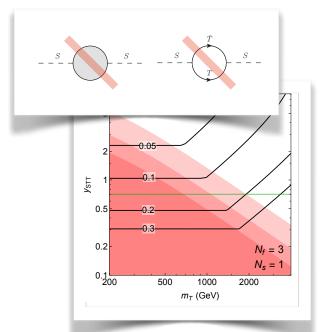
Discovering Neutral Naturalness

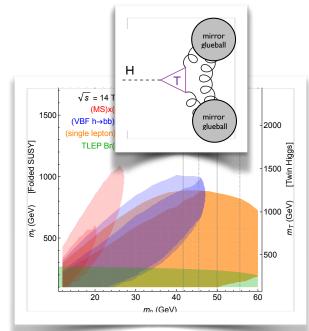
Theory Seminar NHETC Rutgers University

20. October 2015

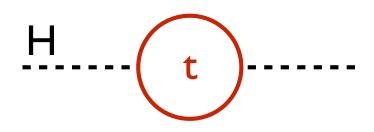
David Curtin University of Maryland

based on DC, Verhaaren 1506.06141 DC, Saraswat 1509.04284 Chacko, DC, Verhaaren, 1512.XXXXX

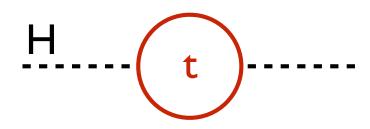




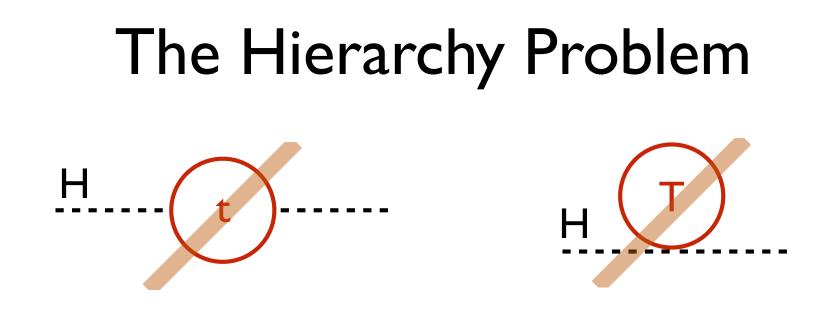
The Hierarchy Problem



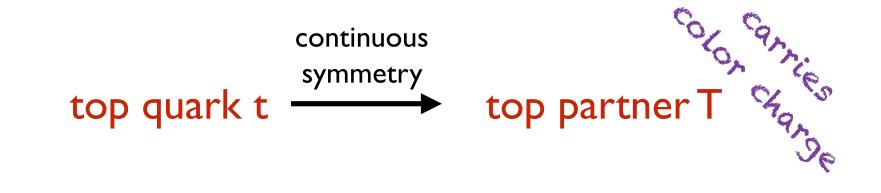
The Hierarchy Problem



... can be solved by top partners

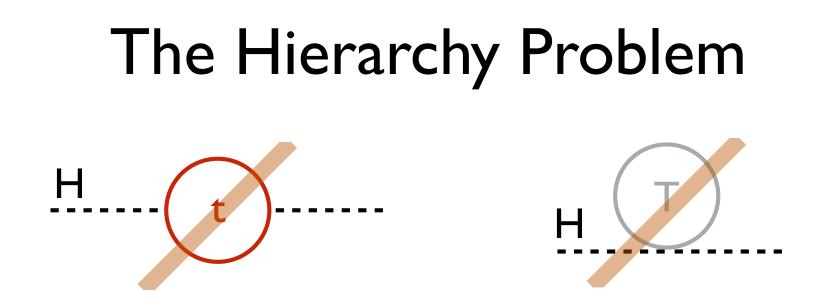


... can be solved by top partners

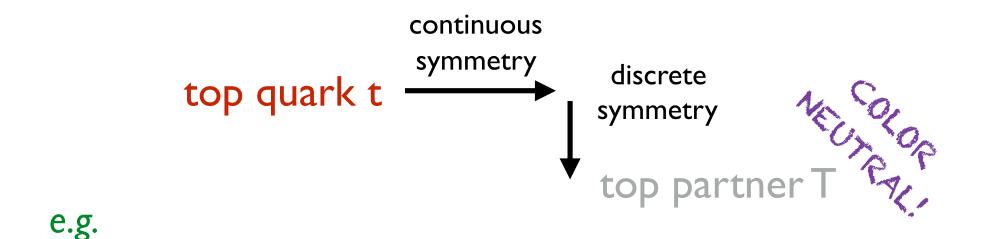


e.g.

Supersymmetry, modern composite Higgs models, etc...



The symmetry need not commute with SM color!

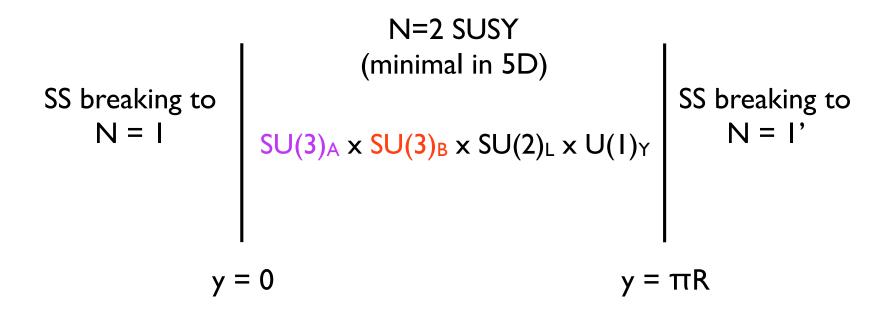


Folded SUSY (EW-charged stops), Twin Higgs (SM singlet T-partners)

hep-ph/0609152 Burdman, Chacko, Goh, Harnik

hep-ph/0506256 Chacko, Goh, Harnik

Theory Example: Folded SUSY



Boundary conditions break $A \leftrightarrow B$ symmetry and globally break N=2 to N=0 SUSY.

Normal MSSM EW sector.

SU(3) sectors: only zero modes are A-fermions, B-sfermions.

'Accidental supersymmetry' protects Higgs @ I-loop with EW charged top partners. hep-ph/0609152

hep-ph/0609152 Burdman, Chacko, Goh, Harnik

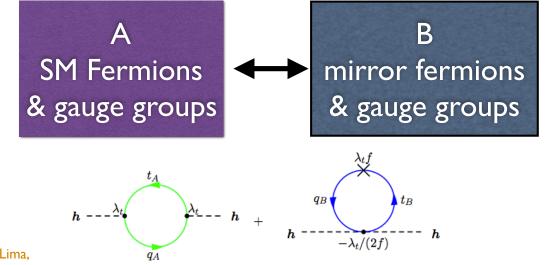
Theory Example: Twin Higgs

 $SM_A \times SM_B$ (mirror sector) particle content with Z_2 symmetry

Higgs sector: SU(4), broken by Gauge + Yukawa interactions to $SU(2)_A \propto SU(2)_B \propto Z_2$, which generate mass for goldstone boson.

$$\Delta V = \frac{3}{8\pi^2} \Lambda^2 \left(\lambda_A^2 H_A^{\dagger} H_A + \lambda_B^2 H_B^{\dagger} H_B \right) \qquad \qquad \lambda_A = \lambda_B \equiv \lambda + \Delta V = \frac{3\lambda^2}{8\pi^2} \Lambda^2 \left(H_A^{\dagger} H_A + H_B^{\dagger} H_B \right) = \frac{3\lambda^2}{8\pi^2} \Lambda^2 H^{\dagger} H_B$$

