

Dark Sectors with a Mass Gap

Josh Ruderman (NYU) @Rutgers 5/17/2016



- Raffaele D'Agnolo, JTR, 1505.07107
- Duccio Pappadopulo, JTR, Gabriele Trevisan, 1602.04219

Towards the Neutrino Floor



- XENON1T, **1512.07501**
- Snowmass, **1310.8327**

Weakly Interacting Dark Matter



Hidden Sector Dark Matter



goal: explore possible cosmologies for thermal relics in hidden sectors

Gapped Hidden Sector



• LDP = DM

Cannibal DM

- LDP nonrelativistic at DM freezeout
- dark sector thermally decoupled from SM

Towards the Neutrino Floor





1. WIMP Warmup

- 2. Forbidden Dark Matter
- 3. Cannibal Dark Matter

1. WIMP Warmup



WIMP "Miracle"



WIMP "Miracle"



$$\dot{n}_{\chi} + 3Hn_{\chi} = -\langle \sigma v \rangle \left(n_{\chi}^2 - (n_{\chi}^{eq})^2 \right)$$
$$n_{\chi} \langle \sigma v \rangle \approx H$$

$$n_{\chi} \left< \sigma v \right> pprox H$$

$$\Omega_{\chi}h^2 \sim 0.1 \frac{m_{\chi}Y_{\chi}}{T_{eq}} \sim 0.1 \frac{m_{\chi}H}{T_{eq}\,s\,\langle\sigma v\rangle} \sim 0.1 \frac{(T_{eq}M_{pl})^{-1}}{\langle\sigma v\rangle}$$

 $\sqrt{T_{eq}M_{pl}} \sim 60 \text{ TeV}$

Models of Light (Thermal) DM

$m_{DM} \ll m_h$

Models of Light (Thermal) DM

1. weakly coupled

- Pospelov, Ritz, Voloshin 0711.4866
- Feng, Kumar 0803.4196

2. asymmetric

- Nussinov, **1985**
- Kaplan, Luty, Zurek, 0901.4117
 - 3. SIMPs
- Hochberg, Kuflik, Volansky, Wacker, 1402.5143



$$m_{\chi} \sim \alpha_{eff} \left(T_{eq}^2 M_{pl}\right)^{1/3}$$

~ 100 MeV

$$m_{\chi} \approx 5 \text{ GeV}\left(\frac{n_B - n_{\bar{B}}}{n_{\chi} - n_{\bar{\chi}}}\right)$$

$$\alpha_d \ll 1$$

 $\langle \sigma v \rangle \sim \frac{\alpha_d^2}{m_{\gamma}^2}$

CMB limit



• Planck, 1502.01589

2. Forbidden Dark Matter



• Raffaele D'Agnolo, JTR, **1505.07107**



- Griest and Seckel, 1991: "Forbidden Channel"
- evades CMB when: $T_{
 m rec} \ll m_X + m_Y 2m_{DM}$

example model

$G_{SM} \times U(1)_d$



forbidden cross section



forbidden cross section

$$\langle \sigma v \rangle_{\psi \bar{\psi}} = \frac{(n_{\gamma_d}^{eq})^2}{(n_{\psi}^{eq})^2} \langle \sigma v \rangle_{\gamma_d \gamma_d} \qquad \gamma_d$$

$$\gamma_d \longrightarrow \psi$$

 $\gamma_d \longrightarrow \bar{\psi}$

$$n^{eq} = g\left(\frac{m\,T}{2\pi}\right)^{3/2} e^{-m/T}$$

$$\langle \sigma v \rangle_{\gamma_d \gamma_d} \sim \frac{\alpha_d^2}{m_{\gamma_d}^2}$$

$$\langle \sigma v \rangle_{\psi \bar{\psi}} \sim \frac{\alpha_d^2}{m_\psi^2} e^{-2x\Delta}$$

 $\Delta \equiv \frac{m_{\gamma_d} - m_{\psi}}{m_{\psi}}$ $x \equiv \frac{m_{\psi}}{T}$

forbidden relic density



 $m_{\psi} \sim \alpha_d \sqrt{T_{eq} M_{pl}} e^{-x_f \Delta}$

Three exceptions in the calculation of relic abundances

Kim Griest Center for Particle Astrophysics and Astronomy Department, University of California, Berkeley, California 94720

> David Seckel Bartol Research Institute, University of Delaware, Newark, Delaware 19716 (Received 15 November 1990)

