ddot{a} + 3H\dot{a} + m_a^2 a = 0$ $(H = T^2/m_{pl})$

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• After $3H = m_a$, coherent oscillations ~ NR matter

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• Today:
$$\Omega_a h^2 \sim 0.1 \left(\frac{f_a}{10^{12} \text{ GeV}} \right)^{7/6} \theta_i^2$$



• After $3H = m_a$, coherent oscillations $\sim NR$ matter

• Today:
$$\Omega_a h^2 \sim 0.1 \left(\frac{f_a}{10^{12} \text{ GeV}}\right)^{7/6} \theta_i^2$$

• $f_a = 10^{16} \text{ GeV} \rightarrow |\theta_i| \lesssim 10^{-3} - 10^{-2}$ (e.g., Tegmark, Aguirre, Rees, Wilczek '05)

Preliminary! In progress with Andrey Katz

 $\ddot{a} + (3H + \gamma_{\text{QCD}})\dot{a} + m_a^2 a = 0$

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QCD Damping rate (McLerran et. al. 1990) :

$$\begin{split} \gamma_{\text{QCD}} &= \frac{1}{f_a^2 T} \int d^4 x \langle \frac{\alpha_s}{4\pi} \text{tr}[G_{\mu\nu} \tilde{G}^{\mu\nu}(x)] \frac{\alpha_s}{4\pi} \text{tr}[G_{\mu\nu} \tilde{G}^{\mu\nu}(0)] \rangle_T \\ &= \frac{\Gamma_{\text{sphaleron}}}{f_a^2 T} \\ &\propto (\text{large coefficient}) \times \frac{T^3}{f_a^2} \end{split}$$

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► QCD Damping rate (McLerran et. al. 1990) :

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• Important if $\gamma_{\text{QCD}} \sim H$ at $T \sim 1$ GeV:

$$\frac{(1 \text{ GeV})^3}{f_a^2} \sim \frac{(1 \text{ GeV})^2}{10^{18} \text{ GeV}}$$

• Important for $f_a \lesssim 10^{10}$ GeV (with the $\mathcal{O}(1)$ numbers)

Probably not if $f_a \gtrsim 10^{11} \text{ GeV}$



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Likely yes if $f_a \lesssim 10^{10} \text{ GeV}$



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How can we probe axion dark matter?



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Axion dark matter modifies Maxwell's equations

Recall axions also couple to QED:

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Scott Thomas and Blas Cabrera (2010), Sikivie et. al. (2013)

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 Magnetoquasistatic approximation: new electric current that follows B-field lines

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 Magnetoquasistatic approximation: new electric current that follows B-field lines

$$\nabla \times \mathbf{B} = \frac{g_{a\gamma\gamma}}{\partial t} \mathbf{B} \frac{\partial a}{\partial t}$$

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• Locally:
$$a(t) \approx a_0 \sin(m_a t)$$
 and $\frac{1}{2}m_a^2 a_0^2 = \rho_{\text{DM}}$

$$\mathbf{F} \mathbf{J}_{\text{eff}} = g_{a\gamma\gamma} \sqrt{2 \rho_{\text{DM}}} \mathbf{B} \sin(m_a t)$$

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Axion dark matter generates magnetic flux



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 Estimate *B*-field induced through pickup loop (r = a = h = R)

Axion dark matter generates magnetic flux



- ► Estimate B-field induced through pickup loop (r = a = h = R)
- Axion effective current: $I_{\text{eff}} \sim R^2 J_{\text{eff}}$
- $\blacktriangleright \ B \sim \frac{I_{\rm eff}}{R} \sim R g_{a\gamma\gamma} \sqrt{2 \, \rho_{\rm DM}} \mathbf{B_0} \sin(m_a t)$
- ► $f_a = 10^{16} \text{ GeV}, \mathbf{B_0} \sim 5 \text{ T}, R \sim 4 \text{ m}: B \sim 10^{-22} \text{ T} \text{ (KSVZ)}$

Two readout strategies



Better at low frequency

Better at high frequency

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Two readout strategies



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Two readout strategies





Better at high frequency

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- Example from MRI application: (Myers et. al. 2007)
 - B-field sensitivity: $S_B^{1/2} \approx 6.4 \times 10^{-17} \text{ T}/\sqrt{\text{Hz}}$

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- ▶ R ≈ 3.3 cm
- Scale to $R \approx 4 \text{ m}$
 - $\blacktriangleright ~S_B^{1/2} \approx 5 \times 10^{-20} ~{\rm T}/\sqrt{\rm Hz}$



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 - Coherence time: $\tau \sim 2\pi/(m_a v^2) \sim 10$ s ($v \sim 10^{-3}$)

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 - S/N = 1 for $B = S_B^{1/2} (t\tau)^{-1/4} \sim 10^{-22} \text{ T}$

