

Recent progress in N=2 4d field theory

DG, G. Moore, A. Neitzke: [arXiv:0907.3987](https://arxiv.org/abs/0907.3987)

DG: [arXiv:0904.2715](https://arxiv.org/abs/0904.2715)

L.F.Alday, DG, Y.Tachikawa: [arXiv:0906.3219](https://arxiv.org/abs/0906.3219)

L.F.Alday, DG, S.Gukov, Y.Tachikawa, H.Verlinde:
[arXiv:0909.0945](https://arxiv.org/abs/0909.0945)

Non-rational 2d CFT
compute protected quantities
in 4d $N=2$ gauge theory

A strange relation

Pick a 2d CFT correlation function

- Pick a conformal Toda theory
 - ADE classification, Liouville theory is A_1
- Pick a Riemann surface
- Place vertex operators at punctures
- Add twist lines for outer automorphisms

Each choice of correlation function
selects a specific N=2 4d field theory

What does it compute?

4d N=2 gauge theory

- S^4 partition function
- Instanton partition function(s)
- Surface operators
- 't Hooft-Wilson loops

2d Toda CFT

- Correlation function
- Conformal blocks
- Degenerate field insertions
- Loop operators

Liouville theory examples

Sphere four point function

- SU(2) gauge theory with $N_f=4$

Sphere n point function

- Linear quiver of SU(2) gauge groups

Torus one point function

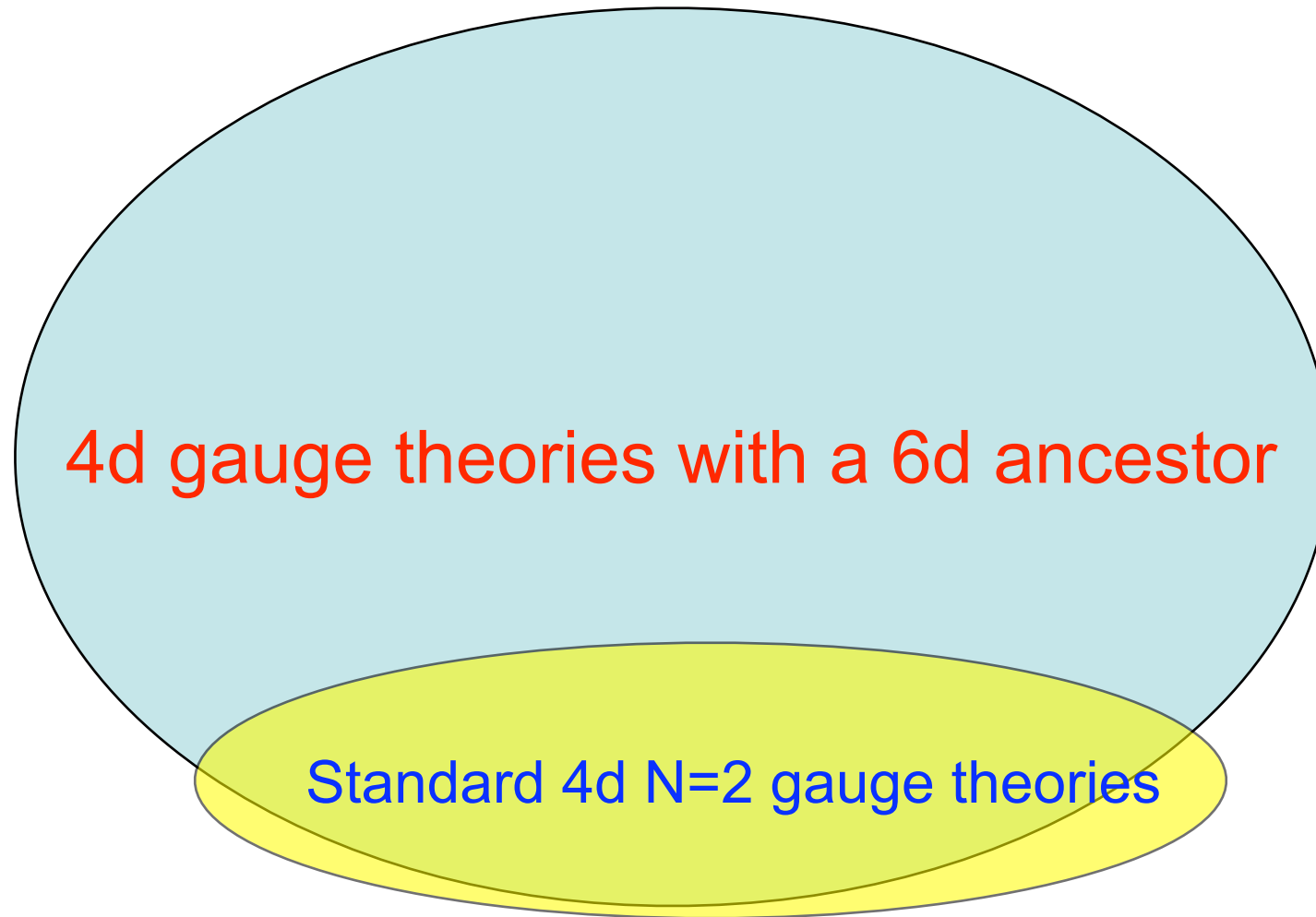
- SU(2) gauge theory coupled to one adjoint

Which $N=2$ gauge theories?

4d gauge theories with a “6d ancestor”

