

# PRECISION HIGGS PHYSICS AT COLLIDERS

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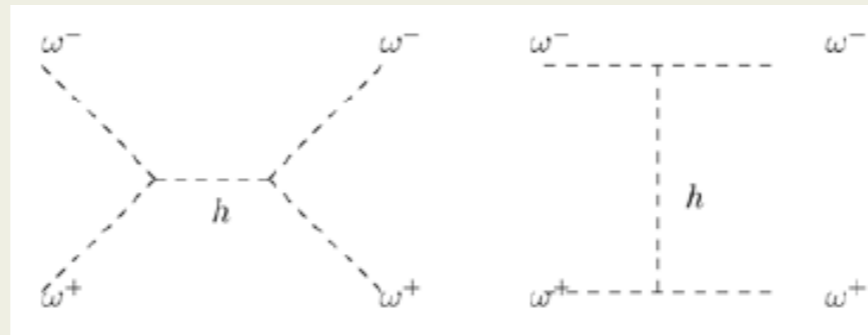
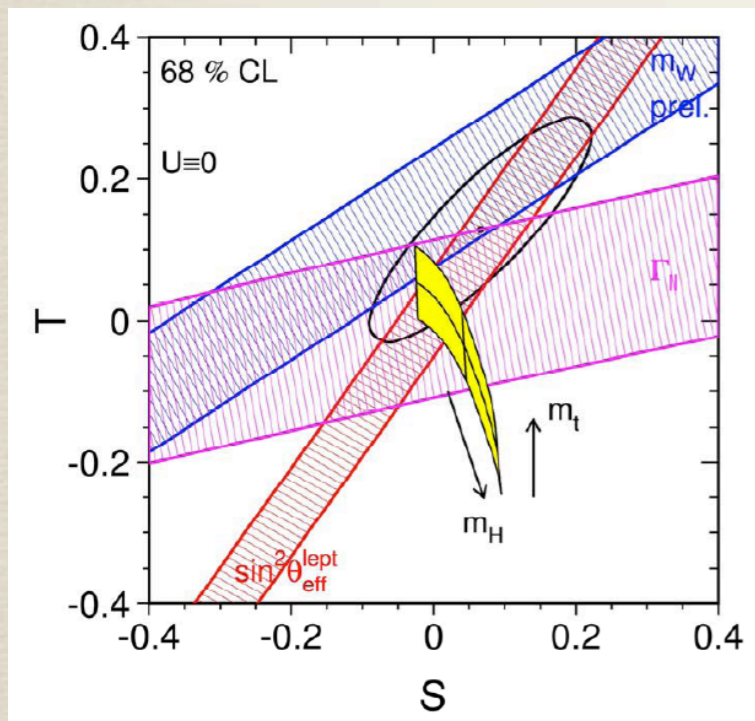
October 20, 2009

# Outline

- Review of experiment, theory for SM Higgs
- The gluon-fusion mechanism: a precision QFT playground
- Electroweak corrections and the Higgs effective theory at 3-loops  
Anastasiou, Boughezal, FP 0811.3458
- Updated numerics, the Tevatron exclusion limit, and fun with PDFs  
Anastasiou, Boughezal, FP 0811.3458
- Electroweak and quark-mass effects at high Higgs  $p_T$   
W.-Y. Keung, FP 0905.2775
- Low- $p_T$  resummation using soft-collinear effective theory  
S. Mantry, FP, in progress

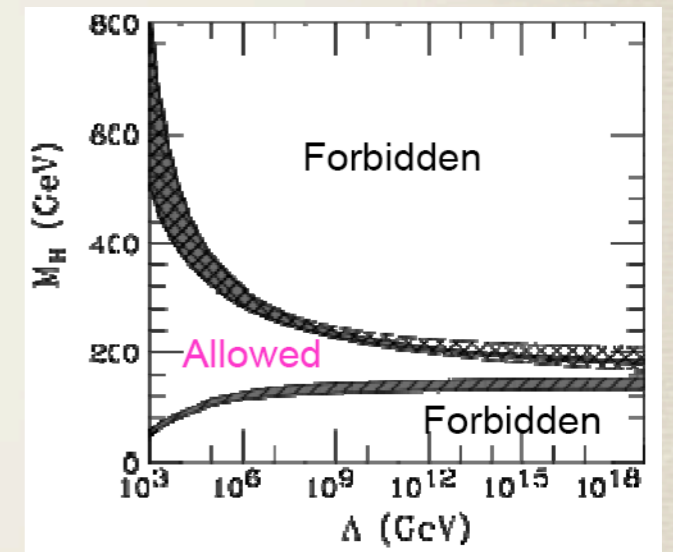
# Why we expect a TeV scale Higgs

- Last undiscovered particle of the SM
- Many reasons to expect it (or something else) to be observed soon



$$a_0^0 \rightarrow -\frac{s}{32\pi v^2}$$

$$\Lambda_{NP} \leq 1.7 \text{ TeV}$$

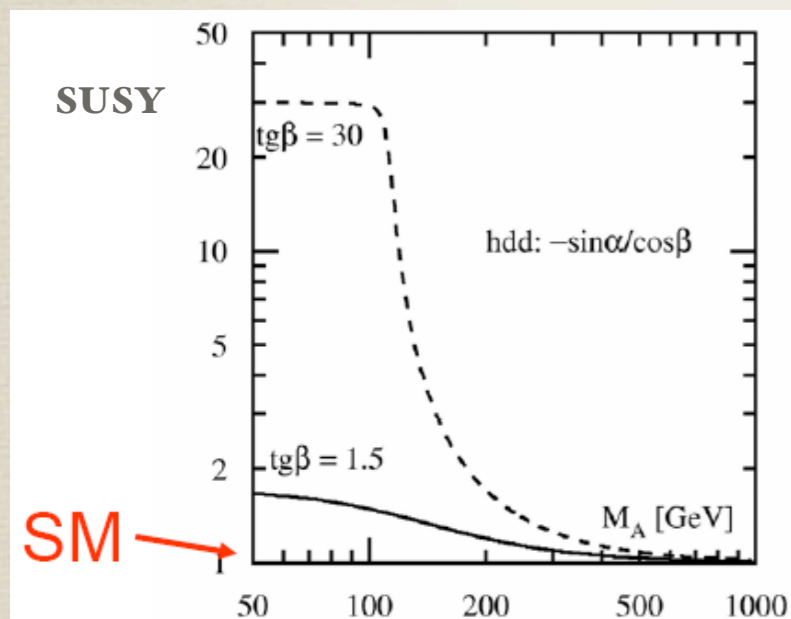


$$M_H^2 < \frac{32\pi^2 v^2}{9 \log\left(\frac{Q^2}{v^2}\right)}$$

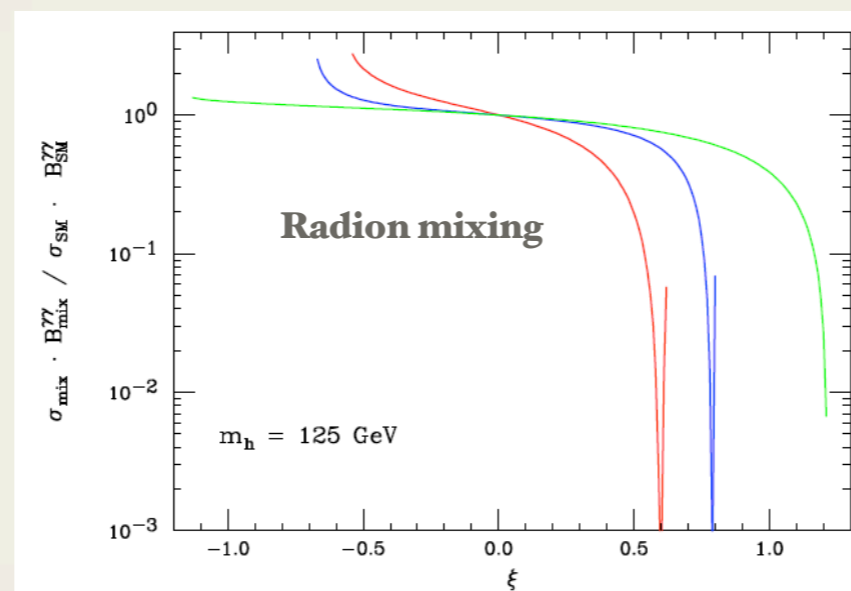
$$M_H^2 > \frac{3v^2}{2\pi^2} \log\left(\frac{\Lambda^2}{v^2}\right)$$

# Higgs in SM extensions

The uncertainty in EWSB mechanism makes Higgs a portal into new physics at the TEV scale



S. Dawson



Hewett, Rizzo hep-ph/0202155

## NMSSM

$m_{h_1}/m_{a_1}$ (GeV)	Branching Ratios			$n_{\text{obs}}/n_{\text{exp}}$ units of $1\sigma$	$s_{95}$	$N_{SD}^{LHC}$
	$h_1 \rightarrow b\bar{b}$	$h_1 \rightarrow a_1 a_1$	$a_1 \rightarrow \tau\bar{\tau}$			
98.0/2.6	0.062	0.926	0.000	2.25/1.72	2.79	1.2
100.0/9.3	0.075	0.910	0.852	1.98/1.88	2.40	1.5
100.2/3.1	0.141	0.832	0.000	2.26/2.78	1.31	2.5
102.0/7.3	0.095	0.887	0.923	1.44/2.08	1.58	1.6
102.2/3.6	0.177	0.789	0.814	1.80/3.12	1.03	3.3
102.4/9.0	0.173	0.793	0.875	1.79/3.03	1.07	3.6
102.5/5.4	0.128	0.848	0.938	1.64/2.46	1.24	2.4
105.0/5.3	0.062	0.926	0.938	1.11/1.52	2.74	1.2