Axion dark matter projected reach



Axion dark matter projected reach



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Axion dark matter projected reach



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ABRACADABRA-10 cm at MIT (experiment)

The MIT prototype: ABRACADABRA-10 cm

- ABRACADABRA: A Broadband/Resonant Approach to Cosmic Axion Detection with an Amplifying B-field Ring Apparatus
- ▶ Dimensions: 12×12 cm² (R = 3 cm, h = 12 cm), B = 1 T
- People (LNS+CTP, PSFC, +1 Princeton): Janet Conrad, Joe Formaggio, Sarah Heine, Yoni Kahn, Joe Minervini, Jonathan Ouellet, Kerstin Perez, Alexey Radovinsky, B.S., Jesse Thaler, Daniel Winklehner, Lindley Winslow

- Lindley's dilution refrigerator (< 100 mK)</p>
 - Workable space: $R \sim 25 \text{ cm}, h \sim 25 \text{ cm}$

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- Lindley's dilution refrigerator (< 100 mK)</p>
 - Workable space: $R \sim 25 \text{ cm}, h \sim 25 \text{ cm}$
- Funded by the NSF (as of this week)

ABRACADABRA-10 cm



ABRACADABRA-10 cm



Thanks Daniel Winklehner for CAD model and slides.

ABRA-10 cm: vertical cut



ABRA-10 cm: pickup cylinder



ABRA-10 cm: reach after 1 month



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Complementary proposals for axion dark matter experiments

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CASPEr: oscillating neutron EDM



Light bosonic dark matter future

- MIT: ABRA-10 cm followed by ABRA-1 m ($B \sim 10$ T)
- ABRA-1 m: multiple experiments at different locations
 - Preliminary discussions with Korean Center for Axion and Precision Physics (Yannis Semertzidis)
- Axions and light bosonic dark matter well motivated by high-scale physics (e.g., compactified string theory)
- Detection may provide window to high-scale physics (GUT scale, inflation, ...)
- New ideas to search for ultra-light scalars, dark-photons, etc. (laboratory experiments + astrophysics)
 - e.g., CASPEr experiment
 - Black Hole superradiance

Questions?

Axion Backup Slides

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Magnetic field sensitivity calculation

- $B(t) = B_0 \sin[\omega_0 t + \phi(t)] + B_n(t)$
- $\phi(t)$: evolves over coherence time τ

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$$P(\omega) \equiv \frac{1}{\sqrt{T}} \int_0^T dt B(t) \sin(\omega t) = P_0(\omega) + P_n(\omega)$$

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- $T < \tau$:
 - $|P_0(\omega_0)|^2 \propto B^2 T \to B^2 = S_B^{1/2}(\omega_0)/T$

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 $\blacktriangleright \ T > \tau$

$$\blacktriangleright |P_0(\omega_0)|^2 \propto \frac{B^2}{T} \times T\tau = B^2 \tau$$

▶ But, line-width is broad and can resolve $N = T/\tau$ different frequencies

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- $\bullet \quad B(t) = B_0 \sin[\omega_0 t + \phi(t)] + B_n(t)$
- $\phi(t)$: evolves over coherence time τ

$$P(\omega) \equiv \frac{1}{\sqrt{T}} \int_0^T dt B(t) \sin(\omega t) = P_0(\omega) + P_n(\omega)$$

- ► Spectral density: $\lim_{T \to \infty} |P_n(\omega)|^2 \to S_B^{1/2}(\omega) [T / \sqrt{Hz}]$
- $T < \tau$:

•
$$|P_0(\omega_0)|^2 \propto B^2 T \to B^2 = S_B^{1/2}(\omega_0)/T$$

 $\blacktriangleright \ T > \tau$

$$\blacktriangleright |P_0(\omega_0)|^2 \propto \frac{B^2}{T} \times T\tau = B^2 \tau$$

▶ But, line-width is broad and can resolve $N = T/\tau$ different frequencies

• $B^2 = S_B^{1/2}(\omega_0)/\tau/\sqrt{N} = S_B^{1/2}(\omega_0)/\sqrt{T\tau}$

Broadband: detailed calculation

Cryogenic Current Comparator



Sese et. al., 1999

Axion DM: Broadband Readout



CASPEr: BBN and tuning bounds

$$\mathcal{L}_{\text{axion}} = -\left(\bar{\theta} + \frac{a}{f_a}\right) \frac{g^2}{32\pi^2} G_{\mu\nu} \tilde{G}^{\mu\nu}$$

- QCD generates minimum m_a
- Effective operator changes neutron-proton mass difference in early universe (Phys. Lett. B. 2014; K. Blum, R. D'Agnolo, M. Lisanti, B.S.)