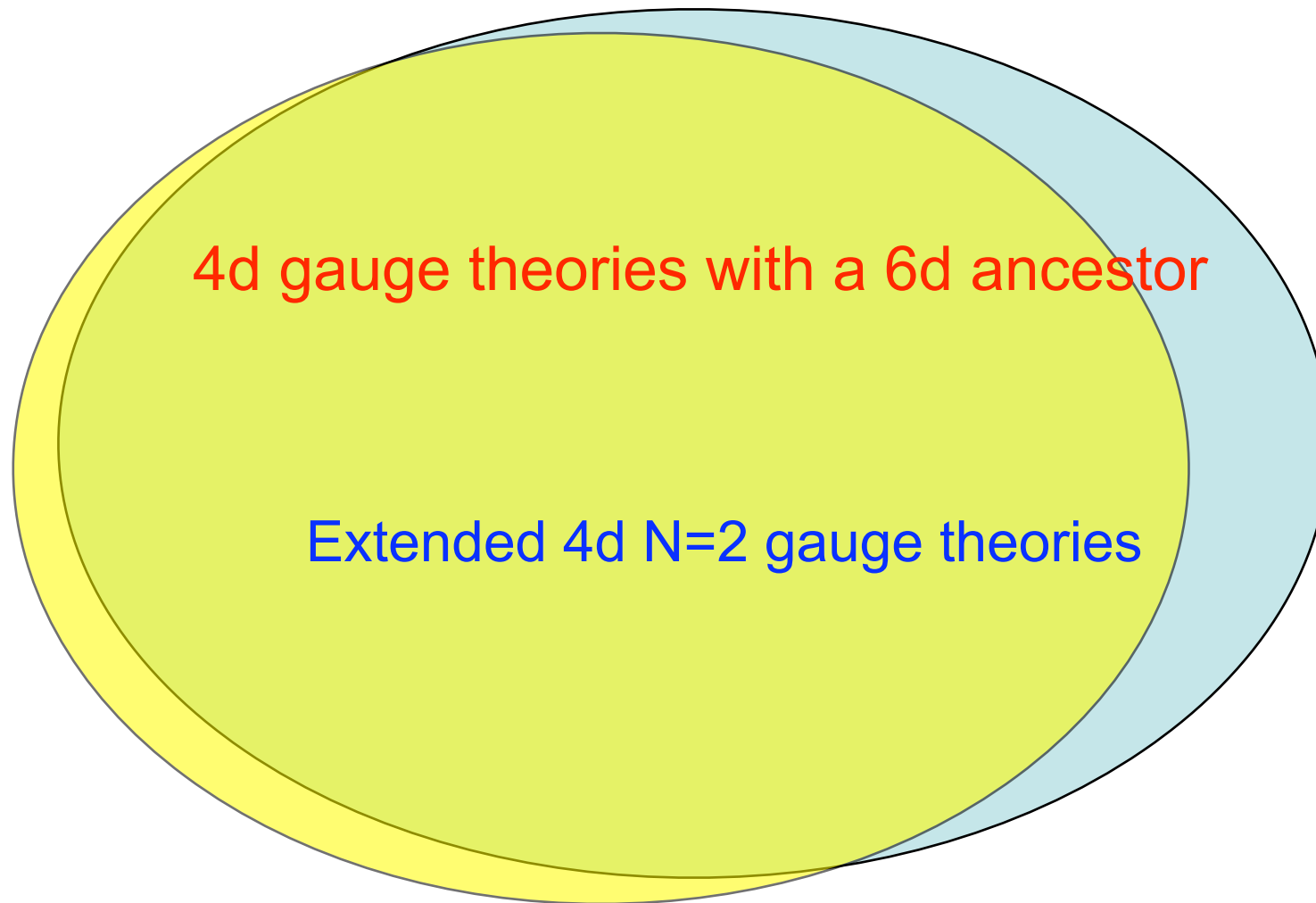
- Twisted compactification of 6d $(2,0)$ ADE SCFT
 - Pick a Riemann surface C
 - Place defects at points in C
 - Add twist lines if appropriate
- Flows in IR to a 4d theory with $N=2$ SUSY

Which N=2 gauge theories?



Do all N=2 theories have a 6d ancestor?

Which N=2 gauge theories?



A step back

Seiberg and Witten solution

- $N=2$ $SU(2)$ gauge theory
 - Exact low energy massless Lagrangian
 - Lagrangian described through SW curve and SW differential
 - Exact spectrum of massive dyonic BPS particles
- SW curve/differential for many more theories is available
 - No systematic field theory approach
 - M-theory, Type IIB engineering!
 - BPS spectrum is poorly understood

A step back

KS Wall Crossing Formula

- Recent mathematical progress
 - About “motivic Donaldson-Thomas invariants”
- Predicts behavior of BPS spectrum in $N=2$ field theories

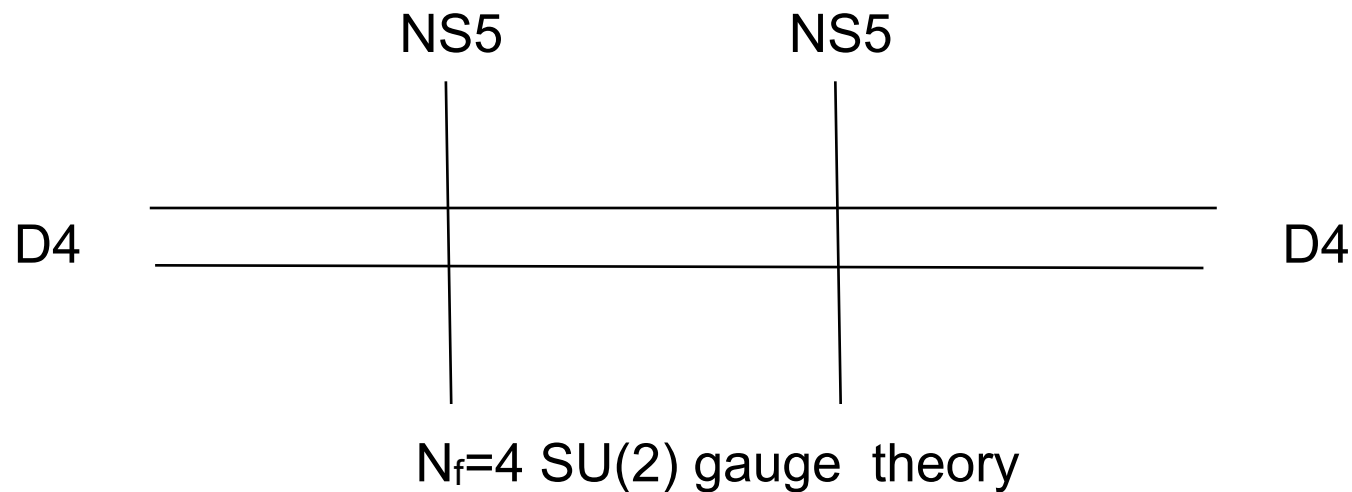
Physical explanation? GMN

- Compactify the $N=2$ 4d theory on a circle.
- BPS particles loops correct the metric on moduli space
- BPS spectrum jumps across walls of marginal stability
- Metric is continuous! Wall crossing formula follows

A step back

How to test this idea?