Dermisek, Gunion hep-ph/0510322

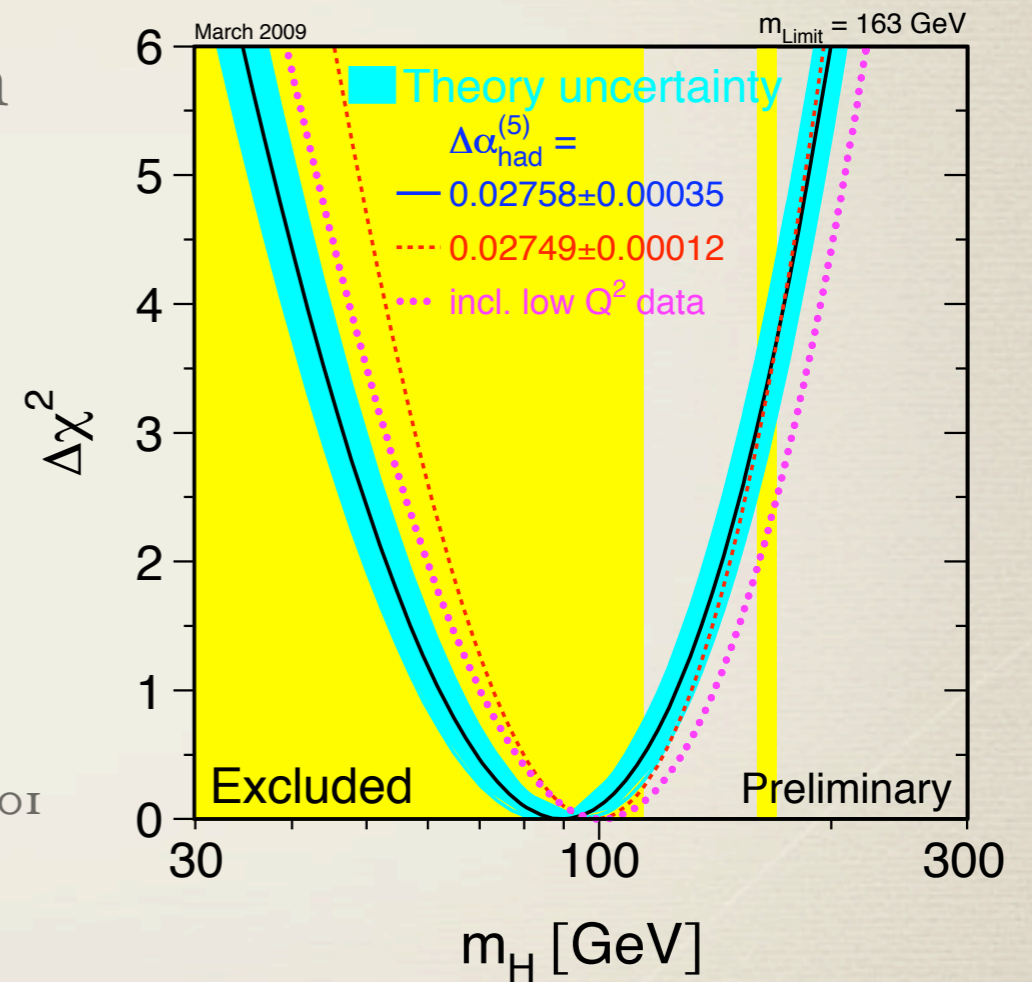
- Loop-induced gluon, photon modes can have  $O(1)$  deviations
- Non-standard decays can drastically change collider signals

# SM Higgs circa 2009

Precision EW upper bound and direct search lower bound at 95% CL:

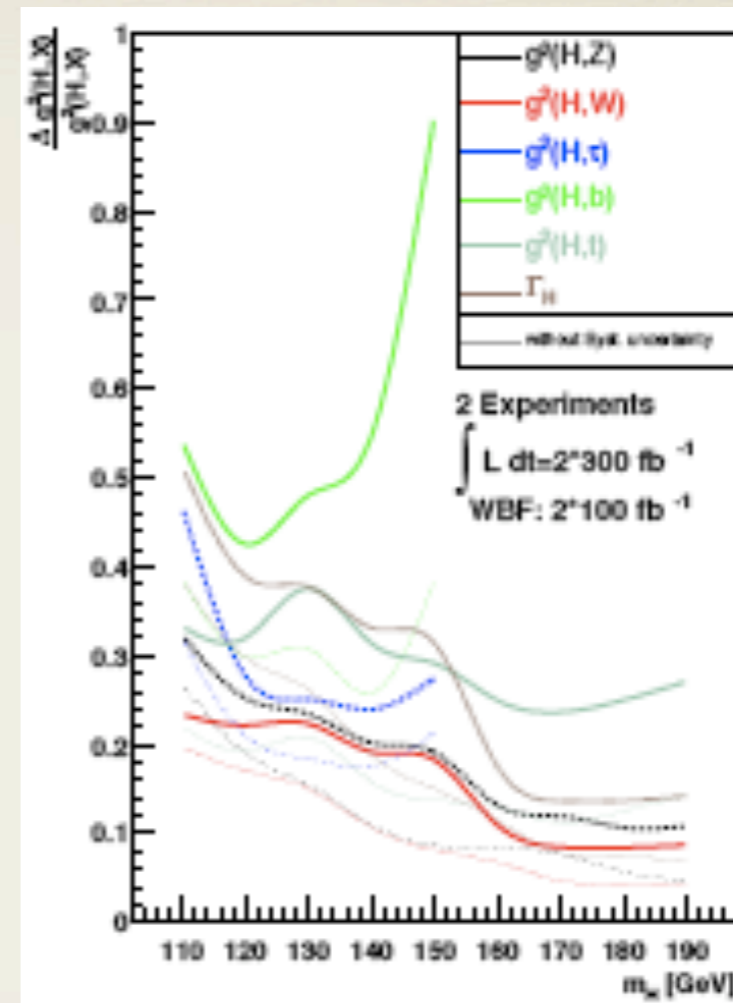
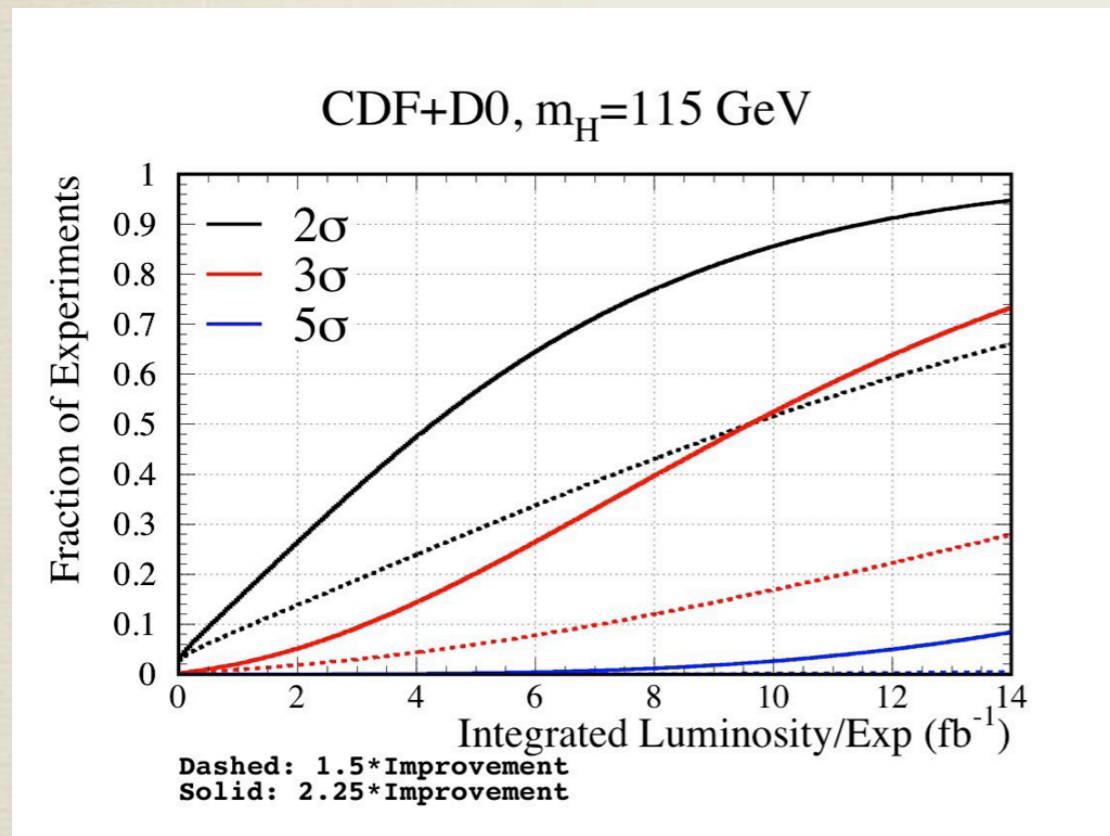
$$114 < M_H / \text{GeV} < 163$$

**News from the Tevatron:** First exclusion in 2008; new combined results exclude 160-170 GeV SM Higgs at 95% CL [arXiv:0903.4001](https://arxiv.org/abs/0903.4001)



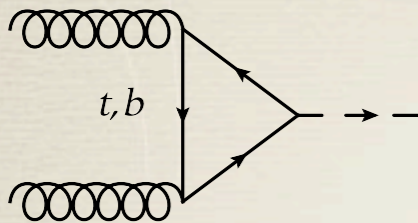
# The Higgs in the future

Duhrssen et al., hep-ph/0406323

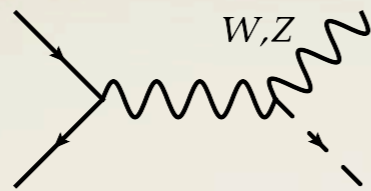


- Discovery program over entire mass range at Tevatron+LHC
- Coupling measurements to 10% possible; spin, CP