- 1. coannihilation
- 2. forbidden channels
- 3. annihilation near pole

forbidden relic density

$$\Omega \propto \frac{m_{\psi}^2}{\alpha_d^2} e^{2x_f \Delta}$$

 $m_{\psi} \sim \alpha_d \sqrt{T_{eq} M_{pl}} e^{-x_f \Delta}$

Three exceptions in the calculation of relic abundances

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The calculation of relic abundances of elementary particles by following their annihilation and freeze-out in the early Universe has become an important and standard tool in discussing particle dark-matter candidates. We find three situations, all occurring in the literature, in which the standard methods of calculating relic abundances fail. The first situation occurs when another particle lies near in mass to the relic particle and shares a quantum number with it. An example is a light squark with neutralino dark matter. The additional particle must be included in the reaction network, since its annihilation can control the relic abundance. The second situation occurs when the relic particle lies near a mass threshold. Previously, annihilation into particles heavier than the relic particle was considered kinematically forbidden, but we show that if the mass difference is $\sim 5-15\%$, these "forbidden" channels can dominate the cross section and determine the relic abundance. The third situation occurs when the annihilation takes place near a pole in the cross section. Proper treatment of the thermal averaging and the annihilation after freeze-out shows that the dip in relic abundance caused by a pole is not nearly as sharp or deep as previously thought.

forbidden relic density





 e^+

ŀ

V

self-interactions





(velocity independent)

self-interactions

bullet cluster:



 $\frac{\sigma_{SI}}{-} < 1.25 \text{ cm}^2/\text{g}$ m_{ψ}

• Randall et al., 0704.0261

self-interactions





 $\frac{\sigma_{SI}}{m_{\psi}} < 0.47 \ \mathrm{cm}^2/\mathrm{g}$

• Harvey et al., **1503.07675**

self-interactions

sensitivity:

$$\frac{\sigma_{SI}}{m_{\psi}} \sim 1 \ \mathrm{cm}^2/g \sim 5 \times 10^{-6} \ \mathrm{MeV}^{-3}$$

thermal annihilation rate:

$$\langle \sigma v \rangle \sim 3 \times 10^{-3} \text{ TeV}^{-2}$$

ratio:
$$\frac{\sigma_{SI}}{\langle \sigma v \rangle} \sim 10^9 \left(\frac{m_{\psi}}{1 \text{ MeV}} \right)$$









indirect detection:

direct detection:







forbidden parameter space



direct detection reach



Essig, Mardon, Volansky 1108.5383
 Snowmass, 1310.8327







Hidden Sector Taxonomy

non-gapped

 T_d

gapped

cannibalism



Non-Gapped Hidden Sector

• entropy per comoving volume is separately conserved:

$$s_{d} = \frac{2\pi^{2}}{45} g_{*S}^{d} T_{d}^{3} \qquad s_{SM} = \frac{2\pi^{2}}{45} g_{*S}^{SM} T_{\gamma}^{3}$$
$$\xi = \frac{s_{SM}}{s_{d}}$$
• temperature ratio:
$$\frac{T_{\gamma}}{T_{d}} = \xi^{1/3} \left(\frac{g_{*S}^{d}}{g_{*S}^{SM}}\right)^{1/3} \sim \mathcal{O}(1)$$

• Feng, Tu, Yu 0808.2318

Cannibalism Conditions

1. hidden sector is kinetically decoupled from SM:

 $T_d
eq T_\gamma$

2. hidden sector has a mass gap: m_{ϕ}

3. number changing interactions are in equilibrium when the hidden sector is non-relativistic:

 $T_d < m_{\phi}$

4. no chemical potential:

 $\mu_{\phi} = 0$

Simplest Hidden Sector







Cannibal Sector Temperature

• entropy:

$$s_d = \frac{\rho_d + p_d}{T_d} \approx \frac{m_\phi n_\phi}{T_d} \approx \frac{m_\phi^{5/2} T_d^{1/2}}{(2\pi)^{3/2}} e^{-m_\phi/T_d} \qquad s_{SM} = \frac{2\pi^2}{45} g_{*S}^{SM} T_\gamma^3$$

• temperature ratio:

$$\xi = \frac{s_{SM}}{s_d} \quad \Longrightarrow \quad \frac{T_{\gamma}}{T_d} \approx 0.5 \, \xi^{1/3} \, (g_*^{SM})^{-1/3} \left(\frac{m_{\phi}}{T_d}\right)^{5/6} \left[e^{-m_{\phi}/3T_d} \right]^{1/3}$$