- We need both the BPS spectra and the 3d metric
 - Certain N=2 theories can be engineered in IIA

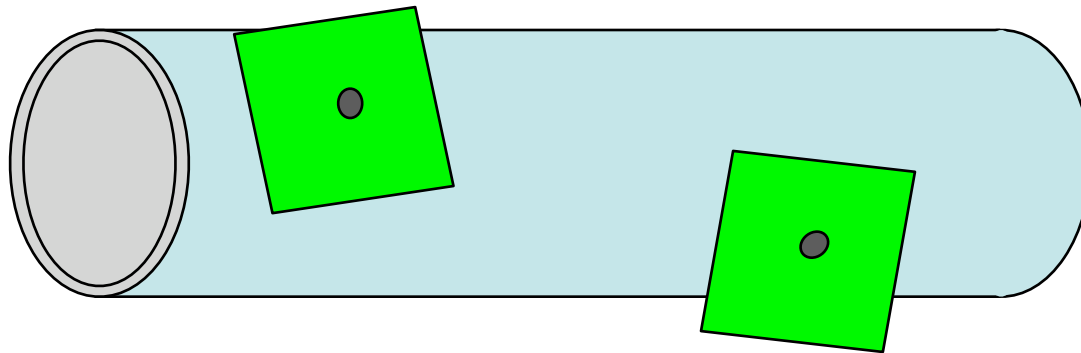


- SW curve/differential from lift to M-theory (Witten)
- BPS spectra and the 3d metric can also be derived!

A step back

IIA to M-theory

- M-theory lift: M5 branes wrapping a cylinder or torus C



- 4d gauge theory determined by choice of defects
- SW curves, differential, are constructed systematically

A broad generalization

Restrictive lifts

- Only special combinations of defects from IIA lift
- Only genus $g=0,1$ surfaces

Why not any surface, any defects combination?

- The M-theory setup is sensible
- We can apply the same analysis
 - Reasonable SW curve, BPS spectrum, etc.
 - Are these actual 4d field theories? Which ones?

A broad generalization

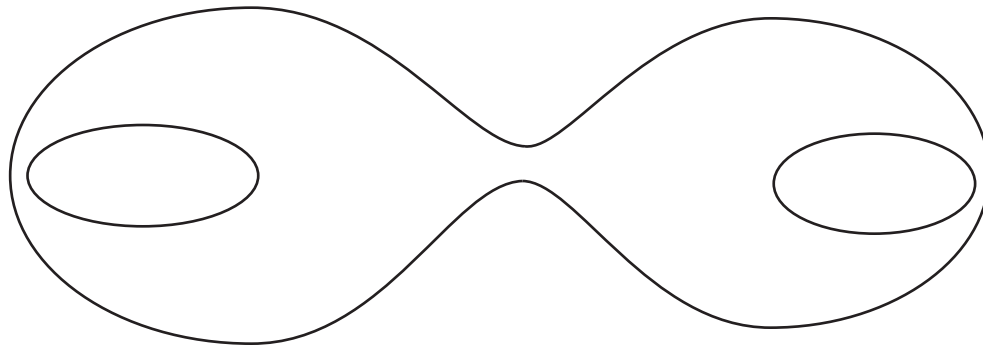
Divide and conquer

- Identify all UV gauge couplings
- Make them infinitely weak
- Read off weakly coupled gauge groups
- Read off matter theories

A broad generalization

Divide and conquer

- Gauge couplings control complex structure of C
- At weak coupling, C degenerates



“Pair of pants decomposition of C ”

- Matter theories associate to fragments of C
- Matter theories are often strongly interacting SCFTs!

Basic SCFTs in A_1 class

A simple example: $T_{g,n}[A_1]$

- Only one type of defect, carries $SU(2)$ flavor symmetry
- No twist lines
- Admit Lagrangian descriptions: matter theories are free.

$T_{0,3}$ the three punctured sphere

- Four free hypermultiplets
- $SU(2)_a \times SU(2)_b \times SU(2)_c$ flavor symmetry (3 punctures!)
- In components, (Q_{abc}, Q^*_{abc})
- A “pair of pants”

Sewing pants together

Connecting the pants

- To join pant leg a and pant leg b....
- Gauge diagonal $SU(2)$ in $SU(2)_a \times SU(2)_b$
- Instanton factor $q = \exp 2\pi i \tau$ is sewing parameter
- Weak coupling when tube is long

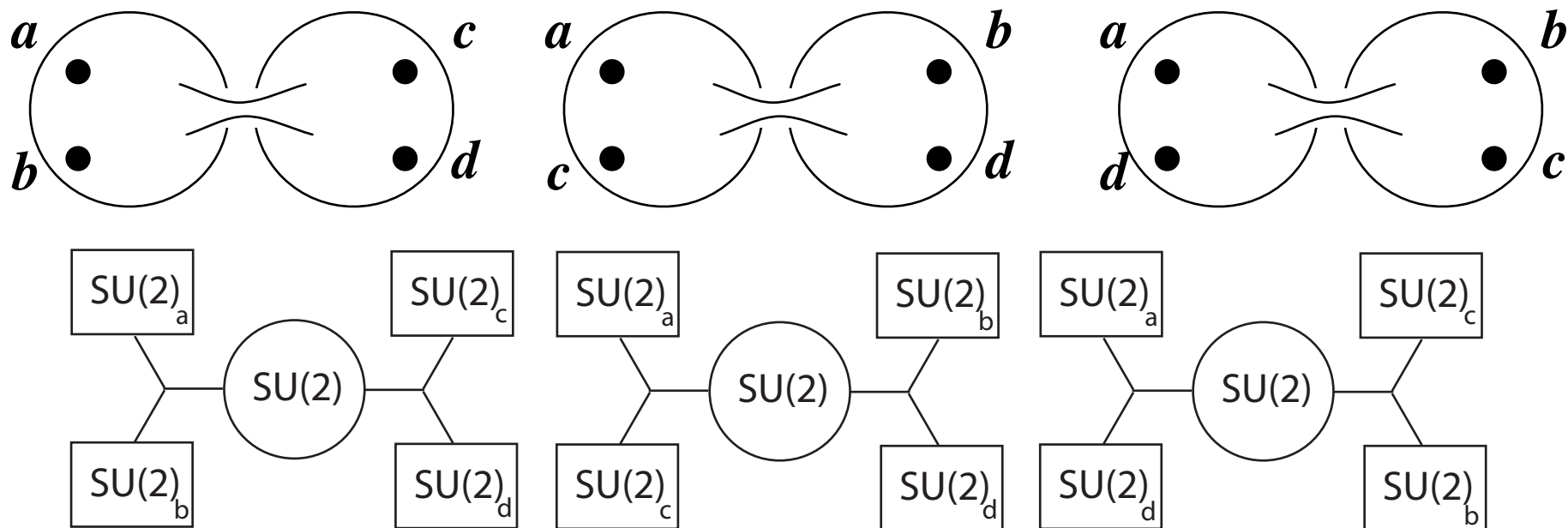
Lagrangians and sewing

- Any Riemann surface can be sewn from pair of pants
- For every sewing graph there is a distinct Lagrangian!
- Every theory $T_{g,n}$ has multiple Lagrangian descriptions!