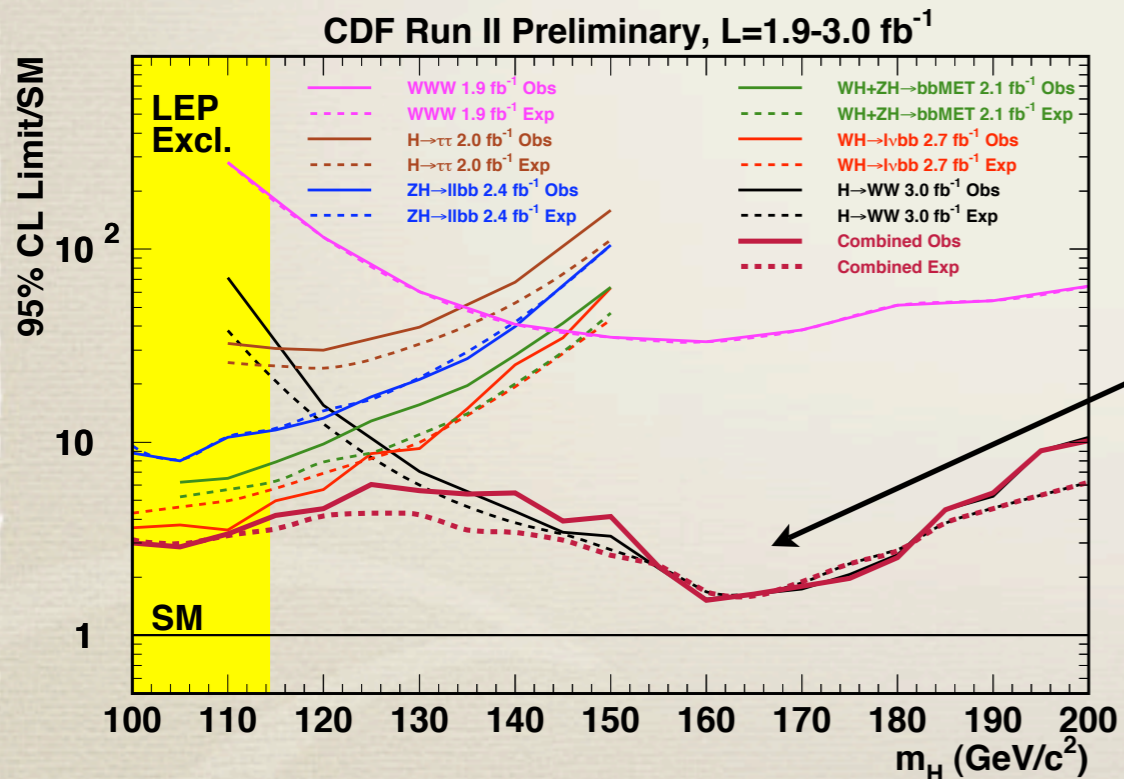
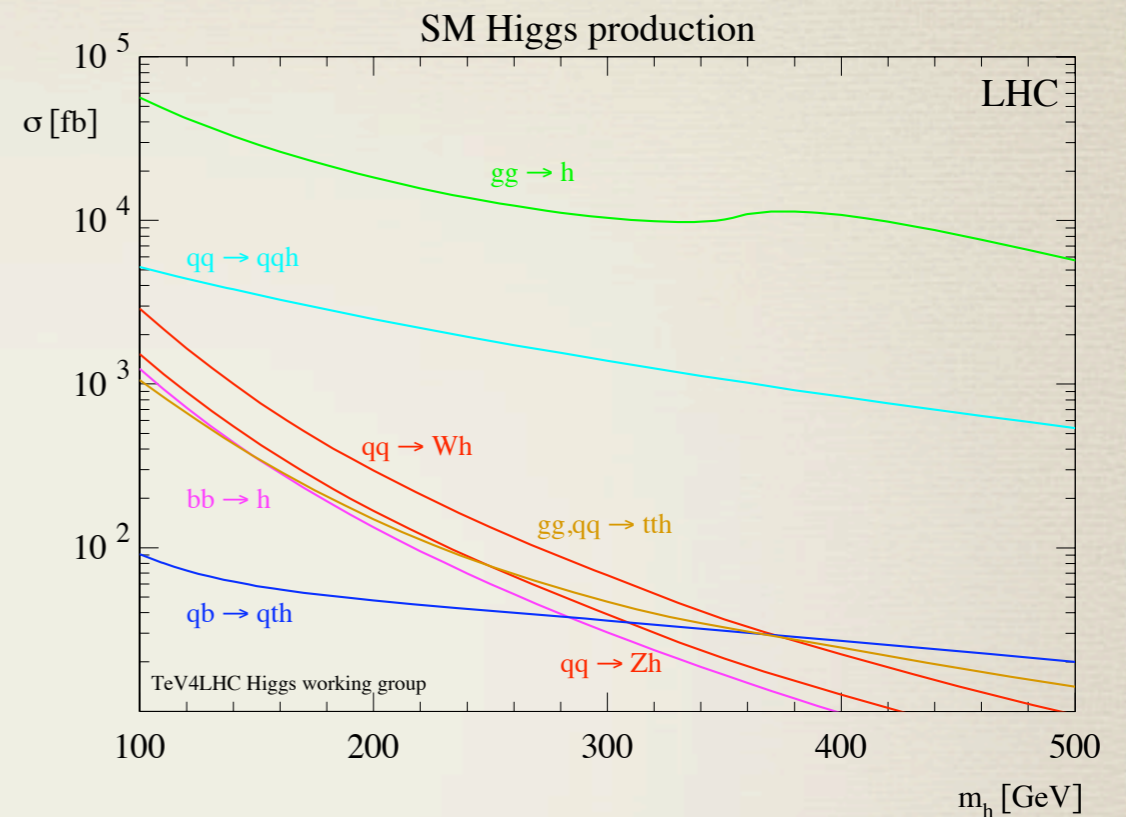
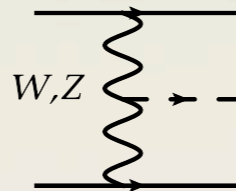
# SM Higgs production



gg fusion  
dominant by  
factor of 10



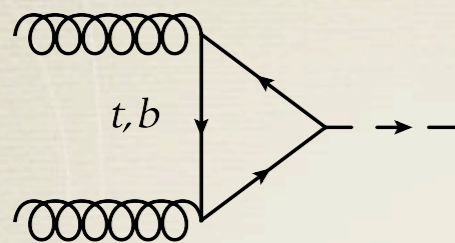
Associated production,  
WBF essential for  
light Higgs



Tevatron exclusion limit entirely  
from  $gg \rightarrow H \rightarrow WW$  above 130 GeV

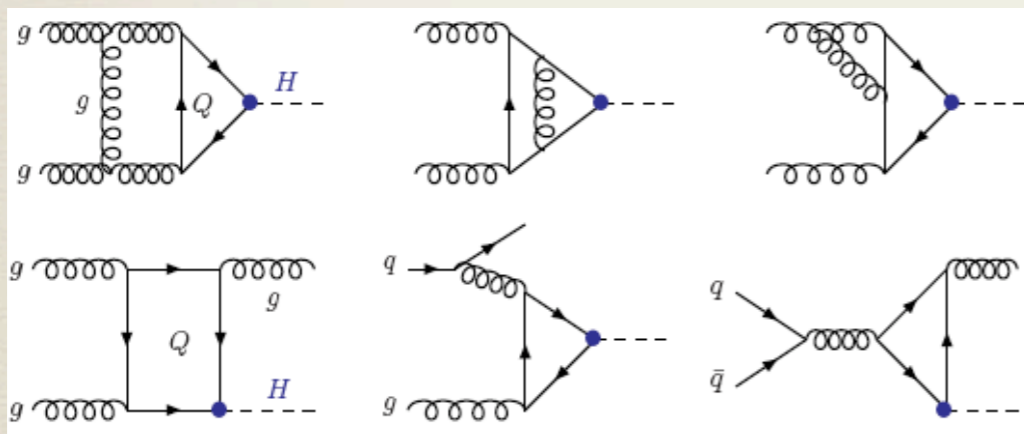
$BR(H \rightarrow WW) > 90\%$  for 160-170  
GeV Higgs

# Gluon-fusion at NLO

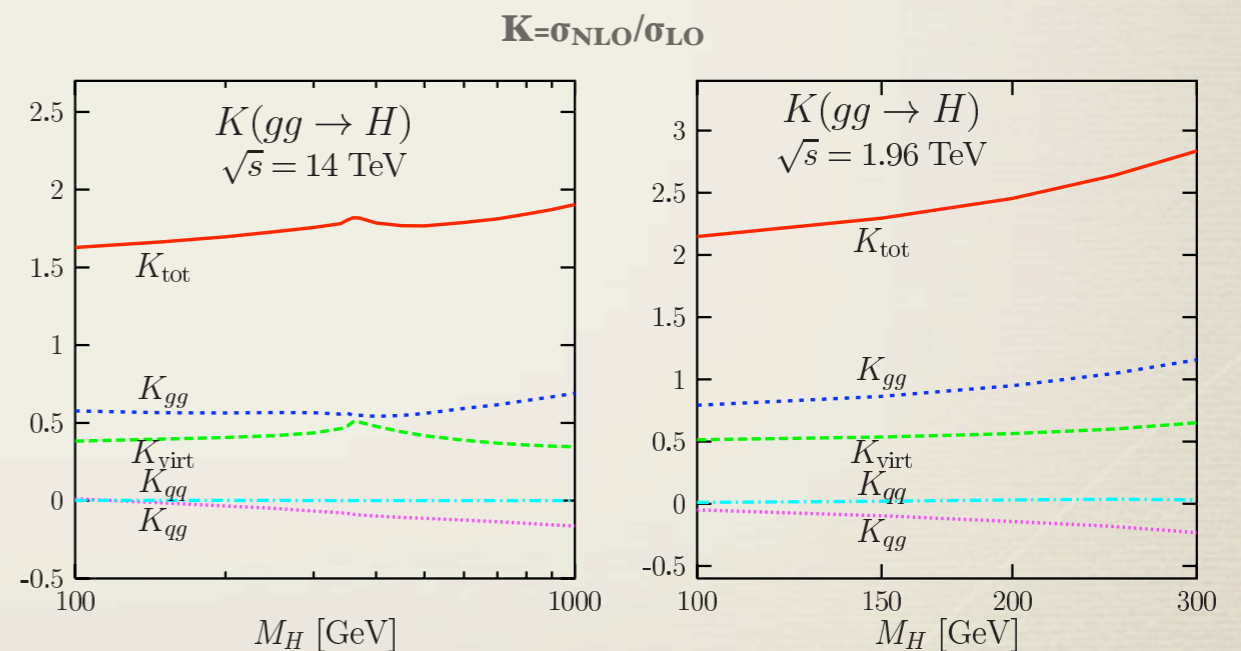


Top-loop dominant; bottom loop gives -5% correction from interference

What makes it sensitive to new physics (begins at 1-loop) also makes it tough to calculate...



Djouadi, Graudenz, Spira, Zerwas PLB 264 (1991), hep-ph/9504378

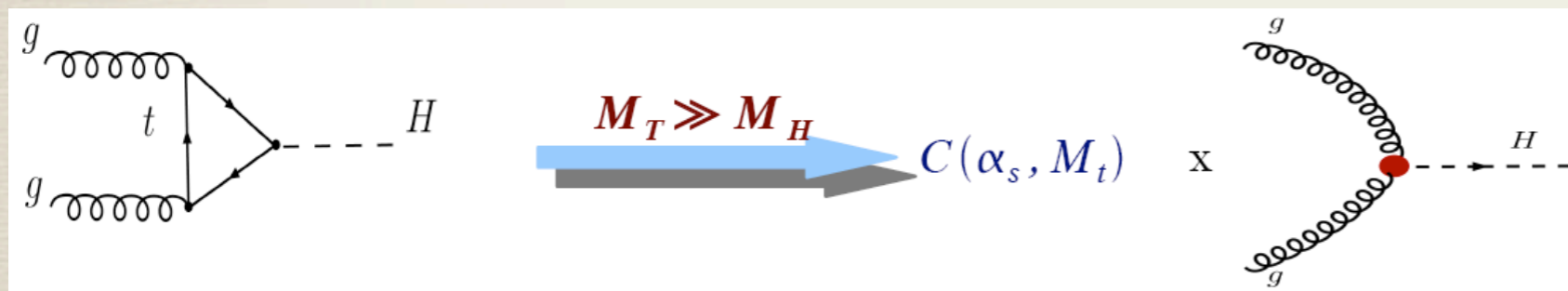


Can reach  $K_{\text{NLO}} = \sigma_{\text{NLO}} / \sigma_{\text{LO}} \approx 2$  at LHC, 3 at Tevatron; why so large?



