

Dark Sectors with a Mass Gap

Josh Ruderman (NYU) @Rutgers 5/17/2016

- Raffaele D'Agnolo, JTR, **[1505.07107](http://arxiv.org/abs/1505.07107)**
- Duccio Pappadopulo, JTR, Gabriele Trevisan, **[1602.04219](http://arxiv.org/abs/1602.04219)**

Towards the Neutrino Floor

- XENON1T, **[1512.07501](http://arxiv.org/abs/1512.07501)**
- Snowmass, **[1310.8327](http://arxiv.org/abs/1310.8327)**

Weakly Interacting Dark Matter

Hidden Sector Dark Matter

goal: explore possible cosmologies for thermal relics in hidden sectors

Gapped Hidden Sector

 \bullet LDP = DM

Cannibal DM

- LDP nonrelativistic at DM freezeout
- dark sector thermally decoupled from SM

Towards the Neutrino Floor

1. WIMP Warmup

- 2. Forbidden Dark Matter
- 3. Cannibal Dark Matter

1. WIMP Warmup

WIMP "Miracle"

WIMP "Miracle"

$$
\dot{n}_{\chi} + 3Hn_{\chi} = -\langle \sigma v \rangle \left(n_{\chi}^2 - (n_{\chi}^{eq})^2 \right)
$$

$$
n_{\chi} \langle \sigma v \rangle \approx H
$$

$$
\Omega_{\chi} h^2 \sim 0.1 \frac{m_{\chi} Y_{\chi}}{T_{eq}} \sim 0.1 \frac{m_{\chi} H}{T_{eq} s \langle \sigma v \rangle} \sim 0.1 \frac{(T_{eq} M_{pl})^{-1}}{\langle \sigma v \rangle}
$$

 $\sqrt{T_{eq}M_{pl}} \sim 60 \text{ TeV}$

Models of Light (Thermal) DM

$m_{DM} \ll m_h$

Models of Light (Thermal) DM

1. weakly coupled

- Pospelov, Ritz, Voloshin **[0711.4866](http://arxiv.org/abs/0711.4866)**
- Feng, Kumar **[0803.4196](http://arxiv.org/abs/0803.4196)**

2. asymmetric

- Nussinov, **[1985](http://inspirehep.net/search?p=recid:217983&of=hd)**
- Kaplan, Luty, Zurek, **[0901.4117](http://arxiv.org/abs/0901.4117)**
	- 3. SIMPs
- Hochberg, Kuflik, Volansky, Wacker, **[1402.5143](http://arxiv.org/abs/1402.5143)**

$$
m_{\chi} \sim \alpha_{eff} \left(T_{eq}^2 M_{pl}\right)^{1/3}
$$

 ~ 100 MeV

 $m_\chi \approx 5 \text{ GeV} \left(\frac{n_B - n_{\bar{B}}}{n_a - n_{\bar{B}}}\right)$ $n_\chi - n_{\bar{\chi}}$ ◆

 $\langle \sigma v \rangle \sim \frac{\alpha_d^2}{m^2}$ *d* m_χ^2 $\alpha_d \ll 1$

CMB limit

 \blacksquare **Combination, h302.01309** • Planck, **[1502.01589](http://arxiv.org/abs/1502.01589)**

2. Forbidden Dark Matter

• Raffaele D'Agnolo, JTR, **[1505.07107](http://arxiv.org/abs/1505.07107)**

- Griest and Seckel, 1991: "Forbidden Channel"
- evades CMB when: $T_{rec} \ll m_X + m_Y 2m_{DM}$

example model

$G_{SM} \times U(1)_d$

the company of the company of the company of the company of dudit divoo oudidi forbidden cross section

forbidden cross section

$$
\left<\sigma v\right>_{\psi\bar\psi}=\frac{(n^{eq}_{\gamma_d})^2}{(n^{eq}_{\psi})^2}\left<\sigma v\right>_{\gamma_d\gamma_d}
$$

$$
\gamma_d \sim \sim \sim \sqrt{\frac{1}{\psi}}
$$

$$
n^{eq} = g \left(\frac{m T}{2\pi}\right)^{3/2} e^{-m/T}
$$

$$
\left<\sigma v\right>_{\gamma_d\gamma_d}\sim \frac{\alpha_d^2}{m_{\gamma_d}^2}
$$

$$
\left\|\left\langle\sigma v\right\rangle_{\psi\bar{\psi}}\sim\frac{\alpha_d^2}{m_\psi^2}e^{-2x\Delta}\right\|
$$