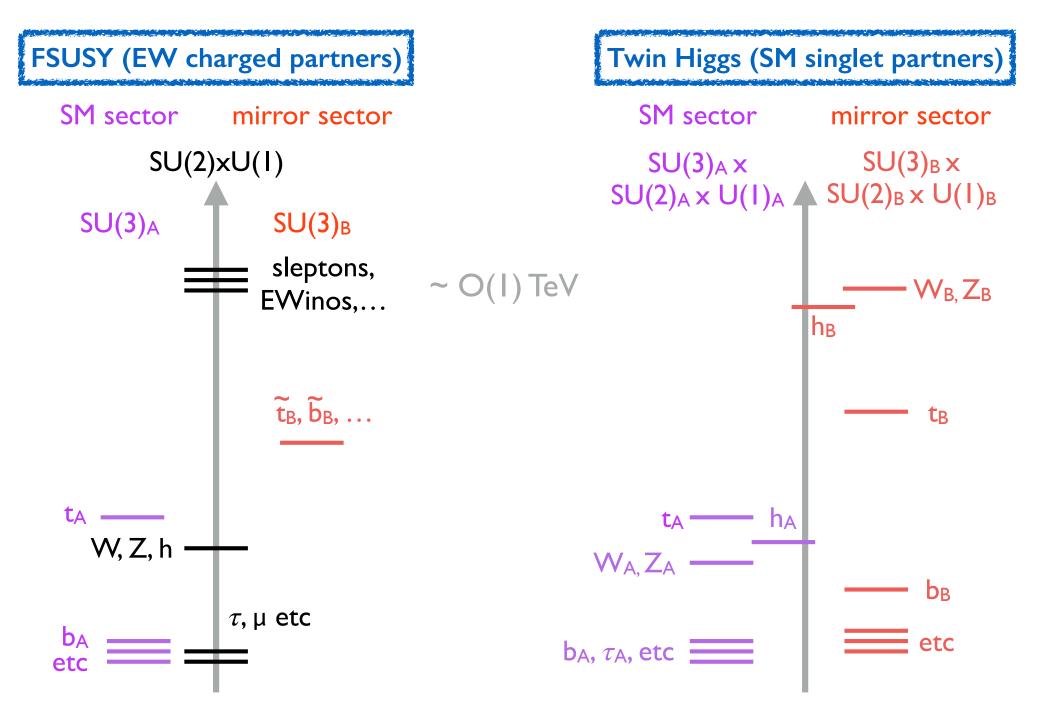
Z₂ symmetry of quadratically divergent contributions mimics full SU(4) symmetry, protects pNGB Higgs mass @ 1-loop.



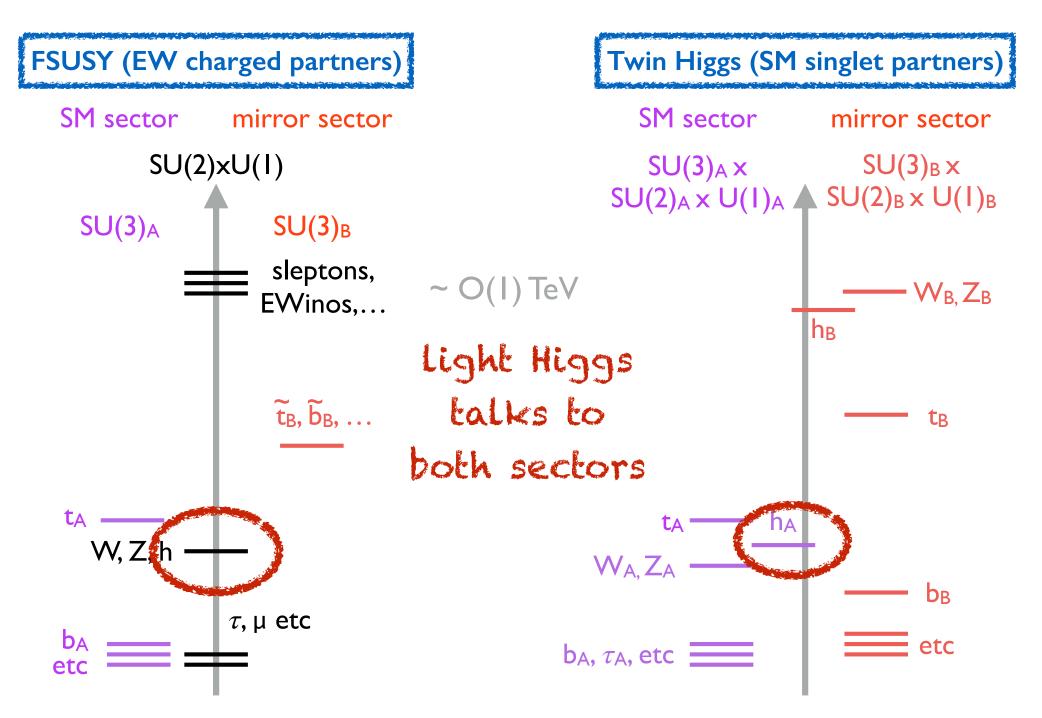
SM singlet top partners.

hep-ph/0506256 Chacko, Goh, Harnik 1411.3310 Burdman, Chacko, Harnik, de Lima, Verhaaren

Typical Low-Energy Spectra



Typical Low-Energy Spectra



Neutral Naturalness

Why would we think about this?

- The LHC is *great* at making colored particles, but so far no top partner discovery...
- 2. Want to examine naturalness as generally as possible: test the mechanism, not the model!

Neutral Naturalness generates radically different phenomenology from colored partners!

What are the most

important questions

right now?

What signals of Neutral Naturalness could we probe today?

. What signals of Neutral Naturalness could we probe today?

2. If the signatures are this malleable... will we be able to probe the general mechanisms underlying naturalness tomorrow?

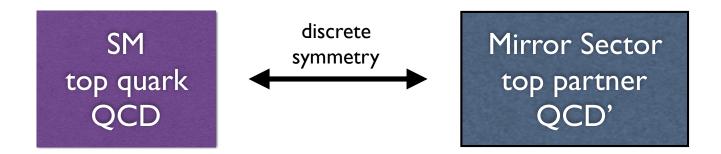
Probing Naturalness today: Signatures of Neutral Naturalness

at the LHC

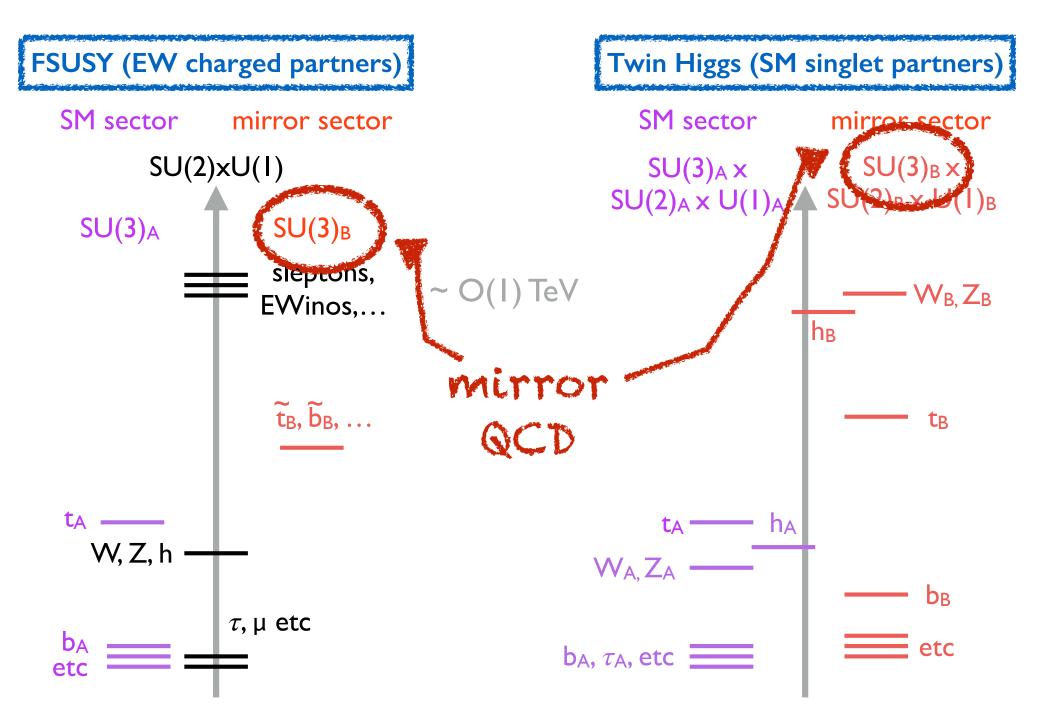
Hidden Valley Phenomenology

c.f. Strassler, Zurek '06 etc..