Examples

$T_{0,4}$: sphere with four defects.

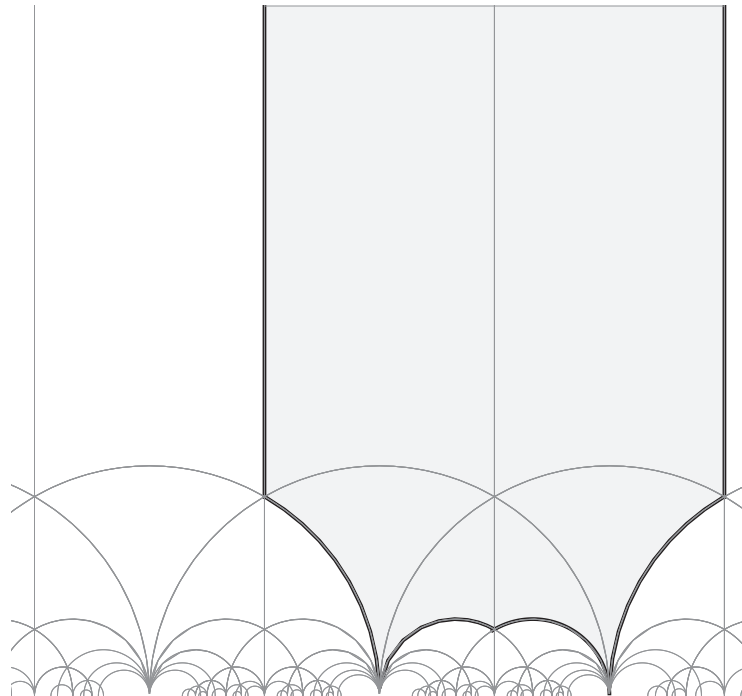
- Two pairs of pants
- One $SU(2)$ gauge group, 8 hypers: $SU(2) N_f=4$



S-duality in $SU(2)$ $N_f=4$

$SO(8)$ flavor symmetry

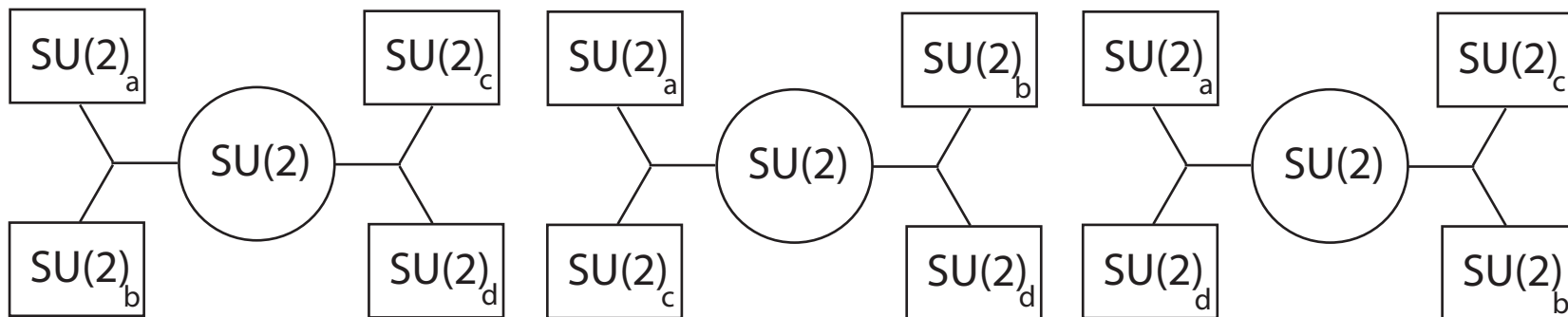
- $SL(2, \mathbb{Z})$ acts through triality on $SO(8)$
- Exchanges electrons in 8_v , monopoles in 8_s , dyons in 8_c



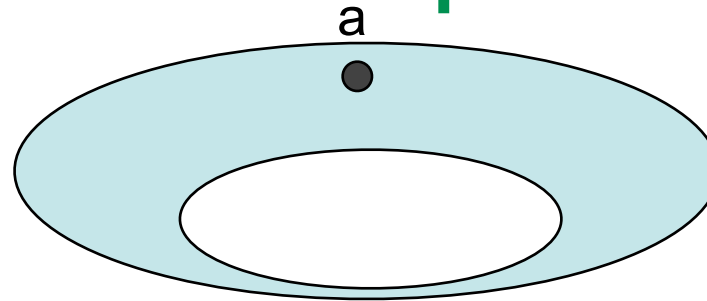
S-duality in $SU(2)$ $N_f=4$

Reformulating triality

- Consider subgroup $SO(4) \times SO(4)$ in $SO(8)$
- Rewrite it as $[SU(2)_a \times SU(2)_b] \times [SU(2)_c \times SU(2)_d]$
- $SL(2, \mathbb{Z})$ permutes (a, b, c, d)
 - $8_v = (2_a \times 2_b) + (2_c \times 2_d)$
 - $8_s = (2_a \times 2_c) + (2_b \times 2_d)$
 - $8_c = (2_a \times 2_d) + (2_c \times 2_b)$