# Effective theory for Higgs

- Difficult to go to NNLO and check convergence of expansion
- NLO analytic expressions unwieldy to study, develop intuition about
- Use EFT instead for top (Shifman et al. 1979; Ellis et al. 1988; S. Dawson; Djouadi, Spira, Zerwas 1991)

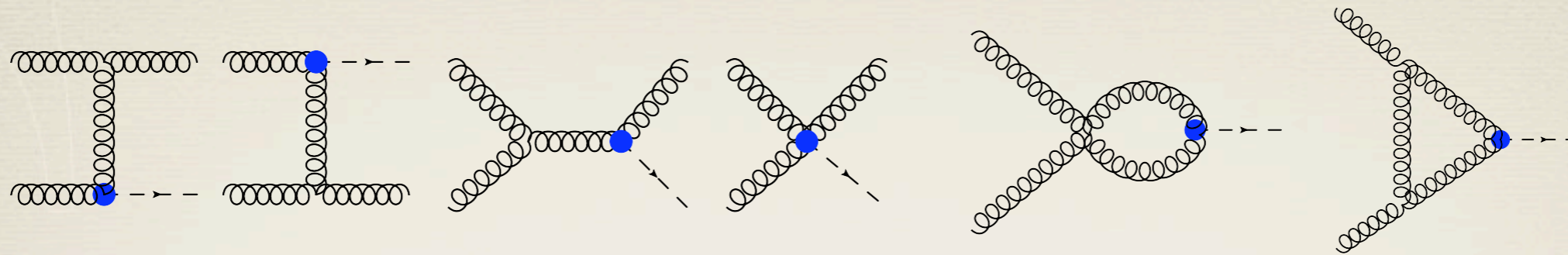


Region of validity:  $M_H < 2m_t$

$$L_{ggh} = \frac{-H}{4v} C(\alpha_s) G_{\mu\nu}^a G_a^{\mu\nu}$$

known through  $O(\alpha_s^5)$ :  
 Schroder, Steinhauser; Chetyrkin,  
 Kuhn, Sturm hep-ph/0512058, 0512060

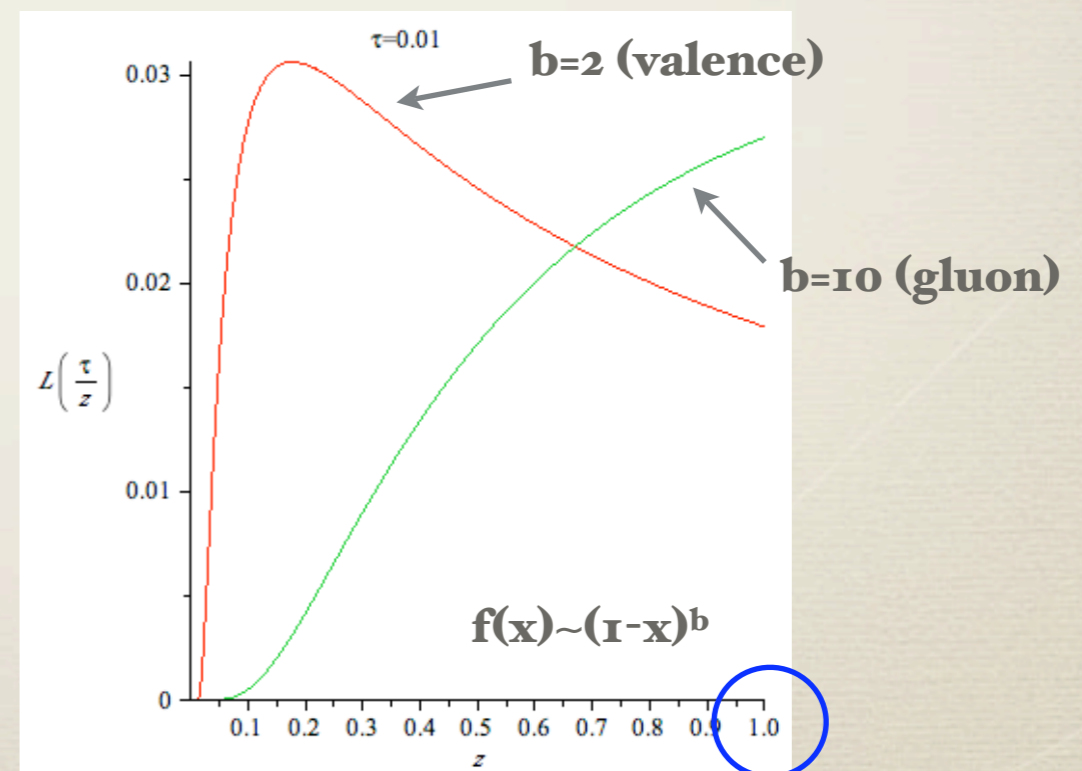
# NLO in the EFT



$$\Delta\sigma = \sigma_0 \frac{\alpha_s}{\pi} \left\{ \left( \frac{11}{2} + \pi^2 \right) \delta(1-z) + 12 \left[ \frac{\ln(1-z)}{1-z} \right]_+ - 12z(-z + z^2 + 2)\ln(1-z) - 6 \frac{(z^2 + 1 - z)^2}{1-z} \ln(z) - \frac{11}{2} (1-z)^3 \right\} \quad (M^2/s \leq z \leq 1)$$

(integration over PDFs  $\Rightarrow$  integration over  $z$ )

- First source of large correction:  
 $11/2 + \pi^2 \Rightarrow 50\%$  increase (can even resum these corrections Magnea, Sterman 1990; Ahrens, Becher, Neubert, Yang 0808.3008)
- Second source: shape of PDFs enhances *threshold* logarithm



# Unreasonably effective EFT

- In the full theory with top quarks, study eikonal approximation for real emission, Sudakov form factor for virtual

$$\Delta\sigma_{top}^{soft} = \sigma_{top}^{(0)} C_A \frac{\alpha_s}{\pi} (1-z)^{-1-2\epsilon} \lambda^{-1-\epsilon} (1-\lambda)^{-1-\epsilon} \Rightarrow \sigma_{top}^{(0)} \frac{\alpha_s}{\pi} 12 \left[ \frac{\ln(1-z)}{1-z} \right]_+$$

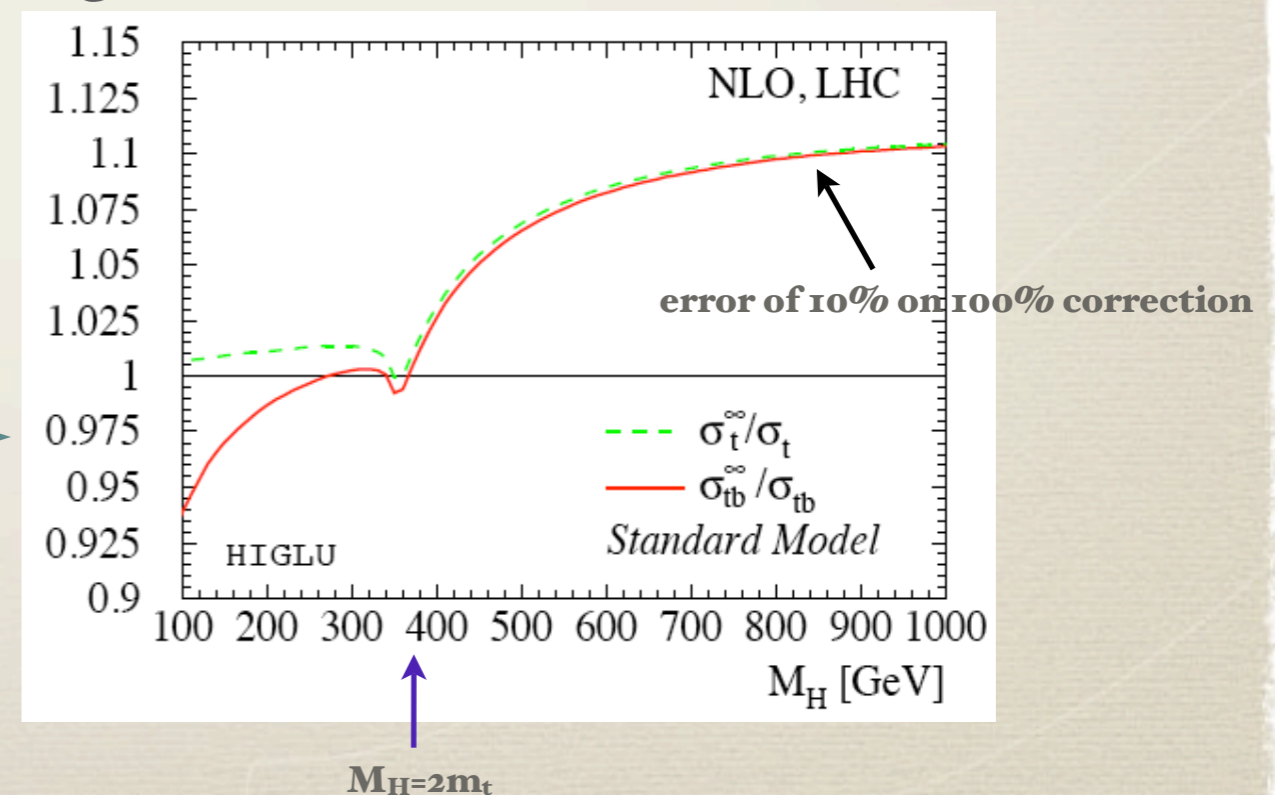
$$\Delta\sigma_{top}^{Sud} = \sigma_{top}^{(0)} 2C_A \frac{\alpha_s}{\pi} \frac{1}{\epsilon^2} \left( \frac{-\mu^2}{\hat{s}} \right)^\epsilon \delta(1-z) \Rightarrow \sigma_{top}^{(0)} \frac{\alpha_s}{\pi} \pi^2 \delta(1-z)$$

- Dominant corrections in full and EFT differ only by tree-level cross section  $\Rightarrow$  caused by shape of gluon PDF

$$\sigma_{NLO}^{approx} = \left( \frac{\sigma_{NLO}^{EFT}}{\sigma_{LO}^{EFT}} \right) \sigma_{LO}^{QCD}$$