In theories of Neutral Naturalness, the partners in the mirror sector are usually charged under a copy of QCD



Typical Low-Energy Spectra

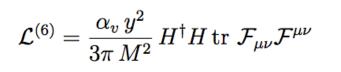


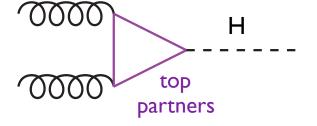
Hidden Valley Phenomenology

c.f. Strassler, Zurek '06 etc..

Mirror gluons talk to the Higgs via top partner loops!

(just like the top quark connects the Higgs to SM QCD)

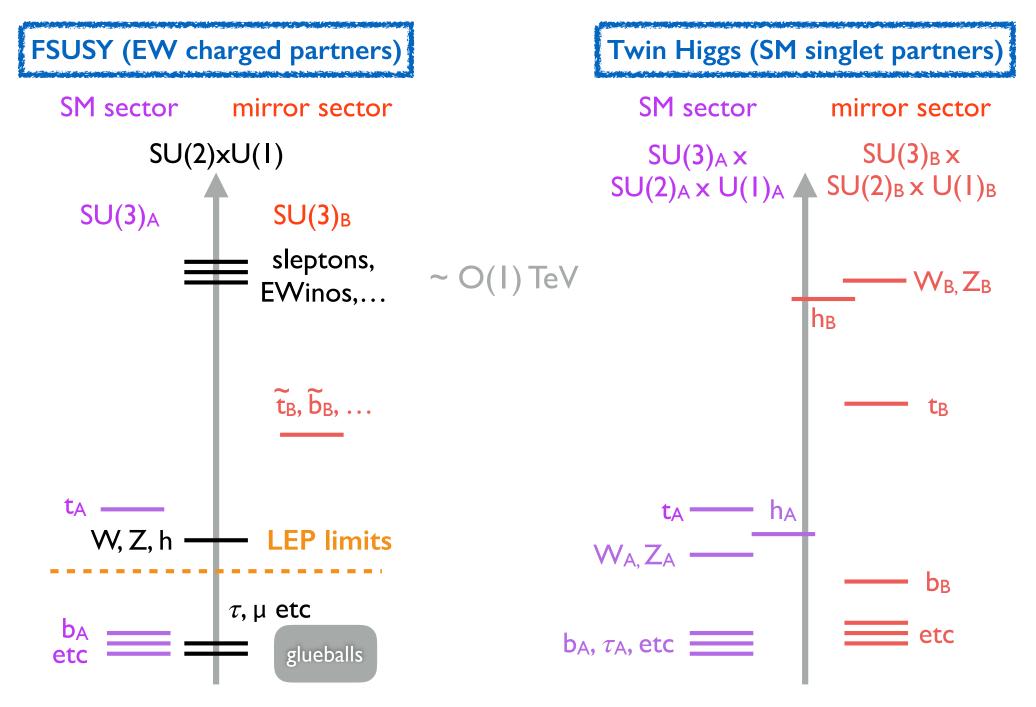


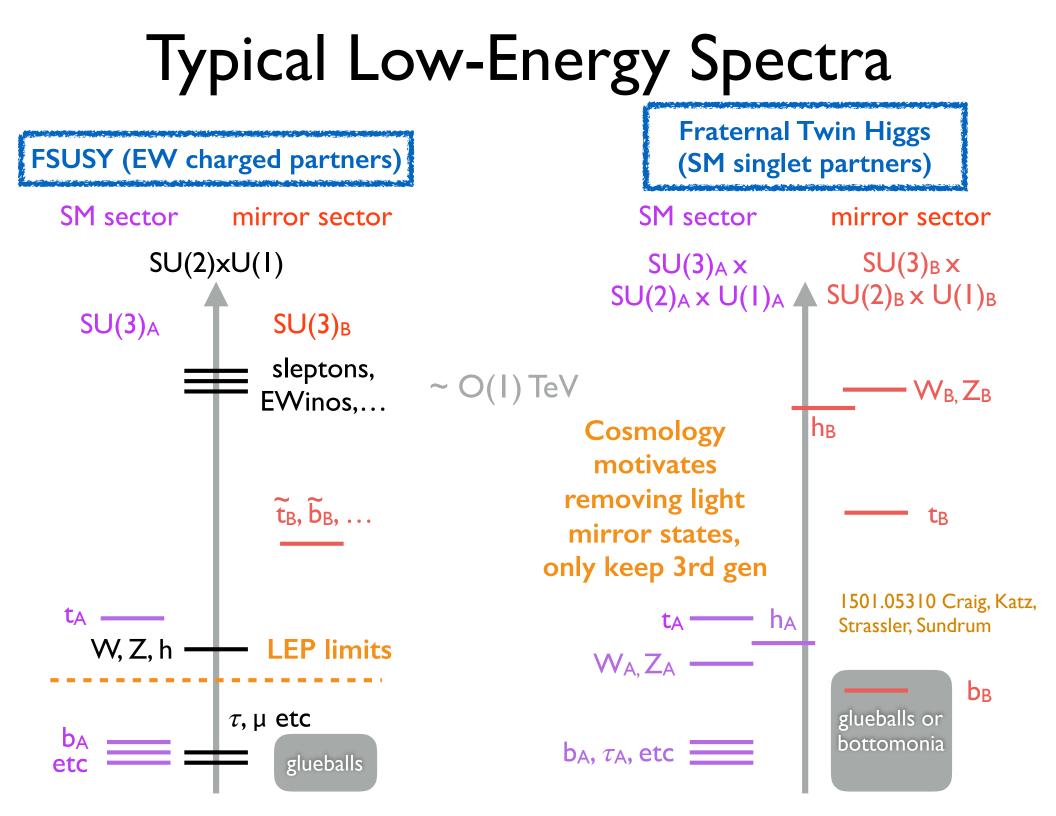


The mirror sector contains mirror hadrons.

Detailed consequences depend on the mirror spectrum: pions? quarkonia? glueballs?

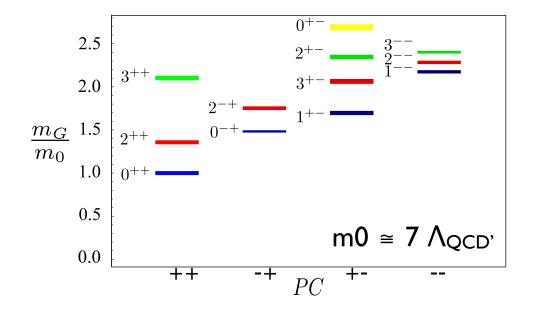
Typical Low-Energy Spectra





Mirror Glueballs

If the mirror sector has no light matter, the mirror QCD hadrons are glueballs.



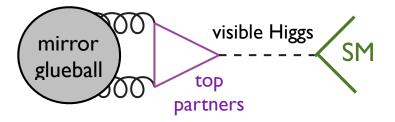
"Required" for EW charged top partners.

Possible (motivated by cosmology) for SM singlet top partners.

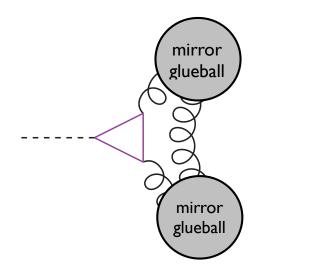
Glueball-Higgs Coupling

0903.0883 Juknevich, Melnikov, Strassler; 0911.5616 Juknevich

Glueballs mix with the Higgs via top partner loop: 0++ would eventually decay back to SM!