Examples



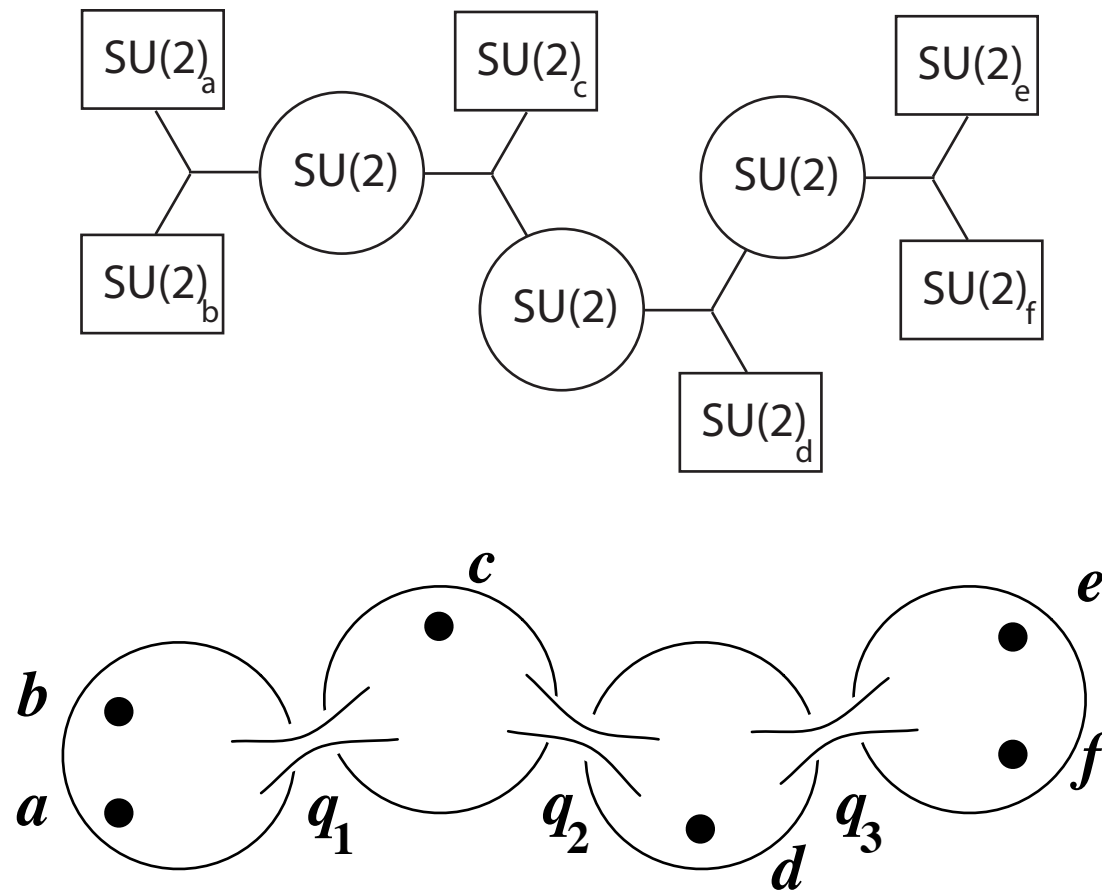
$T_{1,1}$: one-punctured torus

- One pair of pants, glue two legs
 - $(2_a, 2_b, 2_c)$ goes to $(2_a, 3) + (2_a, 1)$
 - One adjoint of $SU(2)$ gauge group and one free hyper

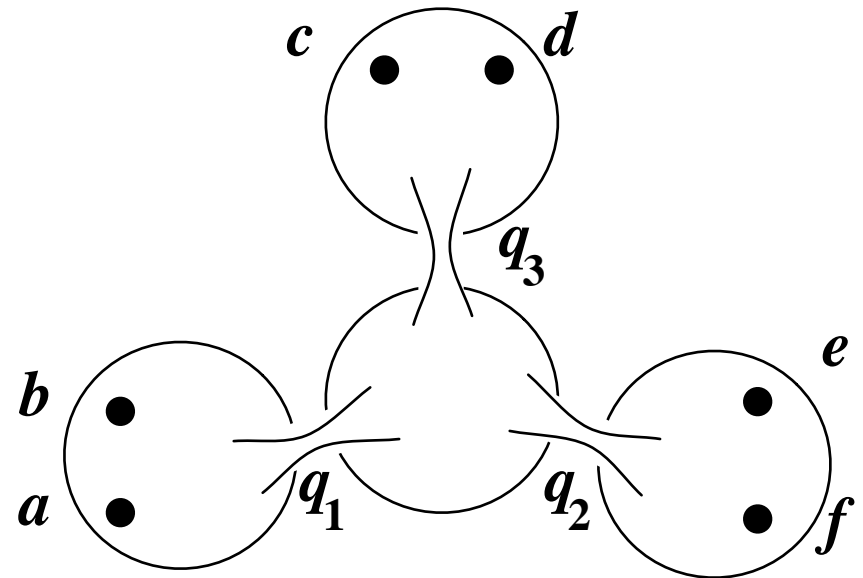
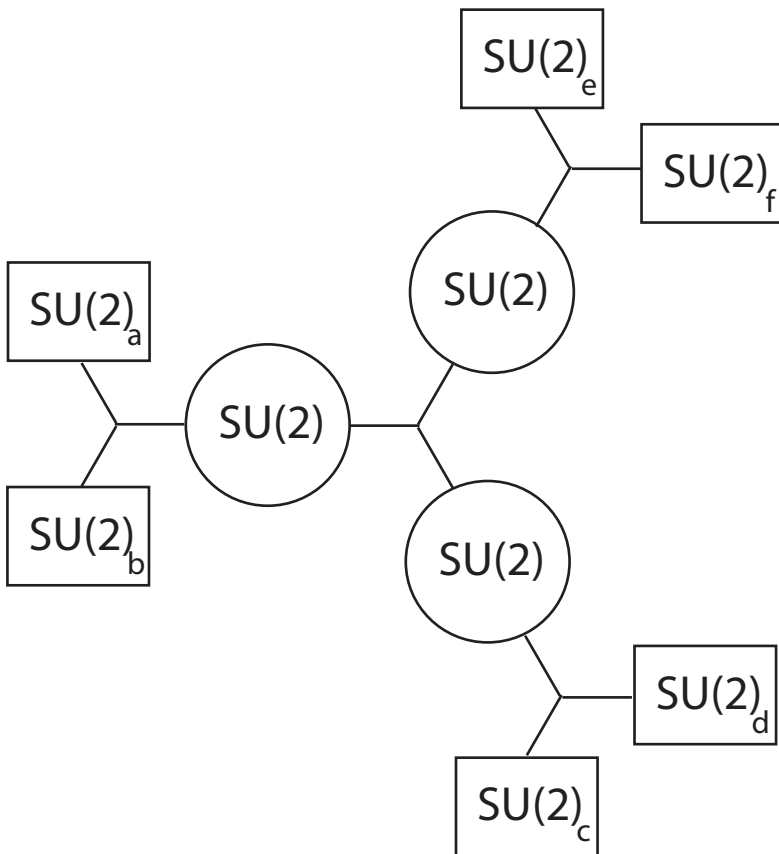
$N=2^*$ theory: mass deformed $N=4$ SYM