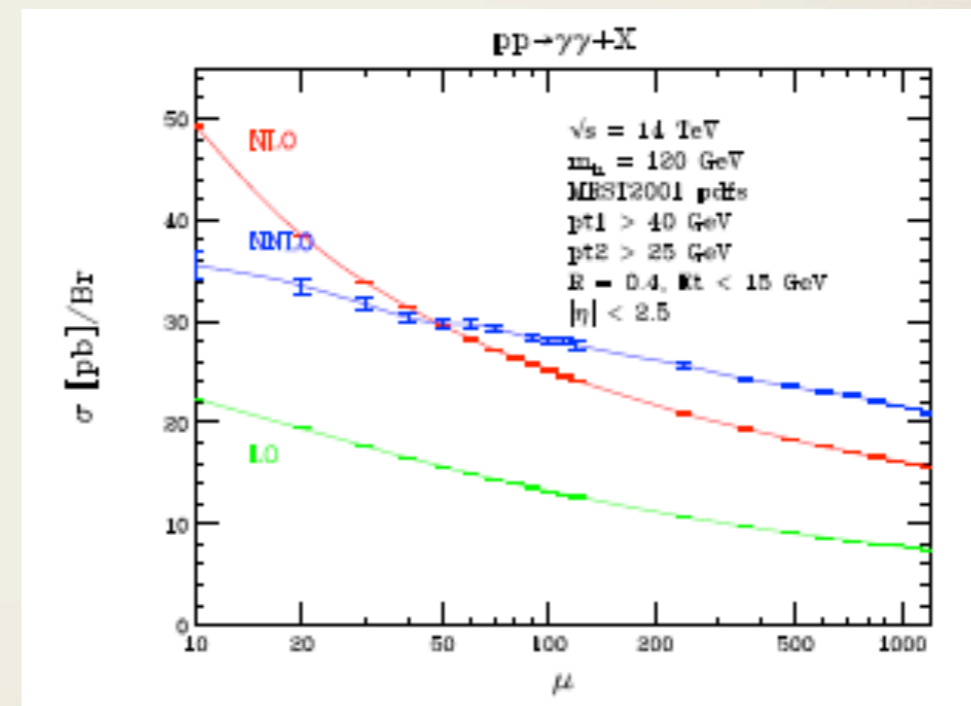
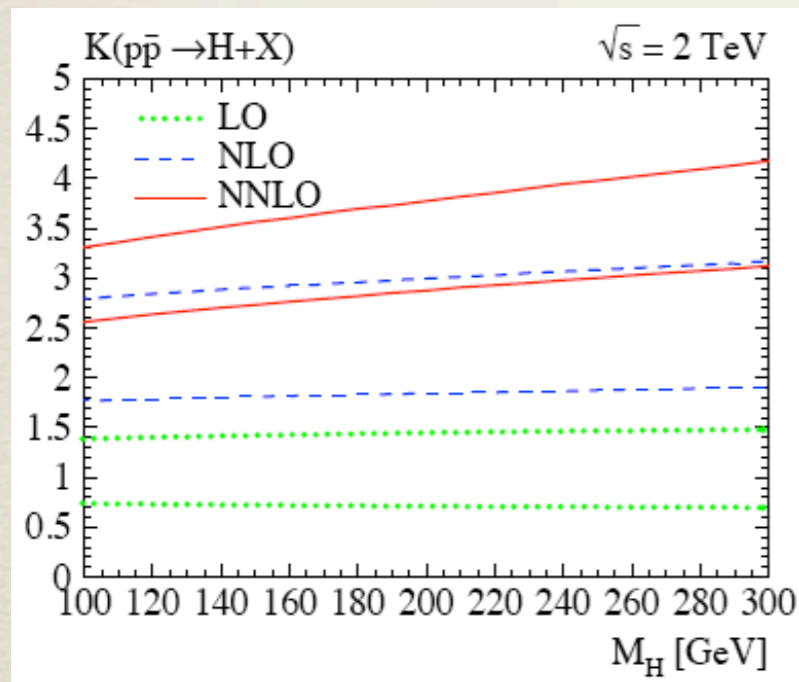
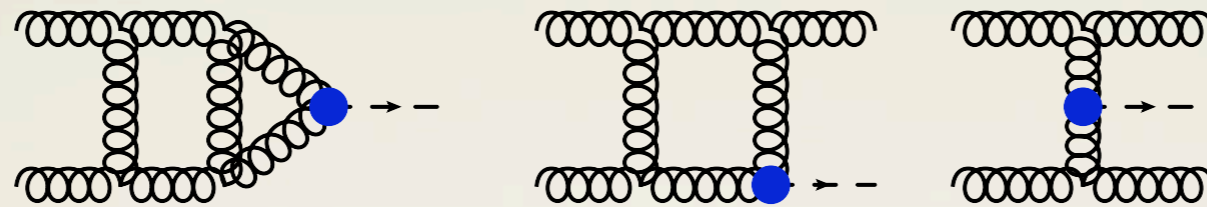
%-level or better for  $M_H < 2m_t$ , even gets >90% of correction above Threshold structure preserved

Initial NNLO study of  $1/m_t$  suppressed operators indicates this persists Harlander, Ozeren 0909.3420



# NNLO in the EFT

- Motivates calculation to NNLO in the EFT



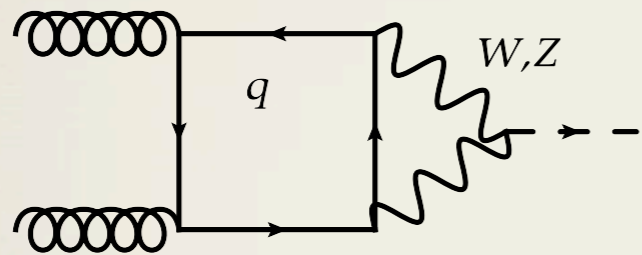
Harlander, Kilgore; Anastasiou, Melnikov;  
Ravindran, Smith, van Neerven 2002-2003

Anastasiou, Melnikov, FP 2005

K-factor: 2 at LHC, 3.5 at Tevatron

# Electroweak corrections

- Residual QCD uncertainty  $\sim 10\%$   $\Rightarrow$  EW corrections potentially important to match QCD and experimental precision
- $N_F$ -enhanced sourced of 2-loop light-quark corrections



$\Rightarrow$  Up to  $\sim 8-9\%$  at threshold relative to LO QCD

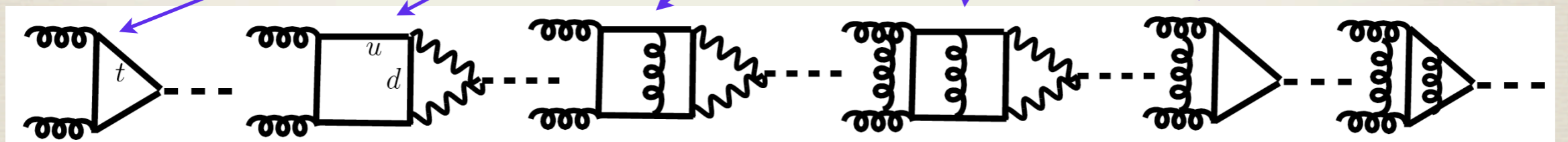
Aglietti, Bonciani, Degrassi, Vicini hep-ph/0404071;  
Actis, Passarino, Sturm, Uccirati 0809.1301

- K-factor? Values between 1-4 assumed in literature; do these get same K-factor of top-quark piece?
- First goal: check with 3-loop calculation in EFT

# EFT formulation

$$\mathcal{L} = -\alpha_s \frac{C_1}{4v} H G_{\mu\nu}^a G^{a\mu\nu}$$

$$C_1 = -\frac{1}{3\pi} \left\{ 1 + \lambda_{EW} \left[ 1 + a_s C_{1w} + a_s^2 C_{2w} \right] + a_s C_{1q} + a_s^2 C_{2q} \right\}$$



- Radius of convergence:  $M_H \leq M_W \dots$
- However, dominant corrections from threshold logs and analytic continuation of Sudakov form factor identical in full and EFT
- Calculate K-factor in EFT, normalize to exact 2-loop EW result

# Factorization in the EFT

- If the K-factor for light-quark pieces is the same as the top quark, then the Wilson coefficient in the EFT “factorizes”

$$C_1 = -\frac{1}{3\pi} \left\{ 1 + \lambda_{EW} \left[ 1 + a_s C_{1w} + a_s^2 C_{2w} \right] + a_s C_{1q} + a_s^2 C_{2q} \right\}$$

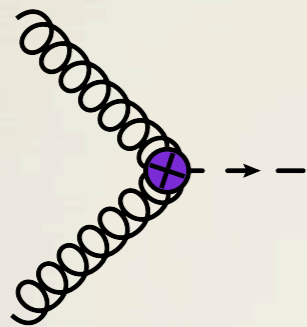


$$C_1^{fac} = -\frac{1}{3\pi} (1 + \lambda_{EW}) \left\{ 1 + a_s C_{1q} + a_s^2 C_{2q} \right\}$$

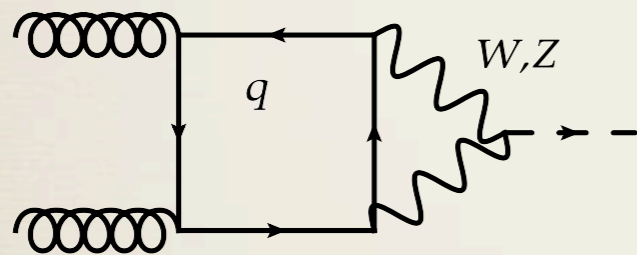
- Factorization holds if  $C_{1w} = C_{1q}$ ;  $C_{1q} = 11/4$
- Calculate  $C_{1w}$  from 3-loop diagrams, check deviation from  $C_{1q}$ , study numerical effect

# Matching to the EFT I

• Matching at  $O(\alpha\alpha_s)$ :



$$= -\frac{1}{3\pi} \frac{\alpha_s}{v} \lambda_{EW} \mathcal{M}_0$$



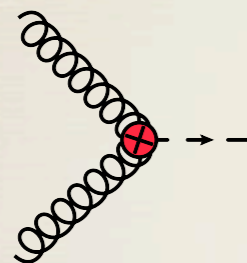
$$= \mathcal{A}^{(2)}(M_H^2 = 0) \mathcal{M}_0 + \mathcal{O}\left(\frac{M_H^2}{M_{W,Z}^2}\right)$$