The Higgs could also decay to these glueballs: exotic Higgs decays with displaced vertices!



Key signature of uncolored naturalness!

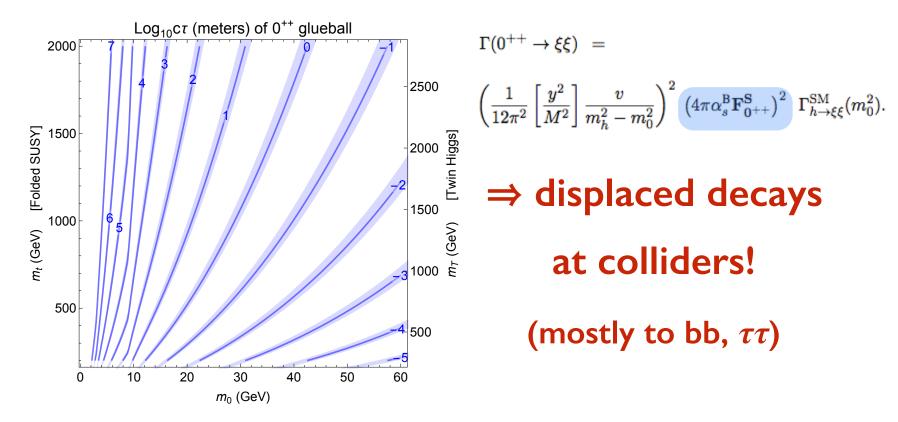
1501.05310 Craig, Katz, Strassler, Sundrum

Is this signature realized?

Mass: $m_0 \sim 7\Lambda_{QCD'} \sim 10 - 60 \text{ GeV}$ from RG arguments, but can move that around in Twin Higgs theories. DC,Verhaaren 1506.06141

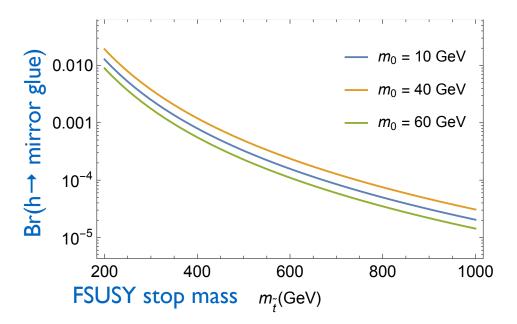
 \Rightarrow can be produced in exotic Higgs decays!

Lifetime of 0^{++} : $c\tau \sim \mu m$ - Ikm (using lattice results)



How many glueballs from Higgs decays?

Estimate inclusive mirror-glue production by rescaling SM Br($h \rightarrow gg$) by top partner loop and mirror α_s (also from RG arguments).



LHC 14 with 300fb-1 makes O(10 million) higgs bosons.

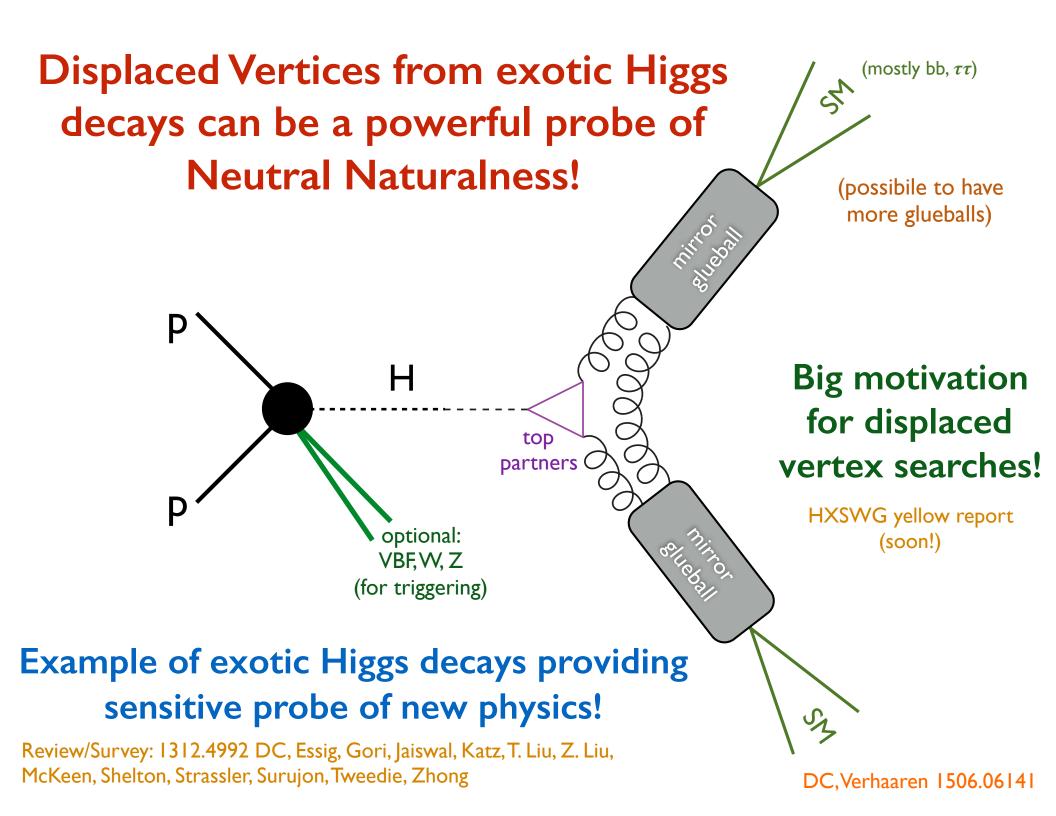
Could probe TeV-scale top partners if exotic Higgs decays conspicuous enough!

Conservatively estimate exclusive production of unstable 0++ glueball by parameterizing our ignorance about mirror hadronization:

$$Br(h \to 0^{++}0^{++}) = Br(h \to \text{mirror glue}) \cdot \kappa \cdot \sqrt{1 - \frac{4m_0^2}{m_h^2}}$$

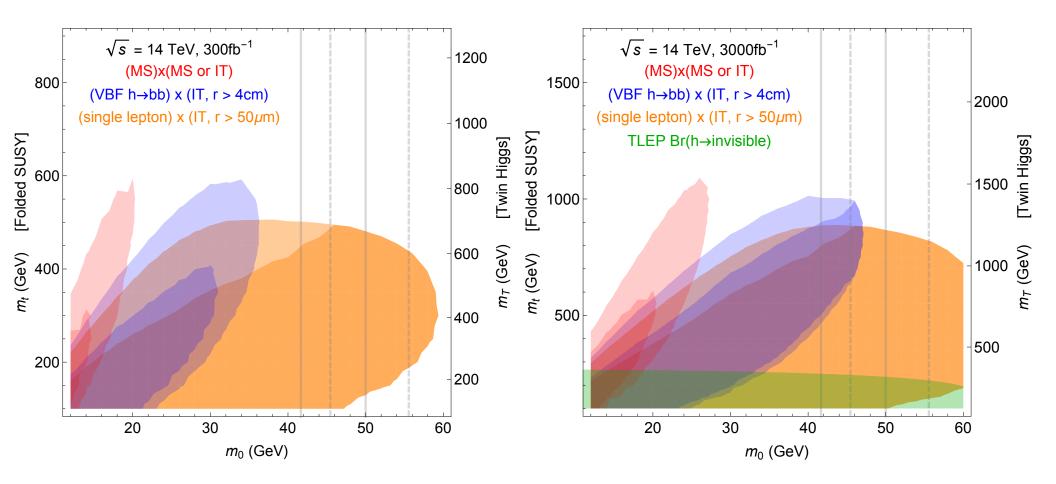
Let κ range from ~ 1/12 (democratic) to ~ 1 (optimistic).