- Standard $SL(2, \mathbb{Z})$ S-duality group acting on torus

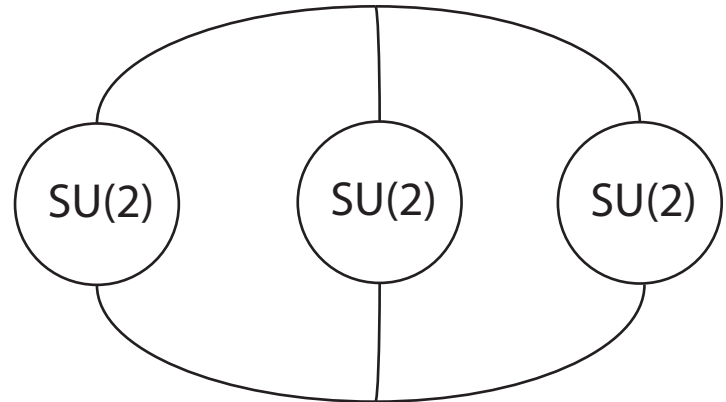
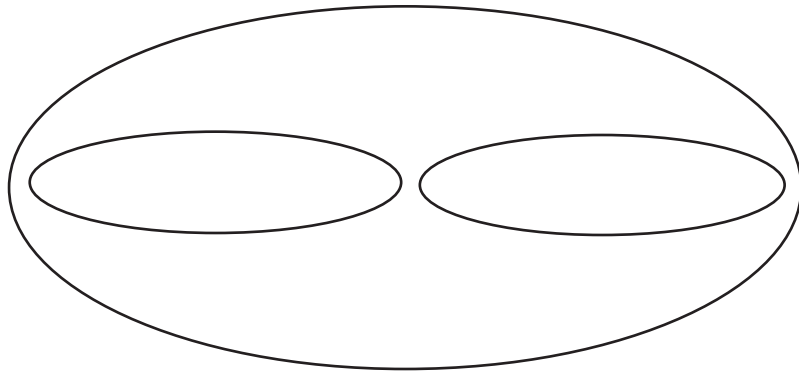
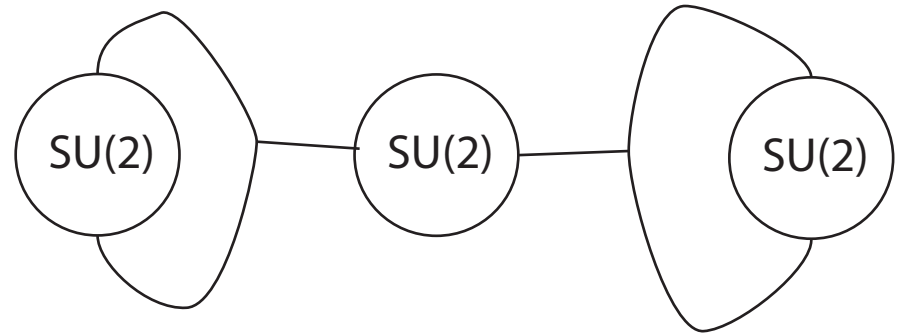
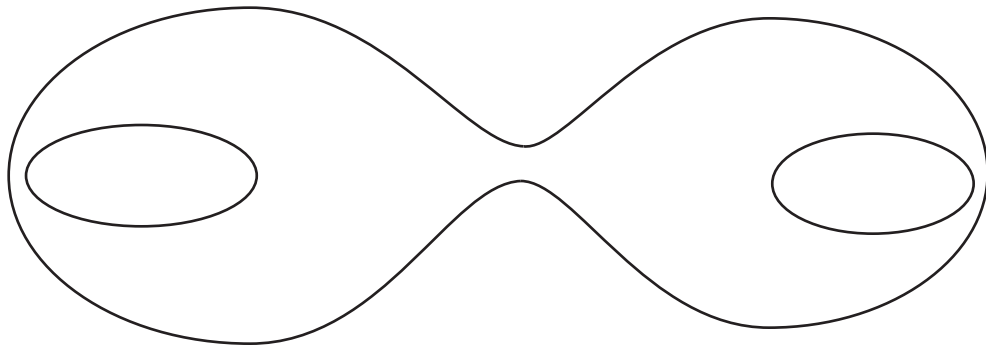
Sphere with six punctures



Sphere with six punctures



Genus two surface



S-dualities

A large S-duality group(oid) relates different sewings, and Lagrangians

- Same as “Moore-Seiberg groupoid” from 2d CFT lore
 - Generated by few basic operations
 - A-move: s-t channel duality in 4 point function ($N_f=4$ S-duality)
 - B-move: permute two legs
 - S-move: duality of one punctured torus ($N=2^*$ S-duality)
 - Crucial relations:
 - pentagon, hexagons at genus 0
 - genus 1 relations

On conformal blocks

Building blocks of 2d CFT correlation functions

- Correlation function computed by sewing
- Glue pair of pants with sum over complete set of states

$$\sum_v \sum_{i,j} |v, i\rangle B_{ij} \langle v, j| q^{L_0} \bar{q}^{\bar{L}_0}$$

- Fix primary fields v , only sum over descendants i, j

$$\mathcal{Z} = \sum_{v_a} \mathcal{F}[v_a](q_a) \bar{\mathcal{F}}[v_a](\bar{q}_a) \prod_n \langle v_a^n v_b^n v_c^n \rangle$$

On conformal blocks

Conformal blocks: a basis of objects which satisfy
Virasoro Ward identities

- Labeled by choice of primaries v_a in the channels
- For RCFT, finite set of primaries
 - finite dimensional space of conformal blocks
 - solutions of hypergeometric-like differential eqns.
- For non-rational theories, continuous labels
 - Hilbert space of conformal blocks
 - Labeled by conformal dimensions Δ
 - Better: Liouville momenta $\Delta = \alpha(Q-\alpha)$ $c = 1+6Q^2$

Examples

Liouville three point function (DOZZ)

$$\frac{\Upsilon(2\alpha_1)\Upsilon(2\alpha_2)\Upsilon(2\alpha_3)}{\Upsilon(\alpha_1 + \alpha_2 + \alpha_3 - Q)\Upsilon(\alpha_1 + \alpha_2 - \alpha_3)\Upsilon(\alpha_1 - \alpha_2 + \alpha_3)\Upsilon(-\alpha_1 + \alpha_2 + \alpha_3)}$$