☑ Equate to get  $\lambda_{EW}$

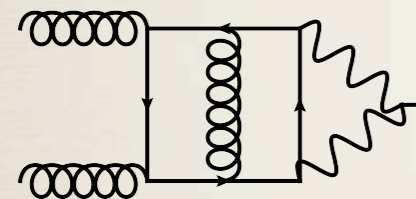


# Matching to the EFT II

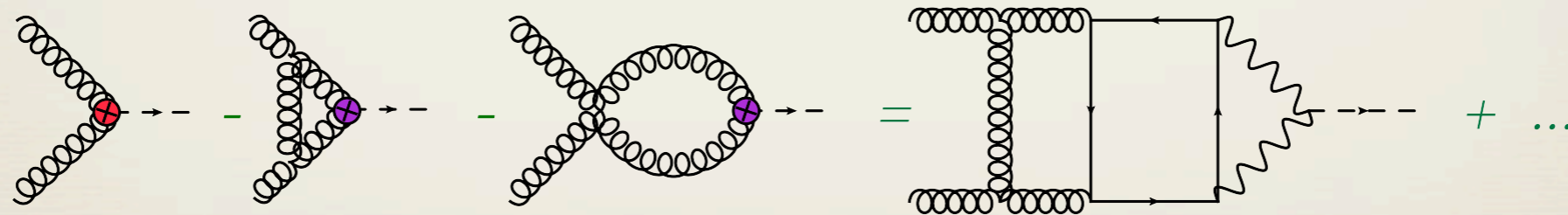
📌 Matching at  $O(\alpha\alpha_s^2)$ :



$$= -\frac{1}{3\pi} \frac{\alpha_s}{v} \lambda_{EW} (\alpha_s C_{1w}) \mathcal{M}_0$$



$$= \mathcal{A}^{(3)}(M_H^2 = 0) \mathcal{M}_0 + \mathcal{O}(M_H^2/M_{W,Z}^2)$$

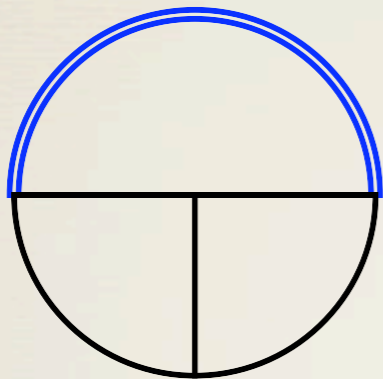


$$= \text{[Diagram 1]} - \text{[Diagram 2]} - \text{[Diagram 3]} = \text{[Diagram 4]} + \dots$$

☑ Equate to get  $C_{1w}$

# Computational strategy

- Expansion in  $M_H/M_W$  reduces diagrams to 3-loop vacuum bubbles



$$\begin{aligned} \mathcal{I}(\vec{\nu}_i) &= \int \prod_{j=1}^3 d^d k_j \frac{1}{k_1^{2\nu_1} k_2^{2\nu_2} (k_3^2 - M_{W,Z}^2)^{\nu_3} (k_1 - k_2)^{2\nu_4} (k_2 - k_3)^{2\nu_5} (k_3 - k_1)^{2\nu_6}} \\ &= \int \prod_{j=1}^3 d^d k_j \mathcal{D} \end{aligned}$$

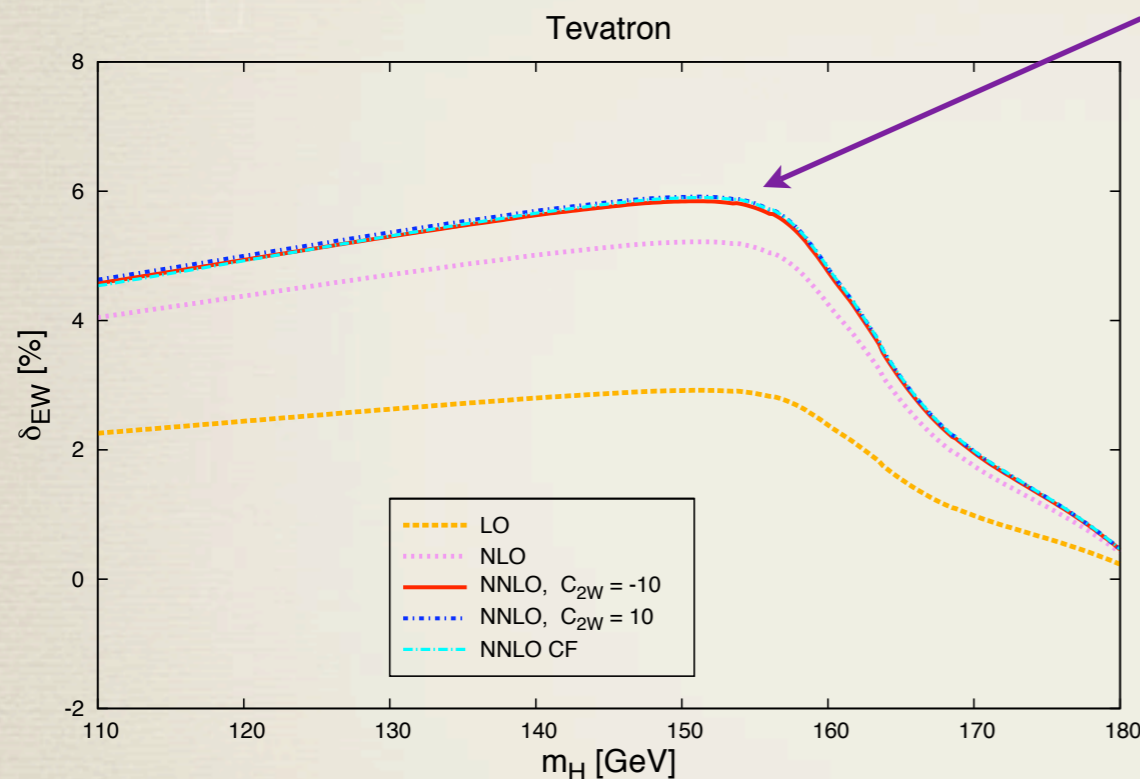
- Use Poincare invariance of loop integrals to facilitate calculation

K. Chetyrkin, F. Tkachov 1981  $\Rightarrow$  left with two such integrals to evaluate

$$\int \prod_{j=1}^3 d^d k_j \partial_i [k_k \mathcal{D}] = 0$$

# Factorization violation

- Analytical result:  $C_{IW} = 7/6$ , compared to  $C_{IQ} = 11/4$



Difference between factorization hypothesis and actual result irrelevantly small (weak violation)

$$\sigma_{3-loop} = \sigma_{2-loop} \left\{ \frac{\alpha_s}{\pi} C_{1w} + G_{EFT}^{(1)} \right\}$$

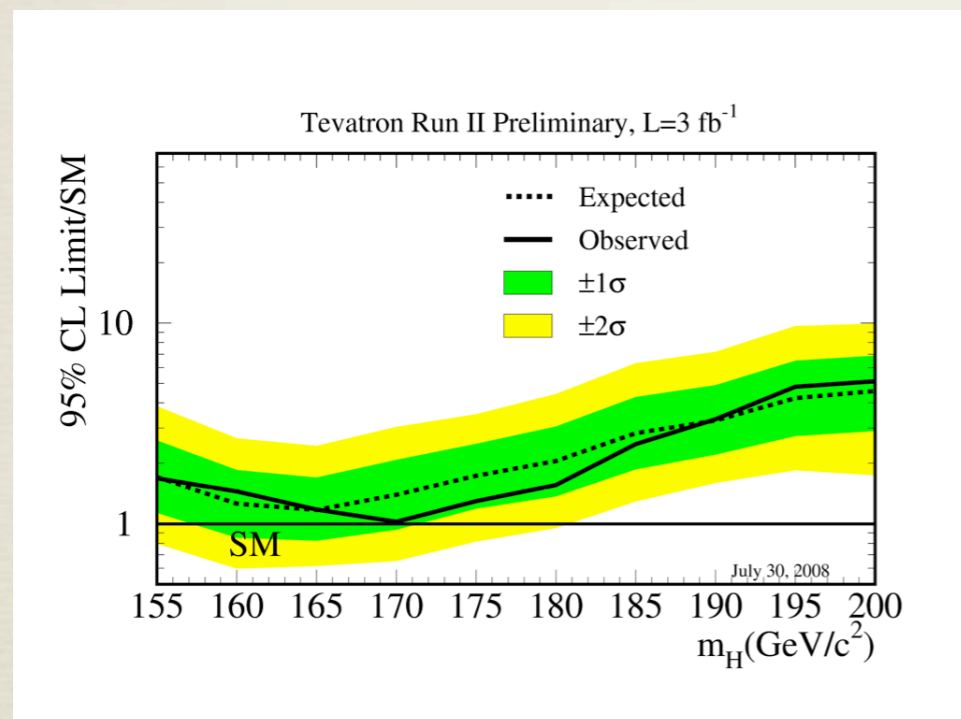
$\alpha_s(C_{IW} - C_{IQ})/\pi \approx 5\%$

$G_{EFT}^{(1)} \approx 100\%$ ; contains  $\pi^2$ ,  $\ln(1-z)$ , the large corrections to HGG operator

- K-factor of 3.5 at Tevatron appropriate

# Circa December 2008

- Combine QCD, EW corrections to derive current best prediction, check what is in Tevatron analysis
- First limit:  $M_H=170$  GeV excluded 0808.0534



- Same K-factors assumed for top, EW contributions ✓
- Same K-factor assumed for top, bottom quarks  
 $K_{tb} \sim 1.5, K_{tt} \sim 3.5 \Rightarrow$  needed updating
- MRST 2002 PDFs used  
significant changes in heavy-quark threshold treatment

What they used →

original	MRST 2006 PDFs	$K_{tb}, K_{bb}$	EW effects
0.3542	0.3650	0.3868	0.3943

← Update #1

$M_H=170$  GeV

10% increase

# Circa January 2009

MSTW 2008 PDF release [arXiv:0901.0002](https://arxiv.org/abs/0901.0002)

- Run II inclusive jet data
- Decrease of  $\alpha_s(M_Z)$  from 0.119  $\rightarrow$  0.117
- Gluon density decreased at  $x \sim 0.1$
- gg luminosity error increased from 5%  $\Rightarrow$  10%

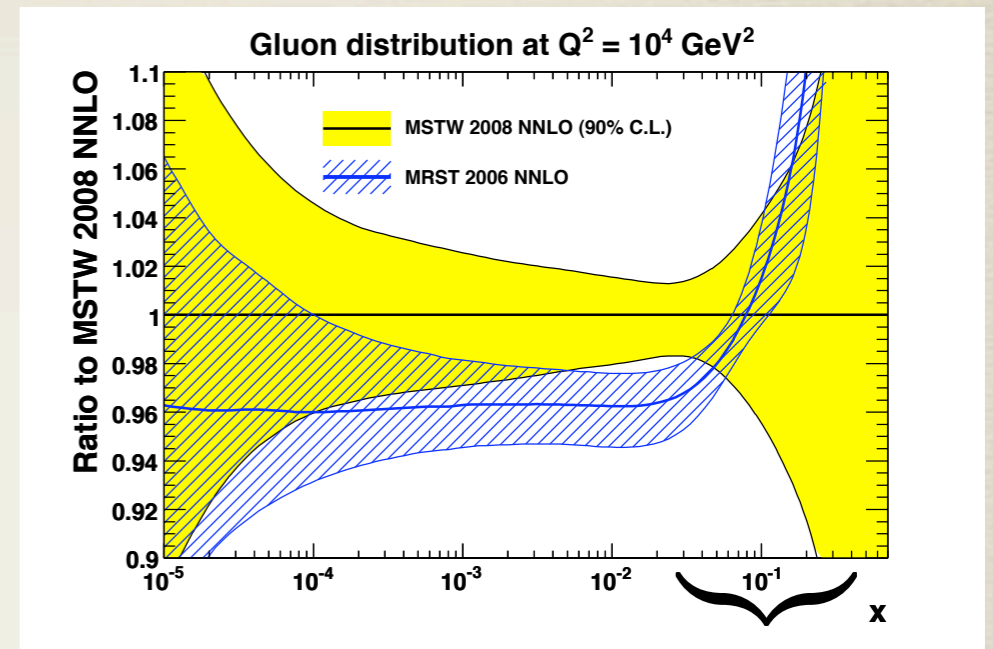
$M_H = 170$  GeV:

MRST 2001	MRST 2004	MRST 2006	MSTW 2008
0.3833	0.3988	0.3943	0.3444

**$\sim 15\%$  decrease** in predicted cross section !

$\sigma \sim \alpha_s^3 \times f_g^2 \Rightarrow$  very sensitive to these changes!