DC, Verhaaren 1506.06141



LHC reach

ATLAS sensitivity projections to LHC14:

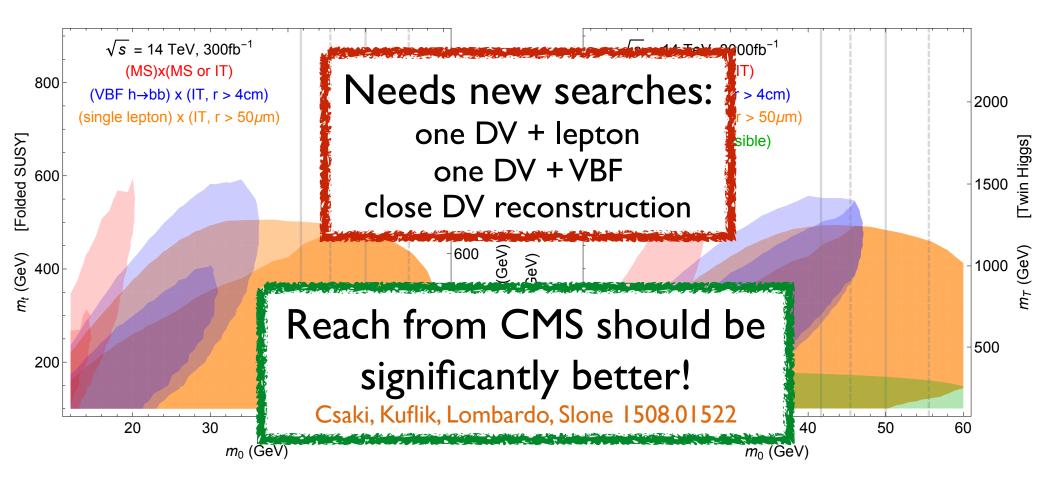


Displaced searches probe TeV-scale uncolored top partners!

DC, Verhaaren 1506.06141

LHC reach

ATLAS sensitivity projections to LHC14:

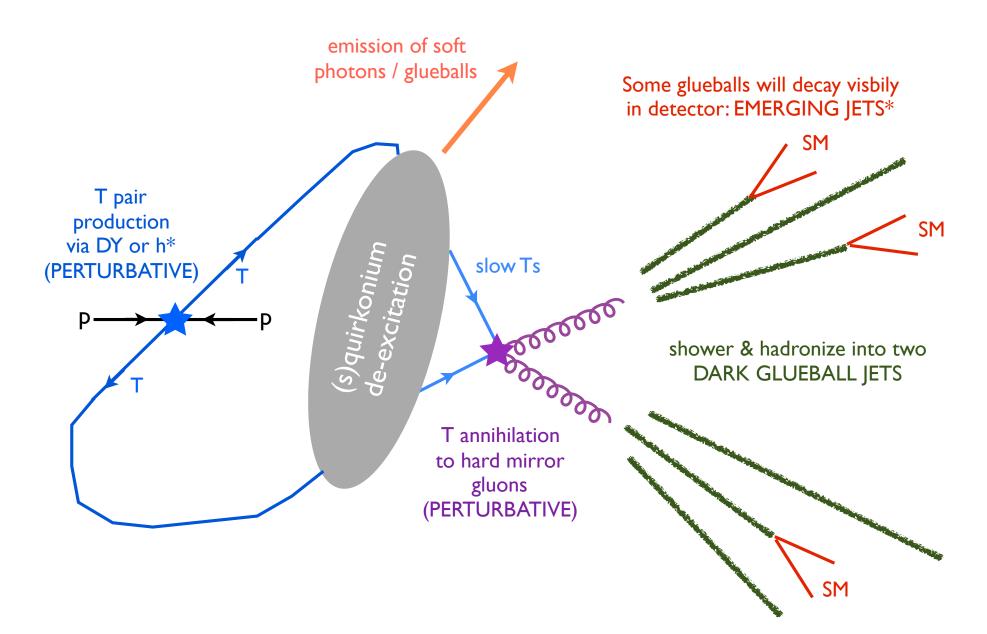


Displaced searches probe TeV-scale uncolored top partners!

DC, Verhaaren 1506.06141

Top partner direct production

Top partner direct production



Top partner direct production

Great opportunity:

- direct evidence of uncolored top partners.
- might have comparable reach to exotic Higgs decays
- could allow measurement of couplings and masses.
- potentiall spectacular signatures: several DVs, or many bb, $\tau\tau$ pairs

Main challenge: how to model mirror hadronization?

Have to parameterize our ignorance again, but this time use DGLAP evolution of hypothetical fragmentation functions to estimate glueball multiplicity and mean pT.

Work in progress!

Chacko, DC, Verhaaren, 1512.XXXXX

Prospects today

Displaced signatures are a great LHC opportunity, and a "smoking gun" for most theories with EW top partners (e.g. FSUSY). Can occur in some TH models.

Many signatures still unexplored, e.g. Flavor....

However: quasi-stable light mirror states are not guaranteed in theories of Neutral Naturalness.

What are the *unavoidable* signatures, at the LHC and at future lepton an 100 TeV colliders?

Probing Naturalness exhaustively: A No-Lose Theorem for Generalized Top

Partners.

Top Partners with SM Charge

Start with TeV-scale top partners that carry SM charge.

If QCD: produce plenty, discover at LHC or 100 TeV.

If partners carry any EW charge, regardless of decay mode etc, will be detectable up to ~ 2+ TeV @ 100 TeV due to RG effects in DY spectrum measurements!

Alves, Galloway, Rudermann, Walsh 1410.6810

TeV-scale SM-charged partners ARE DISCOVERABLE regardless of model details!

Neutral Top Partners

expalin twin higgs signatures how its tree tuning of soft z2 v

is that totally model-indep? no

We really only have one class of models for neutral top partners: Twin Higgs, which predicts Higgs coupling deviations ~ tuning at lepton colliders.

Is this general? Would like to understand signatures of neutral top partners model-independently!

Bottom-Up EFT/Simplified Model Approach!

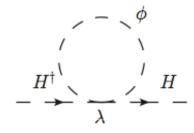
DC, Saraswat 1509.04284

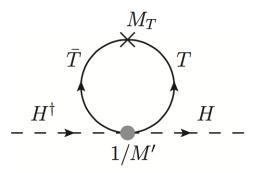
Two distinct low-energy EFTs

Scalar Partners

Fermion Partners

(Vector partners "same" as scalars)



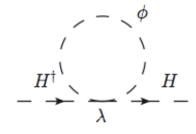


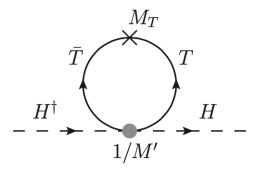
Two distinct low-energy EFTs

Scalar Partners

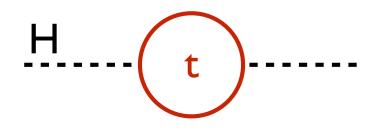
Fermion Partners

(Vector partners "same" as scalars)





Only impose *one* condition on EFT: cancellation of quadratic divergence from top loop

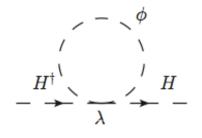


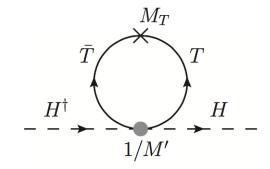
Two distinct low-energy EFTs

Scalar Partners

Fermion Partners

(Vector partners "same" as scalars)





Relevant terms in the HEFT expansion:

$$\mathcal{L}_{\phi} \supset -\sum_{i} \phi_{i}^{2} \left(rac{1}{2} \mu_{\phi_{i}}^{2} + rac{1}{2} \lambda_{i} |H|^{2}
ight) \qquad \qquad \mathcal{L}_{T} \supset \sum_{i} T_{i} ar{T}_{i} \left(M_{T_{i}} - rac{|H|^{2}}{2M_{i}'}
ight)$$

Condition to cancel one-loop quadratic divergence from top quark:

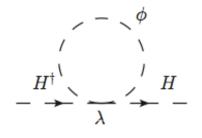
$$\lambda_{\phi} = \frac{12}{N_r} |y_t|^2 \qquad \qquad \frac{M_T}{M'} = \frac{3}{N_f} y_t^2$$

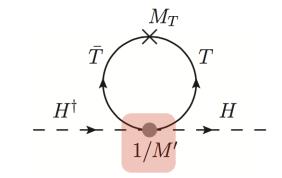
Two distinct low-energy EFTs

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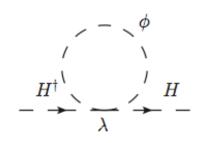
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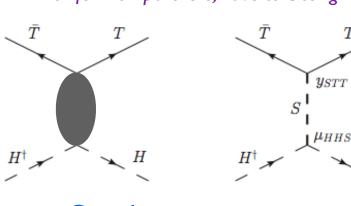
Non-renormalizable term limits what we can compute. Need partial UV completion for fermion partners!