Four point conformal blocks

$$\mathcal{F}[V_1, V_2, V_3, V_4](q) = \frac{\langle V_1 V_2 V_a \rangle \langle V_a V_3 V_4 \rangle}{\langle V_a | V_a \rangle} + \frac{\langle V_1 V_2 (L_{-1} V_a) \rangle \langle (L_{-1} V_a) V_3 V_4 \rangle}{\langle (L_{-1} V_a) | (L_{-1} V_a) \rangle} q + \dots$$

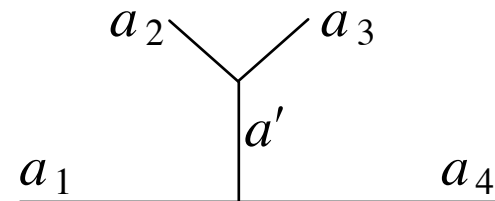
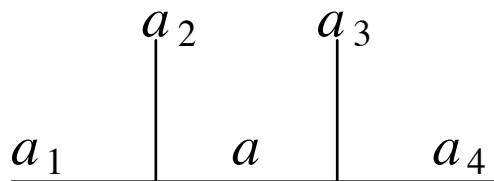
On conformal blocks

Are conformal blocks in different channels related?

- Action of Moore-Seiberg groupoid
- A,B,S moves implemented by “braiding, fusion matrices”
 - For non-rational CFT, really intricate integration kernels
 - Act on primary labels
 - Do NOT depend on moduli of surface!

On conformal blocks

Example: four points on sphere



$$\mathcal{F}_s[a_1, a_2, a_3, a_4](q, a) = \int da' F[a_1, a_2, a_3, a_4](a, a') \mathcal{F}_t[a_3, a_2, a_1, a_4](1 - q, a')$$

First dictionary entry

Two occurrences of Moore-Seiberg. Any relation?

- We need a gauge theory quantity which is holomorphic
 - Ideally, a power series in q
 - sewing parameter, but also instanton factor

Nekrasov instanton partition function

- A deformation of flat space partition function
- Two deformation parameters $\varepsilon_1, \varepsilon_2$
- Precise definition would bring too far
- Coefficient of q^k computed from k -instanton moduli space

Nekrasov partition function and conformal blocks

Gauge theory partition function depends on a_i, m_j

- a_i is Coulomb branch parameter of i -th $SU(2)$ gauge group
- m_j is mass parameter for $SU(2)_j$

“Experimental” observation:

- Z_{inst} coincides with conformal blocks up to irrelevant factor!
- Dictionary:
 - $\varepsilon_1 = b \varepsilon$ $\varepsilon_2 = \varepsilon/b$ $c = 1 + 6(b+b^{-1})^2$
 - $\alpha = b/2 + b^{-1}/2 + a/\varepsilon$ Liouville momenta
 - $\Delta = \alpha(b + 1/b - \alpha)$

Explicit formulae: $SU(2)$ $N_f=4$

Z_{inst}

$$1 - q \left(\frac{(a - m_1 + \epsilon_1 + \epsilon_2)(a - m_2 + \epsilon_1 + \epsilon_2)(a - m_3 + \epsilon_1 + \epsilon_2)(a - m_4 + \epsilon_1 + \epsilon_2)}{2a\epsilon_1\epsilon_2(2a + \epsilon_1 + \epsilon_2)} + a \leftrightarrow -a \right)$$

Conformal block

$$1 + \frac{(\Delta - \Delta_1 + \Delta_2)(\Delta + \Delta_3 - \Delta_4)}{2\Delta} q$$

Relation

$$Z_{\text{inst}} = (1 - q)^\delta \mathcal{F}$$

A neat use of Z_{inst}

Pestun computation of S^4 partition function

- Careful localization procedure

$$\int a^2 da |Z_{tree} Z_{1-loop} Z_{inst}|^2$$

Does it have a 2d CFT interpretation?

- Z_{1-loop} and Vandermonde measure give DOZZs!
- $Z_{tree} Z_{inst}$ provide overall powers of q , conformal blocks
- $Z[S_4]$ is Liouville correlation function! ($b=1$)

S-duality invariance from crossing symmetry

Correlation functions are independent of sewing

- Then $Z[S_4]$ is also S-duality invariant
- Very hard to prove directly!

Can we extend this result?

- Pestun also inserted Wilson loops along equator

$$\int a^2 da \cos \pi a |Z_{tree} Z_{1-loop} Z_{inst}|^2$$

S-duality and line operators

Can we do an S-duality transformation?

- Answer involving integral kernels is not very useful
- Hard to simplify by hand
- We need a direct 2d CFT interpretation of Wilson loop
- A function of the Liouville momentum in a channel

How do you “measure” the primary in a channel?

- Erik Verlinde strategy:
- monodromies of degenerate insertions

Degenerate fields

Primaries of momentum $2\alpha = (1-n)b + (1-m)b^{-1}$

- Null vector at level nm in Verma module
- In correlation functions, satisfy deg. nm differential eqn.

Example

- $V_{1,1}$ is identity operator, $\partial V_{1,1} = 0$
- We'll look at $V_{2,1}$
 - $\partial^2 V_{2,1} + b^2 :T V_{2,1}: = 0$
 - Restricted fusion: $V_{2,1} V_\alpha = V_{\alpha+b/2} + V_{\alpha-b/2}$

Degenerate fields and Conformal blocks

Degenerate fields can be inserted in a channel

- Liouville momentum jumps by $\pm b/2$ from left to right

Degenerate fields can be moved between channels

- “Degenerate” fusion and braiding operators
- Act as 2×2 matrices on $\pm b/2$ choice
- Transport matrices for $\partial^2 V_{2,1} + b^2 :T V_{2,1}: = 0$
- Easy to compute and use! No integral kernels