- ⦿ Central value, but not increased error, accounted for in 2009 analysis
- ⦿ LHC: 25% increase in cross sections at  $M_H = 120$  GeV, 10% at 200 GeV after changing PDFs, bottom-quark description, EW effects

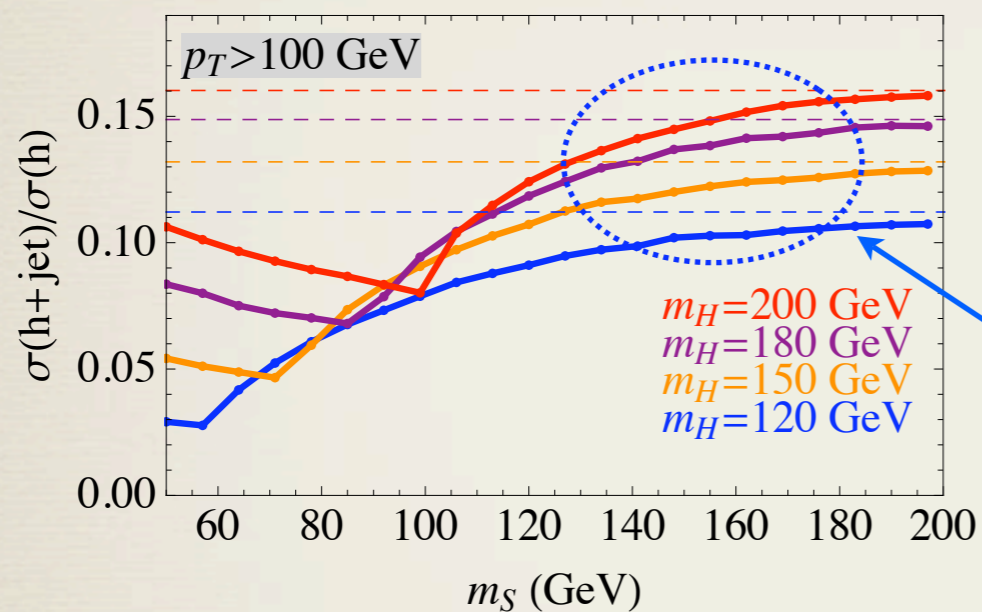


Changes in relevant region for Tevatron analyses

Keep in mind for LHC analyses...

# The Higgs $p_T$ spectrum

- Other surprises, perhaps in differential distributions?
- Many studies of  $p_T$  spectrum as probe of new physics



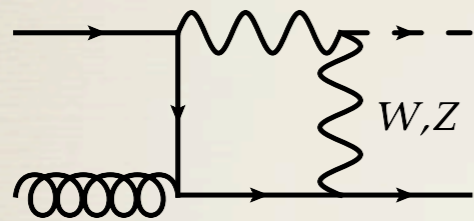
Color-octet scalar induced deviation in integrated ratio Arnesen, Rothstein, Zupan 0809.1429

Motivational region for precision study; new physics or QCD?

- Roughly 45% of Tevatron exclusion from 1,2 jet bins M. Herndon
- Some LHC analyses in  $\tau\tau$ ,  $\gamma\gamma$  select high  $p_T$  to remove background;  $p_T > 100$  GeV typical Abdullin et al. hep-ph/9805341; Mellado, Quayle, Wu hep-ph/0406095

# EW, quark-mass effects

- One possible problem:  $p_T^2/m_t^2$  effects from other EFT operators
- Effects missed previously that contribute to  $qg$ ,  $qq$  channels



• Same loop-order as previously included effects

- Not useful to compute, but think in terms of operators

$$\mathcal{O}_{EW} = \frac{H^\dagger D_\mu H}{v^2} \frac{\bar{q} G^{\mu\nu} \gamma_\nu q}{M_{W,Z}^2}$$

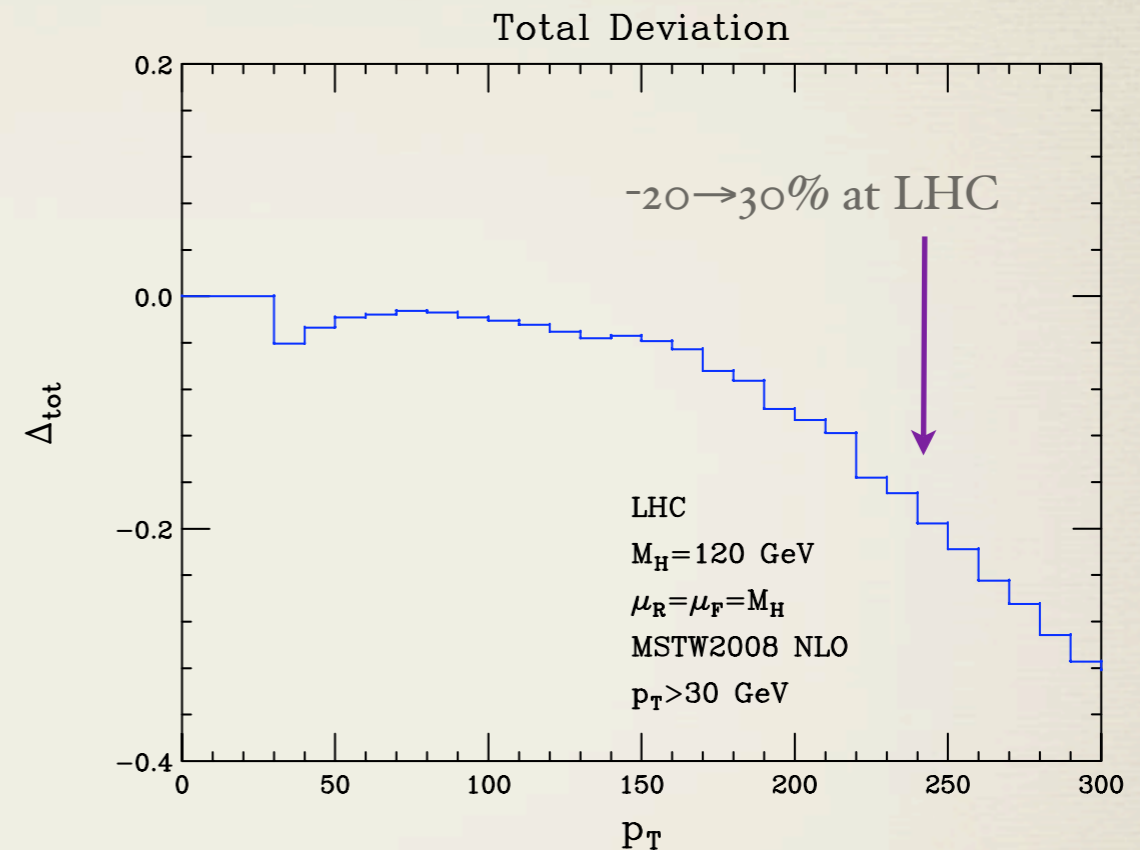
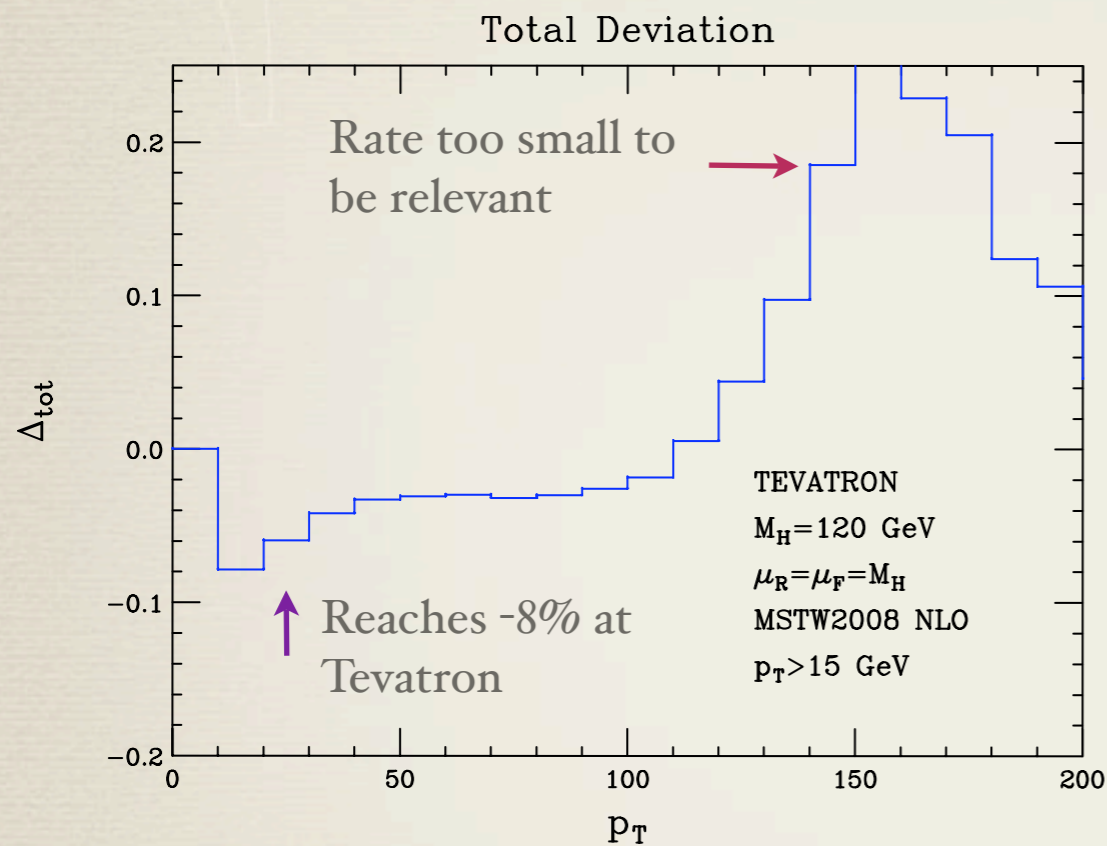
$$\Rightarrow \frac{p_H \cdot p_g \bar{u} \not{\epsilon}_g u - p_H \cdot \epsilon_g \bar{u} \not{p}_g u}{v M_{W,Z}^2}$$