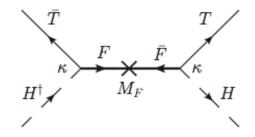
Scalar Partners

Fermion Partners

For fermion partners, have to distinguish how HHTT operator is generated.







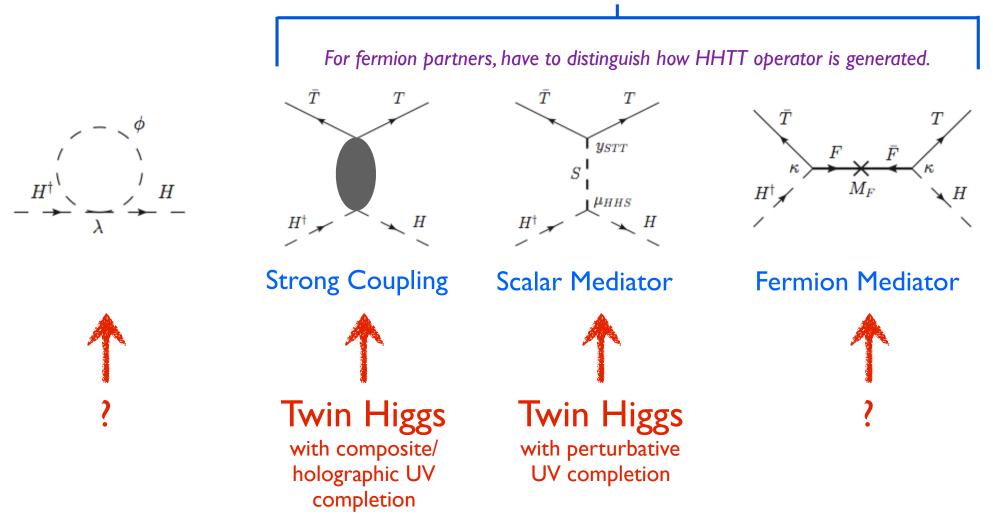
Strong Coupling

Scalar Mediator

Fermion Mediator

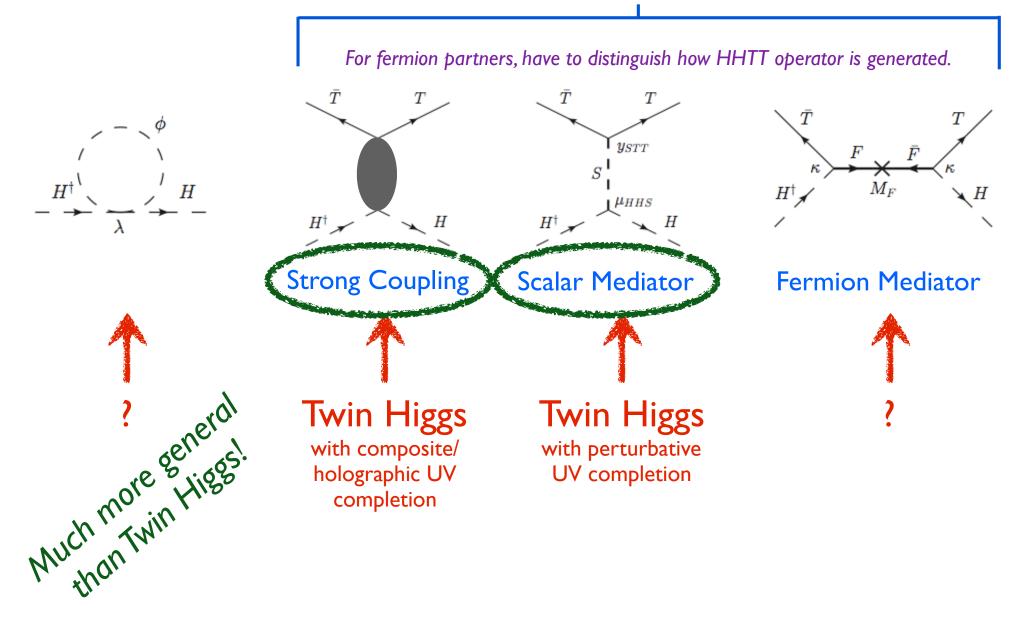
Scalar Partners

Fermion Partners



Scalar Partners

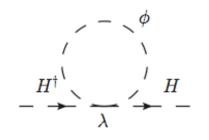
Fermion Partners

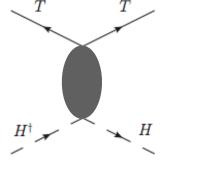


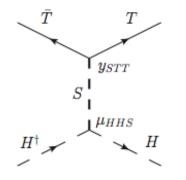
Scalar Partners

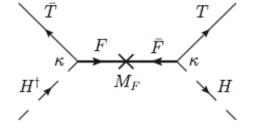
Fermion Partners

For fermion partners, have to distinguish how HHTT operator is generated.









Strong Coupling

Scalar Mediator

Fermion Mediator

For each scenario, analyze:

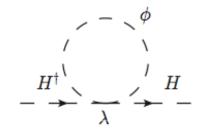
Irreducible low-E signatures:

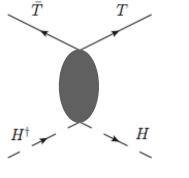
- Zh cross section (lepton collider)
- electroweak precision observables (lepton)
- higgs cubic coupling (100 TeV)
- top partner direct production (100 TeV)

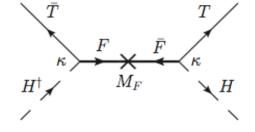
Scalar Partners

Fermion Partners

For fermion partners, have to distinguish how HHTT operator is generated.







Strong Coupling

Scalar Mediator

Fermion Mediator

For each scenario, analyze:

Irreducible low-E signatures:

- Zh cross section (lepton collider)
- electroweak precision observables (lepton)
- higgs cubic coupling (100 TeV)
- top partner direct production (100 TeV)

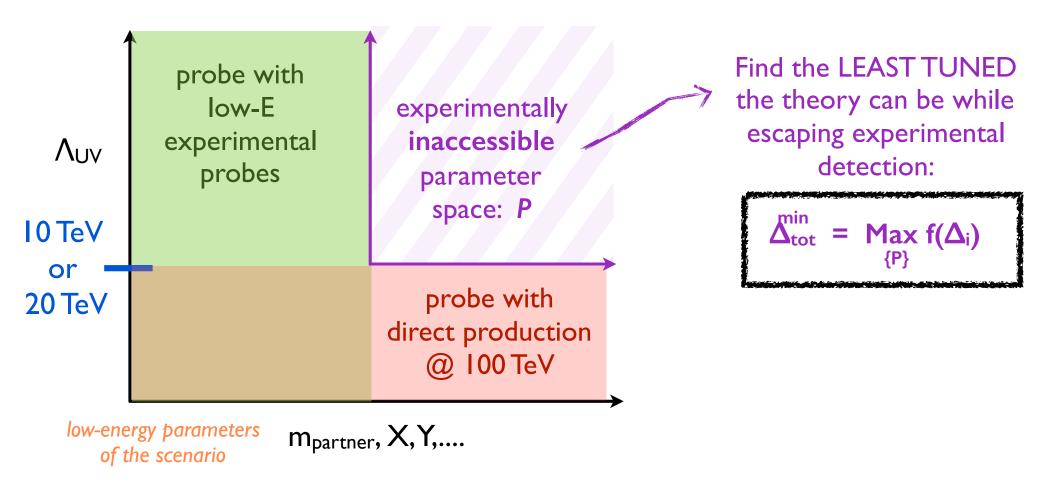
Irreducible tunings $\{\Delta_i\}$ of loop vs tree suffered by scenario $\Rightarrow \Delta_{tot} = f(\Delta_i)$

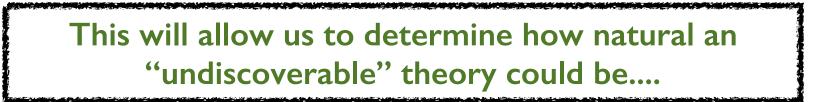
These will relate to UV completion scale Λ_{UV} .