E. Verlinde move

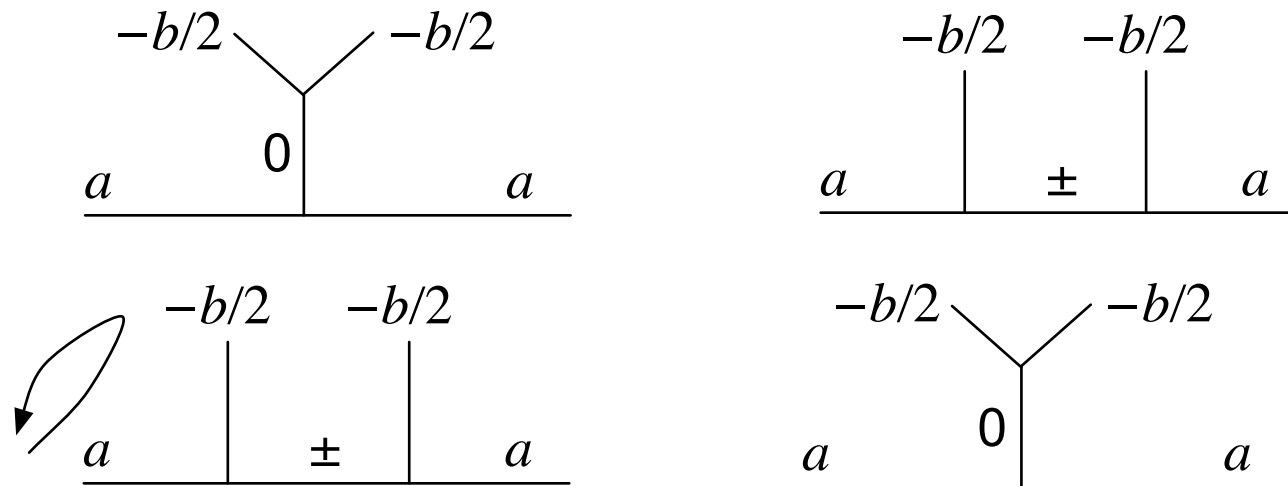
Insert identity in conformal block as fusion $V_{2,1} V_{2,1}$

Transport one $V_{2,1}$ along closed path p on C

Fuse back to identity

Defines a “loop operator” L_p

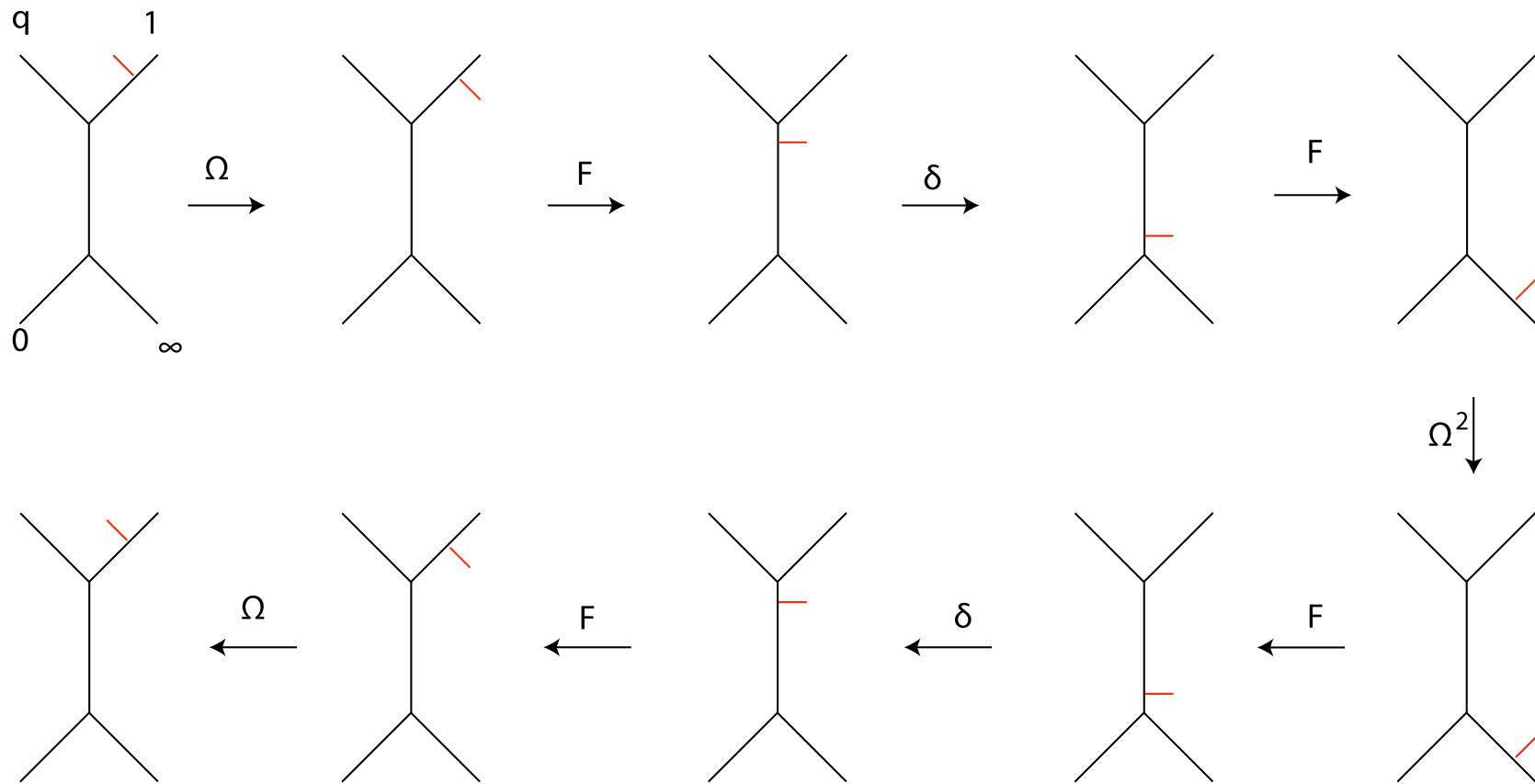
Wilson loops and L_p



Consider p around a leg

- No degenerate fusions, braiding gives a phase $\pm\pi b \alpha$
- Fusion to identity acts as a trace, gives $\cos \pi b \alpha$
- Wilson loop insertion! Z_{inst} goes to $\cos \pi b \alpha Z_{\text{inst}}$

Degenerate fields and Conformal blocks



Wilson loops and L_p

To compute S-dual,

- Pick a different pants decomposition
- Compute L_p on the same path
- More intricate sequence of elementary moves
 - Degenerate fusion matrices appear,
 - L_p involves shifts of α
- $L_p Z_{\text{inst}}(\mathbf{a}) = \sum c_n(\mathbf{a}) Z_{\text{inst}}(\mathbf{a} + n \mathbf{b}/2)$

Time is up

Degenerate insertions are surface operators

- A whole story there.....

Which loops did we compute in S-dual frame?

- Drukker: 'tHooft-Wilson loops labeled by paths p

Is the result sensible?

- A direct gauge theory computation would be neat

Time is up

More open questions

- Direct proof of 2d-4d correspondence
- Why 6d? Find pure 4d gauge theory interpretation
- Classification of $N=2$ theories?