 $\Delta \equiv$ $m_{\gamma_d} - m_\psi$ m_ψ $x \equiv$ m_ψ *T*

forbidden relic density

 $m_\psi \thicksim \alpha_d$ $\sqrt{T_{eq}M_{pl}}e^{-x_f\Delta}$

Three exceptions in the calculation of relic abundances

Kim Griest Center for Particle Astrophysics and Astronomy Department, University of California, Berkeley, California 94720

> David Seckel Bartol Research Institute, University of Delaware, Newark, Delaware 19716 (Received 15 November 1990)

- lies near in mass to the relic particle and shares a quantum number with it. An example is a light in \sim 1. coannihilation
- relic particle lies near a mass threshold. Previously, annihilation into particles heavier than the relic particle was considered that if the mass discussion Ω is the mass discussion of the mass discussion o 2. forbidden channels
- Proper treatment of the thermal averaging and the annihilation after freeze-out shows that the dip 3. annihilation near pole

forbidden relic density

$$
\Omega \propto \frac{m_{\psi}^2}{\alpha_d^2} e^{2x_f \Delta} \qquad m_{\psi} \sim \alpha_d \sqrt{T_{eq} M_{pl}} e^{-x_f \Delta}
$$

 $m_\psi \thicksim \alpha_d$ $\Omega \propto \frac{m_{\psi}}{r^2} e^{2x_f \Delta}$ $m_{\psi} \sim \alpha_d \sqrt{T_{eq} M_{pl}} e^{-x_f \Delta}$

Three exceptions in the calculation of relic abundances \boldsymbol{t}

Kim Griest Center for Particle Astrophysics and Astronomy Department, University of California, Berkeley, California 94720 Center for Particle Astrophysics and Astronomy Department, University of California, Berkeley, California 94720 Kim Griest

> David Seckel David Seckel Bartol Research Institute, University of Delaware, Newark, Delaware 19716 (Received 15 November 1990) (Received 15 November 1990)

The calculation of relic abundances of elementary particles by following their annihilation and The calculation of relic abundances of elementary particles by following their annihilation and freeze-out in the early Universe has become an important and standard tool in discussing particle dark-matter candidates. We find three situations, all occurring in the literature, in which the stan- $\frac{1}{\sqrt{1-\frac{1$ dard methods of calculating relic abundances fail. The first situation occurs when another particle lies near in mass to the relic particle and shares a quantum number with it. An example is a light squark with neutralino dark matter. The additional particle must be included in the reaction network, since its annihilation can control the relic abundance. The second situation occurs when the ϵ its annimization can control the form abundance. The second situation occurs
cla lies near a mass threshold. Previously, annihilation into particles heavier the diance the third situation occurs when the annihilation takes place near a pole in the cross section. ic particle was considered kinematically forbidden, but we show that if the mass difference is \sim 5–15%, these "forbidden" channels can dominate the cross section and determine the relic abunrelic particle lies near a mass threshold. Previously, annihilation into particles heavier than the reldance. The third situation occurs when the annihilation takes place near a pole in the cross section. Proper treatment of the thermal averaging and the annihilation after freeze-out shows that the dip in relic abundance caused by a pole is not nearly as sharp or deep as previously thought.

forbidden relic density

 $\angle e^-$

N N

e+

 \overline{y}

self-interactions *delf-interantion®*

 $\bar{\psi}$

 ψ

 N —

(velocity independent)

self-interactions

bullet cluster:

 σ_{SI} m_ψ < 1.25 cm²/g

• Randall et al., **[0704.0261](http://arxiv.org/abs/0704.0261v1)**

self-interactions

 m_ψ < 0.47 cm²/g

• Harvey et al., **[1503.07675](http://arxiv.org/abs/1503.07675)** $SINLOV$ of all this implies the coordinate system of the coordinate system of the coordinate system, ΔT position of the galaxies. The separation between galaxies and gas, SG, is shown in red. The

self-interactions

sensitivity:

$$
\frac{\sigma_{SI}}{m_{\psi}} \sim 1 \text{ cm}^2/g \sim 5 \times 10^{-6} \text{ MeV}^{-3}
$$

thermal annihilation rate:

$$
\langle \sigma v \rangle \sim 3 \times 10^{-3} \text{ TeV}^{-2}
$$

ratio:
$$
\frac{\sigma_{SI}}{\langle \sigma v \rangle} \sim 10^9 \left(\frac{m_{\psi}}{1 \text{ MeV}} \right)
$$

indirect detection: direct detection: e

forbidden parameter space

direct detection reach

• Essig, Mardon, Volansky **[1108.5383](http://arxiv.org/abs/1108.5383)** • Snowmass, **[1310.8327](http://arxiv.org/abs/1310.8327)**

Hidden Sector Taxonomy

non-gapped gapped

 T_d

cannibalism

Non-Gapped Hidden Sector

• entropy per comoving volume is separately conserved:

$$
s_d = \frac{2\pi^2}{45} g_{*S}^d T_d^3 \qquad s_{SM} = \frac{2\pi^2}{45} g_{*S}^{SM} T_\gamma^3
$$

$$
\xi = \frac{s_{SM}}{s_d}
$$

• temperature ratio:
$$
\frac{T_\gamma}{T_d} = \xi^{1/3} \left(\frac{g_{*S}^d}{g_{*S}^{SM}}\right)^{1/3} \sim \mathcal{O}(1)
$$

• Feng, Tu, Yu **[0808.2318](http://arxiv.org/abs/0808.2318)**

Cannibalism Conditions

1. hidden sector is kinetically decoupled from SM:

 $T_d \neq T_\gamma$

2. hidden sector has a mass gap: m_{ϕ} $\frac{\qquad \qquad }{\qquad \qquad }$

3. number changing interactions are in equilibrium when the hidden sector is non-relativistic:

 ϕ

 $T_d < m_{\phi}$

4. no chemical potential:

 $\mu_{\phi}=0$

Simplest Hidden Sector

Cannibal Sector Temperature

• entropy:

$$
s_d = \frac{\rho_d + p_d}{T_d} \approx \frac{m_\phi n_\phi}{T_d} \approx \frac{m_\phi^{5/2} T_d^{1/2}}{(2\pi)^{3/2}} e^{-m_\phi/T_d} \qquad s_{SM} = \frac{2\pi^2}{45} g_{*S}^{SM} T_\gamma^3
$$

• temperature ratio:

$$
\xi = \frac{s_{SM}}{s_d} \qquad \frac{T_{\gamma}}{T_d} \approx 0.5 \xi^{1/3} \left(g_{*}^{SM} \right)^{-1/3} \left(\frac{m_{\phi}}{T_d} \right)^{5/6} \boxed{e^{-m_{\phi}/3T_d}}
$$

• temperature vs. scale factor:

$$
T_{\gamma} \sim \frac{1}{a} \qquad T_d \sim \frac{1}{\log a}
$$

SELF-INTERACTING DARK MATTER

ERIC D. CARLSON Lyman Laboratory of Physics, Harvard University, Cambridge, MA 02138

MARIE E. MACHACEK Department of Physics, Northeastern University, Boston, MA 02115