• temperature vs. scale factor:

$$T_{\gamma} \sim \frac{1}{a} \qquad T_d \sim \frac{1}{\log a}$$

SELF-INTERACTING DARK MATTER

ERIC D. CARLSON Lyman Laboratory of Physics, Harvard University, Cambridge, MA 02138

MARIE E. MACHACEK Department of Physics, Northeastern University, Boston, MA 02115

AND

LAWRENCE J. HALL

Department of Physics, University of California; and Theoretical Physics Group, Physics Division, Lawrence Berkeley Laboratory, 1 Cyclotron Road, Berkeley, CA 94720 Received 1992 March 17; accepted 1992 April 20

the number density of particles. Hence number changing processes like $3 \rightarrow 2$ or $4 \rightarrow 2$ will tend to deplete the number of dark matter particles. But these processes take nonrelativistic particles in and produce (fewer) relativistic particles out, so that the outgoing particles have much more kinetic energy than the mean (3/2)T'. Hence subsequent $2 \rightarrow 2$ processes will transfer the kinetic energy of these few particles to all the dark matter, increasing the temperature. So as the universe expands, the dark matter cannibalizes itself to keep warm.



End of Cannibalism



 $\phi decays$ $\frac{\phi F^2}{M} \phi \cdots \gamma \phi \gamma$

• during cannibalism:

 $\Gamma_{\phi} \ll H$

• end of cannibalism: $\label{eq:gamma} \frac{\Gamma_{\phi}}{\phi} \approx H$



Domination



Ø Dark Matter?

$$\Omega_{\phi}h^2 \approx \frac{m_{\phi}n_{\phi}}{s_{SM}} (3.5 \text{ eV})^{-1} = \frac{m_{\phi}}{x_f \xi} (3.5 \text{ eV})^{-1}$$
$$x_f = \frac{m_{\phi}}{T_d^f} \quad \xi = \frac{s_{SM}}{s_d}$$

• Carlson, Hall, Machacek, 1992.

• ϕ is too warm: $m_{\phi} = x_f \xi \times 0.4 \ {
m eV} \lesssim 1 \ {
m keV}$ (except for large ξ)

• DM from 2-to-2 freezeout in a cannibalizing sector:



1) χ annihilations are in equilibrium ϕ is relativistic





2) cannibalism starts when: $T_d < m_\phi$



3) χ annihilations freezeout: $\chi = \frac{\chi}{\chi} = \frac{-\phi}{-\phi}$





4) cannibalism ends when:

- ϕ decays - $\phi\phi\phi \rightarrow \phi\phi$ freezeout



Relic Density

$$\Omega_{\chi}h^2 \approx \frac{m_{\chi}n_{\chi}}{s_{SM}} \left(3.5 \text{ eV}\right)^{-1}$$

freezeout:

$$\chi - - - \phi$$

 $\chi - - - \phi$
 $n_{\chi} \langle \sigma v \rangle = H$

 $\sigma_0 = 3 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$

Relic Density

$$\Omega_{\chi} \propto \langle \sigma v \rangle^{-1} \ e^{3m_{\phi}/T_d^f}$$



Indirect Detection



boosted cross: $\langle \sigma v \rangle \sim \sigma_0 \, e^{m_\phi/3T_d^f}$

 $\frac{y}{2} \phi \chi^2$ • s-wave: $\arg(y) \neq 0, \pi$ • p-wave: $\arg(y) = 0, \pi$

Cannibal DM Pheno



Dark Sector Phases

cannibal: chemical: x_{ϕ} $T_{2\leftrightarrow 2}^{FO}$ $T^{FO}_{3\leftrightarrow 2}$ – $T_{3\leftrightarrow 2}^{FO} - T_{2\leftrightarrow 2}^{FO} - T_{\phi}^{decay}$ $T_{3\leftrightarrow 2}^{FO} = T_{\phi}^{\text{decay}}$ chemical cannibal T x_2 decay: chemical & $x_2 = \frac{m_{\chi}}{T_{2\leftrightarrow 2}^{FO}} \quad x_3 = \frac{m_{\chi}}{T_{3\leftrightarrow 2}^{FO}}$ decay $T_{\phi}^{
m decay}$ decay $-T_{2\leftrightarrow 2}^{FO}$ T $x_{\phi} = \frac{m_{\chi}}{T_{\perp}^{\text{decay}}}$ x_3 x_2

• Marco Farina, Duccio Pappadopulo, JTR, Gabriele Trevisan, to appear.



Forbidden DM

Cannibal DM





take away

Forbidden DM

Cannibal DM