• Vanishes for  $p_g \sim p_1, p_2 \Rightarrow$  hard  $p_T$  spectrum

• Interferes, destructively, with EFT contributions

# Numerical results

- Both  $W/Z$  and  $p_T^2/m_t^2$  act destructively to reduce EFT prediction



- All effects being included in updated analysis code FEHiP  
Anastasiou, Boughezal, Bucherer, FP, Stoeckli, in progress



# Low $p_T$ Higgs production

- Other searches restrict the  $p_T$  to low values

- LHC  $H \rightarrow WW$  jet veto:  $p_{TJ} < 20$  GeV Dittmar, Dreiner hep-ph/9608317

LHC 14 TeV		Accepted event fraction		
reaction $pp \rightarrow X$	$\sigma \times BR^2$ [pb]	cut 1-3	cut 4-6	cut 7 <span style="color: blue;">→ jet veto</span>
$pp \rightarrow H \rightarrow W^+W^-$ ( $m_H = 170$ GeV)	1.24	0.21	0.18	0.080
$pp \rightarrow W^+W^-$	7.4	0.14	0.055	0.039
$pp \rightarrow t\bar{t}$ ( $m_t = 175$ GeV)	62.0	0.17	0.070	0.001
$pp \rightarrow Wtb$ ( $m_t = 175$ GeV)	$\approx 6$	0.17	0.092	0.013
$pp \rightarrow ZW \rightarrow \ell^+\ell^-\ell\nu$	0.86	0.23	0.054	0.026
$pp \rightarrow ZZ \rightarrow 4\text{-leptons}$	1.05	0.13	0.016	0.007

- Leading order prediction for  $p_T$  spectrum:

$$\frac{d\sigma}{dp_T^2} \sim \sigma_0 \frac{C_A \alpha_s}{\pi} \frac{1}{p_T^2} \ln \frac{s}{p_T^2} \quad \Rightarrow \text{must resum to all orders}$$

# Impact-parameter formalism

- Classic analysis of Collins, Soper, Sterman NPB250 199 (1985)

$$\begin{aligned}
 \frac{d\sigma}{dQ^2 dy dQ_T^2} &= \frac{4\pi^2\alpha^2}{9Q^2s} (2\pi)^{-2} \int d^2b e^{iQ_T \cdot b} \sum_j e_j^2 \\
 &\times \sum_a \int_{x_A}^1 \frac{d\xi_A}{\xi_A} f_{a/A}(\xi_A; 1/b) \sum_b \int_{x_B}^1 \frac{d\xi_B}{\xi_B} f_{b/B}(\xi_B; 1/b) \\
 &\times \exp\left\{ - \int_{1/b^2}^{Q^2} \frac{d\bar{\mu}^2}{\bar{\mu}^2} \left[ \ln\left(\frac{Q^2}{\bar{\mu}^2}\right) A(g(\bar{\mu})) + B(g(\bar{\mu})) \right] \right\} \\
 &\times C_{ja}\left(\frac{x_A}{\xi_A}; g(1/b)\right) C_{jb}\left(\frac{x_B}{\xi_B}; g(1/b)\right) \\
 &+ \frac{4\pi^2\alpha^2}{9Q^2s} Y(Q_T; Q, x_A, x_B).
 \end{aligned}$$

- b-space arises from treatment of momentum conservation

$$\begin{aligned}
 \delta^{(2)}\left(\vec{p}_T - \sum_i \vec{k}_{Ti}\right) \\
 = \int \frac{d^2b}{(2\pi)^2} e^{ib \cdot \vec{p}_T} \prod_j e^{-i\vec{b} \cdot \vec{k}_{Tj}}
 \end{aligned}$$

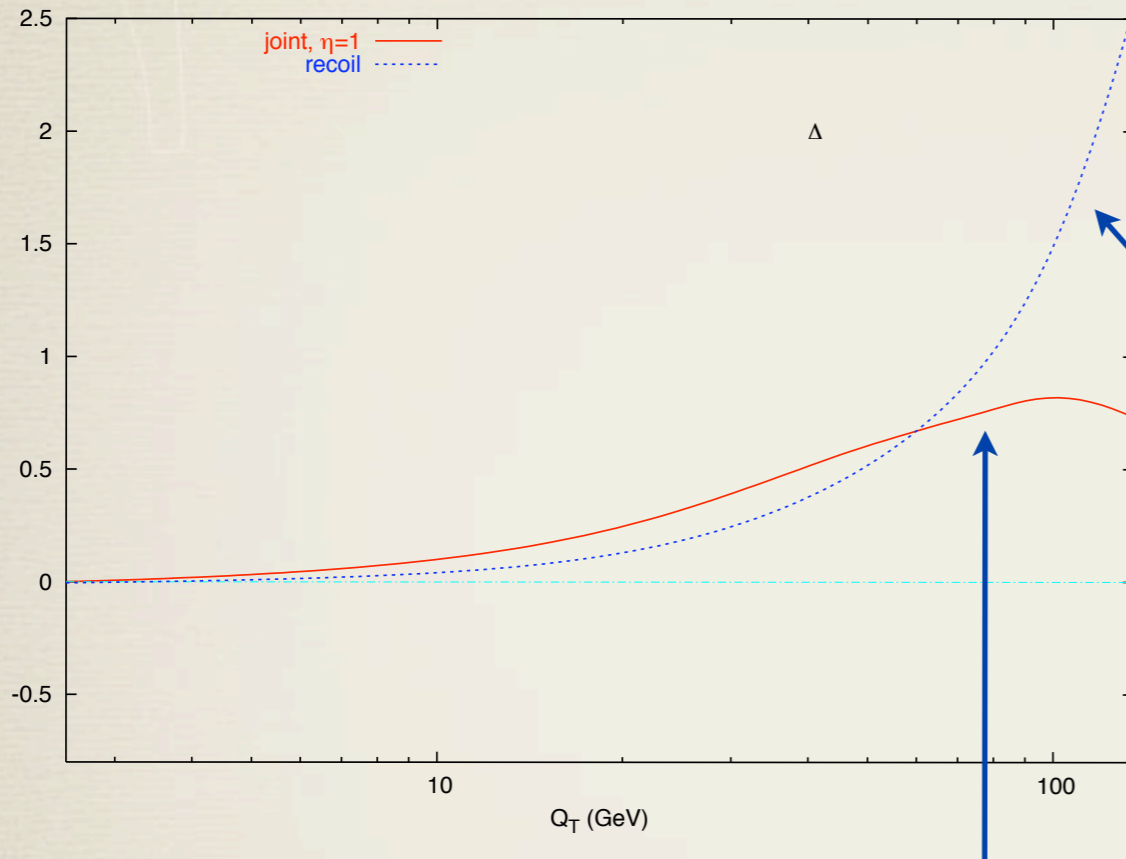
- Sudakov form factor: A, B contain  $\alpha_s(\bar{\mu}^2)$

- As  $b \rightarrow \infty$ , hit Landau pole of QCD coupling; non-perturbative physics enters (why for  $p_T > \Lambda_{\text{QCD}}$ ?)
- A typical approach; others exist (W is the integrand):

$$W(b) \rightarrow W(b_*) \exp\left[-b^2 g_1 - b^2 g_2 \ln \frac{b_{\text{max}} m_D}{2}\right], \quad b_* = \frac{b}{\sqrt{1 + b^2/b_{\text{max}}^2}}$$

# More on b-space resummation

- Other issues occur, in the combination with fixed order



$$\Delta \equiv \left[ \frac{d\sigma^{\text{fixed}(1)}}{dQ_T} - \frac{d\sigma^{\text{exp}(1)}}{dQ_T} \right] / \frac{d\sigma^{\text{fixed}(1)}}{dQ_T}$$

Large deviation of expanded b-space result from fixed-order complicates combining phase-space regions

“Joint resummation” Kulesza, Sterman, Vogelsang hep-ph/0309264

- Claim: more convenient framework to formulate low- $p_T$  resummation is soft-collinear effective theory

# Overview of SCET

- Lightning overview of SCET: split QCD gluon field into several fields with definite momentum scalings ( $\eta \sim p_T/M_H$ ) and their own gauge transformations

- $A = A_H + A_{c1} + A_{c2} + A_{us}$

- $A_H \sim (p^-, p^+, p_T) \sim M_H(1, 1, 1)$  “hard”

- $A_{c1} \sim M_H(\eta^2, 1, \eta)$  “collinear”

- $A_{c2} \sim M_H(1, \eta^2, \eta)$  “collinear”

- $A_{us} \sim M_H(\eta^2, \eta^2, \eta^2)$  “ultrasoft”

- ☑ Integrate out hard modes, match to collinear-invariant operators
- ☑ At leading-power, can decouple u-soft and collinear gluons

$$A_c = Y \circledast A_c^{(0)} Y^\dagger$$

$$Y_c = P \exp \left\{ ig \int_{-\infty}^x ds n \cdot A_{us}(sn) \right\}$$

$A^{(0)}$  has no u-soft couplings;  
factorize matrix element into u-  
soft and collinear components

# Low $p_T$ in SCET

- Sequence of effective field theories:  $\text{QCD} \rightarrow \text{SCET}_{p_T} \rightarrow \text{SCET}_{\Lambda\text{QCD}}$
- Structure after matching to  $\text{SCET}_{p_T}$ :

$$\frac{d^2\sigma}{dp_T dY} \sim \int dx_1 dx_2 |C(x_1, x_2; \mu_Q, \mu_T)|^2 \int dk_n^+ dk_{\bar{n}}^- d^2 k_n^\perp d^2 k_{\bar{n}}^\perp$$

$$\times J_n^{\alpha\beta}(x_1, k_n^+, k_n^\perp, \mu_T) J_{\bar{n}\alpha\beta}(x_2, k_{\bar{n}}^-, k_{\bar{n}}^\perp, \mu_T) S(x_1, x_2, k_{n\bar{n}}, \mu_T)$$

$$|C|^2 \sim \exp \left\{ - \int_{\mu_T^2}^{\mu_Q^2} \frac{d\bar{\mu}^2}{\bar{\mu}^2} \gamma_{\mathcal{O}}(\bar{\mu}^2) \right\}$$

- ☑ Lower limit of evolution set by matching scale  $\mu_T \sim p_T$
- ☑ Exact momentum conservation still implemented
- ☑ Non-perturbative effects in power-suppressed operators
- ☑ Collinear gluon-jet functions valid for other LHC processes; some interesting differences from previous SCET studies...

# Conclusions

- Intricate and large quantum effects on Higgs production
- Effective theory for gluon-fusion valid over a larger range than naively expected, excellent framework for pheno studies
- Combination of 3-loop light-quark terms, PDFs have significant effect on Tevatron exclusion limits
- Previously neglected high- $p_T$  effects calculated
- All result being implemented in up-to-date analysis code FEHiP
- Framework for low- $p_T$  resummation in SCET that should make fixed-order matching more convenient