AND

LAWRENCE J. HALL

Department of Physics, University of California; and Theoretical Physics Group, Physics Division, Lawrence Berkeley Laboratory, 1 Cyclotron Road, Berkeley, CA 94720 Received 1992 March 17; accepted 1992 April 20

the number density of particles. Hence number changing processes like $3 \rightarrow 2$ or $4 \rightarrow 2$ will tend to deplete the number of dark matter particles. But these processes take nonrelativistic particles in and produce (fewer) relativistic particles out, so that the outgoing particles have much more kinetic energy than the mean $(3/2)T'$. Hence subsequent $2 \rightarrow 2$ processes will transfer the kinetic energy of these few particles to all the dark matter, increasing the temperature. So as the universe expands, the dark matter cannibalizes itself to keep warm.

End of Cannibalism

decays

- · during cannibalism:
	- $\Gamma_{\phi} \ll H$
- end of cannibalism: $\Gamma_{\phi} \approx H$

Domination

 ϕ Dark Matter?

$$
\Omega_{\phi} h^2 \approx \frac{m_{\phi} n_{\phi}}{s_{SM}} (3.5 \text{ eV})^{-1} = \frac{m_{\phi}}{x_f \xi} (3.5 \text{ eV})^{-1}
$$

$$
x_f = \frac{m_{\phi}}{T_d^f} \quad \xi = \frac{s_{SM}}{s_d}
$$

• Carlson, Hall, Machacek, **[1992](http://inspirehep.net/record/333829)**.

• ϕ is too warm: $m_{\phi} = x_f \xi \times 0.4 \text{ eV} \lesssim 1 \text{ keV}$ (except for large ξ)

• DM from 2-to-2 freezeout in a cannibalizing sector:

 X annihilations are in equilibrium $\left(\begin{array}{c} 1 \end{array} \right)$ ϕ is relativistic

2) cannibalism starts when: $T_d < m_\phi$

-time \rightarrow

3) χ annihilations freezeout: $\chi \rightarrow \chi - \phi$

4) cannibalism ends when:

 $-\phi$ decays
- $\phi\phi \rightarrow \phi\phi$ freezeout

Relic Density

$$
\Omega_{\chi} h^2 \approx \frac{m_{\chi} n_{\chi}}{s_{SM}} \left(3.5 \text{ eV}\right)^{-1}
$$

$$
\text{freezeout:} \\
\chi \longrightarrow \text{--} \phi \\
\chi \longrightarrow \text{--} \phi \\
n_{\chi} \langle \sigma v \rangle = H
$$

$$
\frac{\Omega_{\chi}}{\Omega_{obs}} \approx 0.3 \ (g_{*}^{SM})^{-1/2} x_{f} \frac{\sigma_{0}}{\langle \sigma v \rangle} \frac{T_{d}}{T_{\gamma}} \sim \frac{\sigma_{0}}{\langle \sigma v \rangle} e^{3m_{\phi}/T_{d}^{f}}
$$

$$
x_{f} \equiv \frac{m_{\chi}}{T_{d}^{f}} \qquad \frac{T_{d}}{T_{\gamma}} \sim e^{m_{\phi}/3T_{d}^{f}}
$$

 $\sigma_0 = 3 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$

Relic Density

$$
\Omega_{\chi} \propto \langle \sigma v \rangle^{-1} e^{3m_{\phi}/T_d^f}
$$

Indirect Detection

 $\frac{y}{2} \phi \chi^2$

boosted cross: $\langle \sigma v \rangle \sim \sigma_0 e^{m_\phi/3T_d^f}$ *d*

- s-wave: $arg(y) \neq 0, \pi$ $\frac{y}{2} \phi \chi^2$ • s-wave: $\arg(y) \neq 0, \pi$
• p-wave: $\arg(y) = 0, \pi$
	- p-wave:

Cannibal DM Pheno

Dark Sector Phases

cannibal:

•Marco Farina, Duccio Pappadopulo, JTR, Gabriele Trevisan, *to appear.*

Forbidden DM Cannibal DM

take away

Forbidden DM

Cannibal DM