Existing UV completions & symmetry arguments suggest SM-charged BSM states at this scale → Assume production at 100 TeV collider!

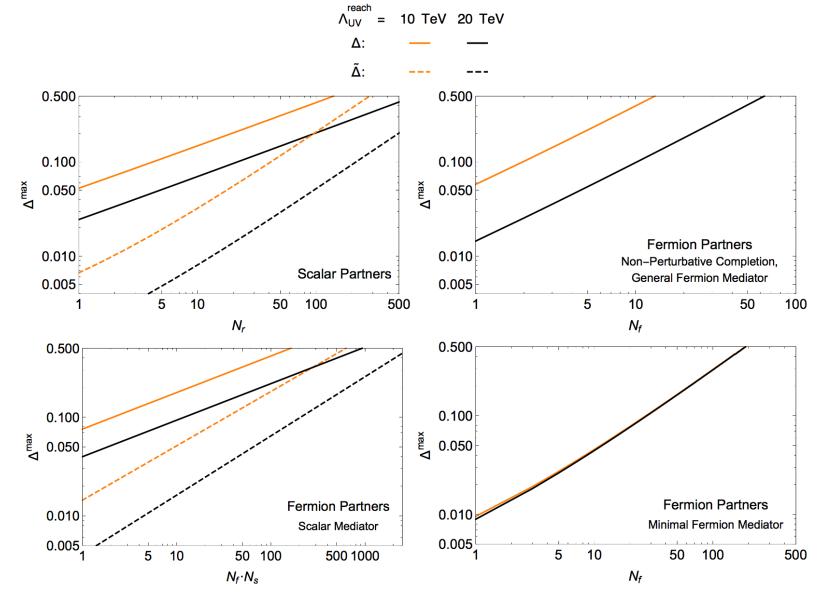
Strategy

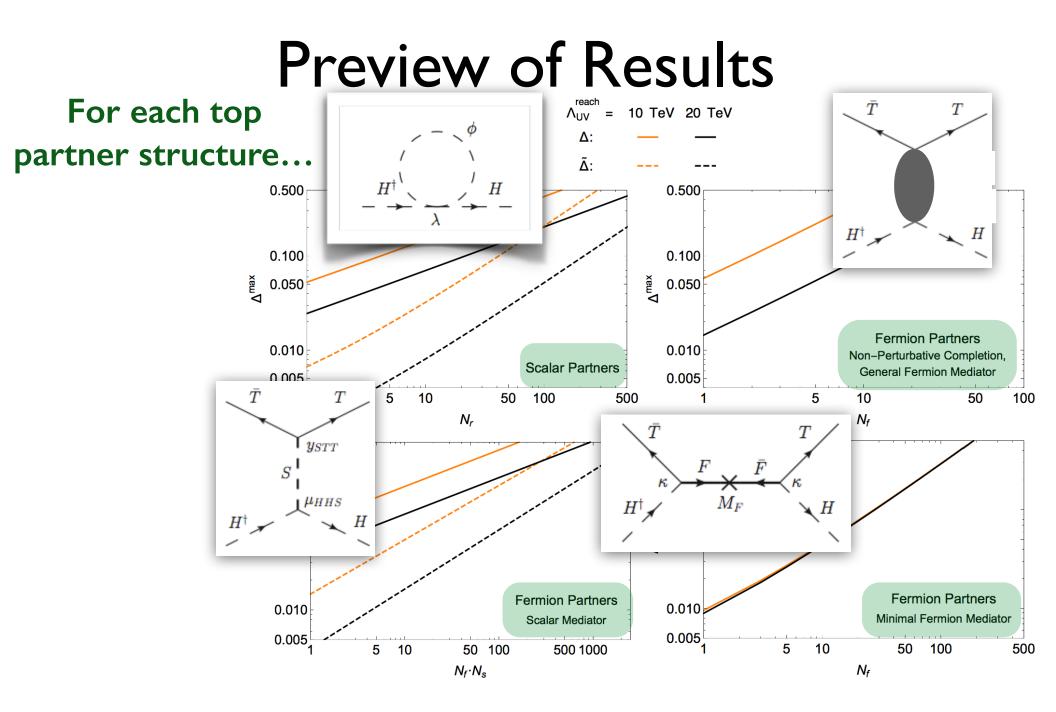
For each scenario:



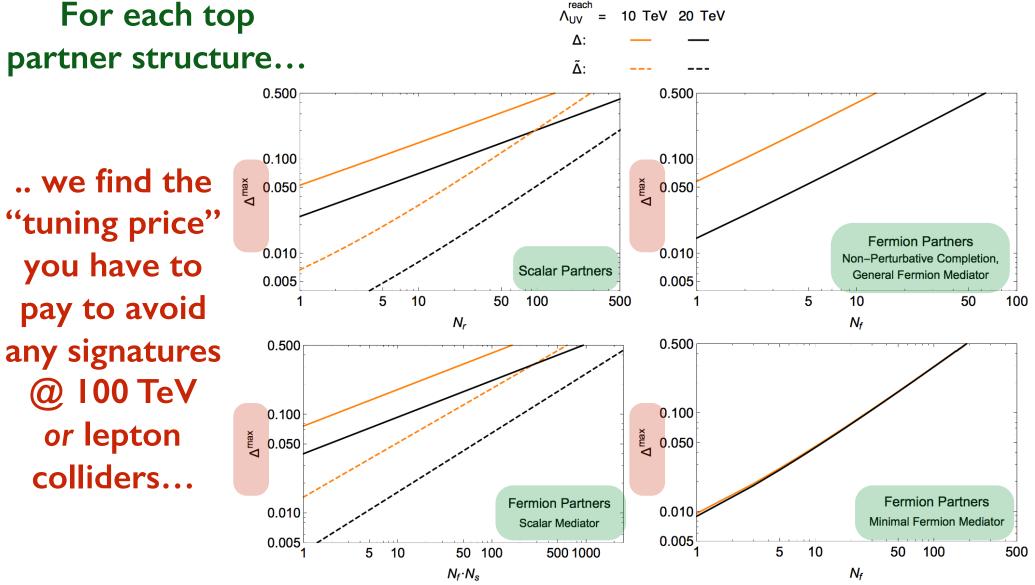


Preview of Results

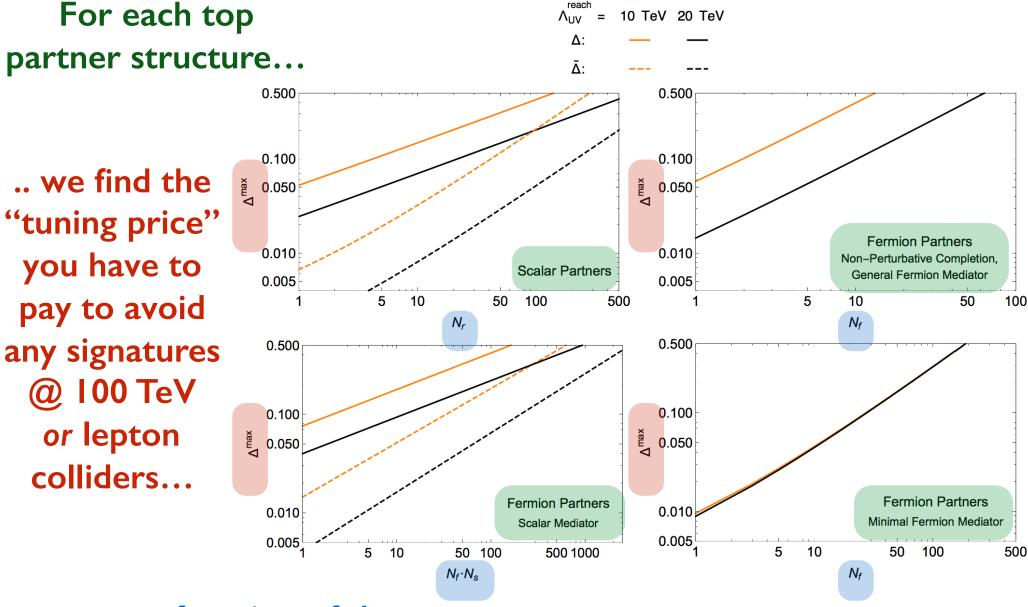




Preview of Results



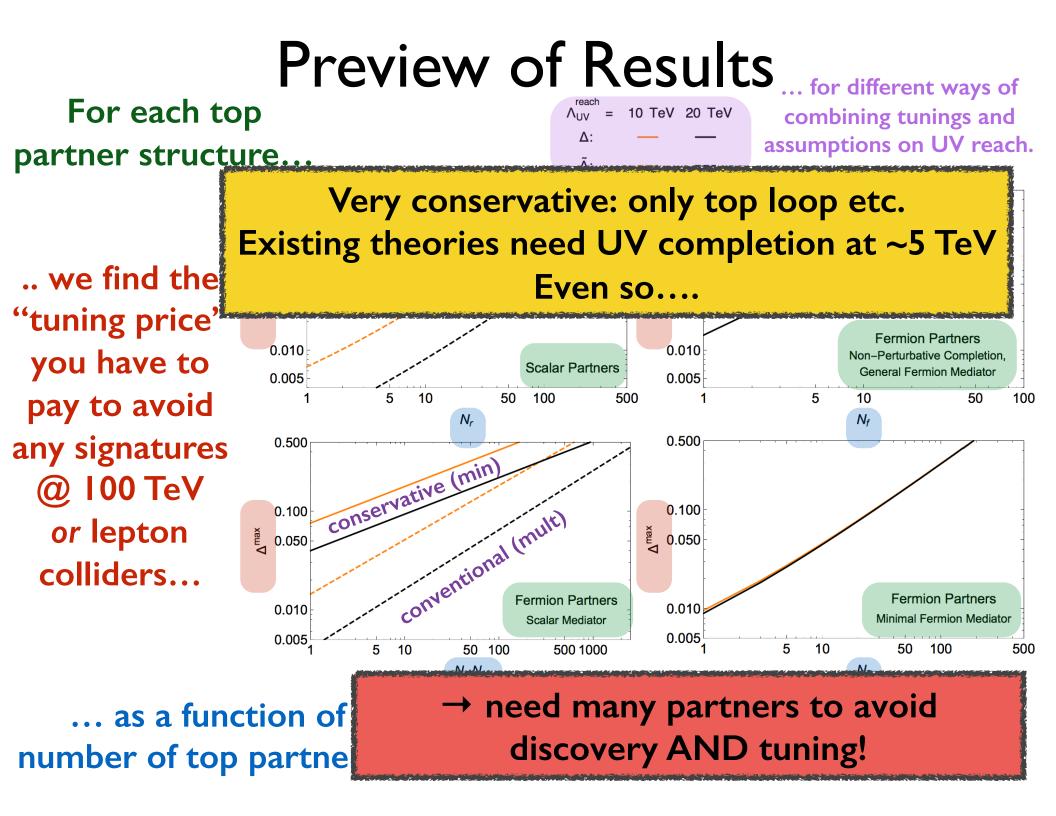
Preview of Results



... as a function of the number of top partner dof...

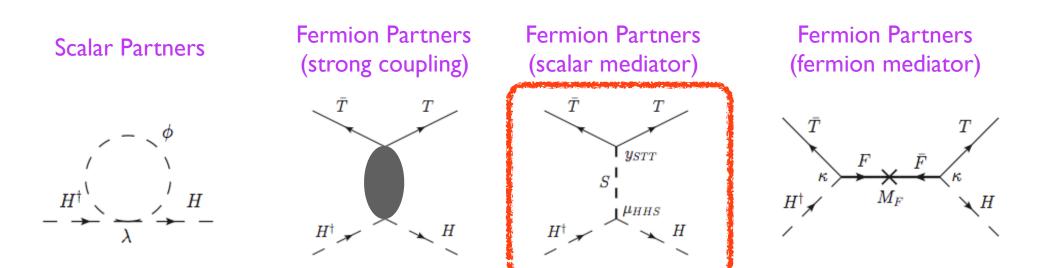
Preview of Results ... for different ways of For each top $\Lambda_{UV} = 10 \text{ TeV} 20 \text{ TeV}$ combining tunings and Δ: assumptions on UV reach. partner structure... Ã: 0.500 0.500 0.100 0.100 .. we find the A^{max} Δ^{max} 0.050 0.050 "tuning price" **Fermion Partners** 0.010 0.010 you have to Non-Perturbative Completion Scalar Partners General Fermion Mediator 0.005 0.005 pay to avoid 5 10 50 100 500 5 10 50 100 1 Nr Nf 0.500 any signatures 0.500 conservative (min) @ 100 TeV conventional (mult) 0.100 0.100 or lepton 0.050 ∆^{max} 0.050 colliders... Fermion Partners **Fermion Partners** 0.010 0.010 Minimal Fermion Mediator Scalar Mediator 0.005 0.005 5 500 5 10 50 100 500 1000 10 50 100 $N_f \cdot N_s$ Nf

... as a function of the number of top partner dof...



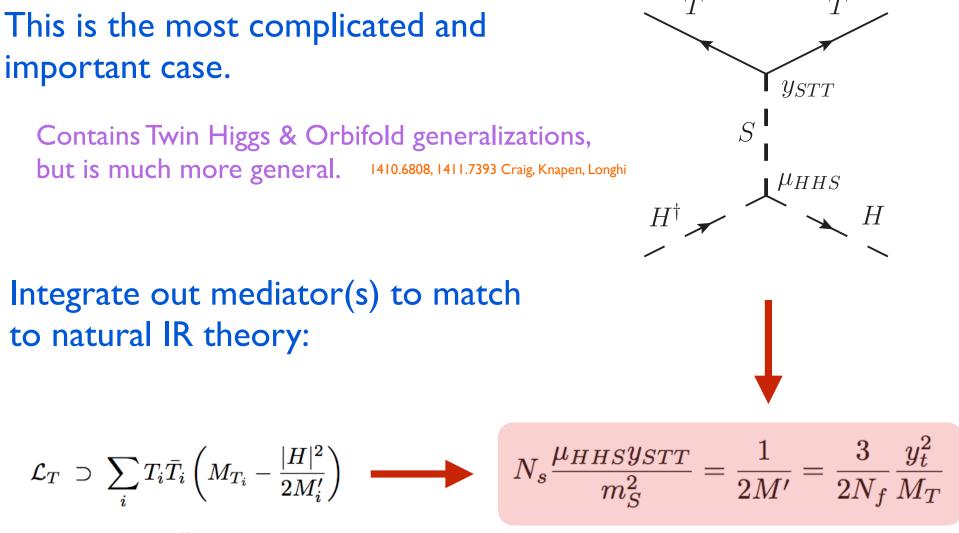
How do we get there?

Neutral Naturalness Scenarios



Trickiest/most interesting case to analyze in complete generality...

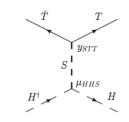
Fermion Partner - Scalar Mediator



low-energy effective Lagrangian to cancel top loop

naturalness matching condition

The Scalar Mediator

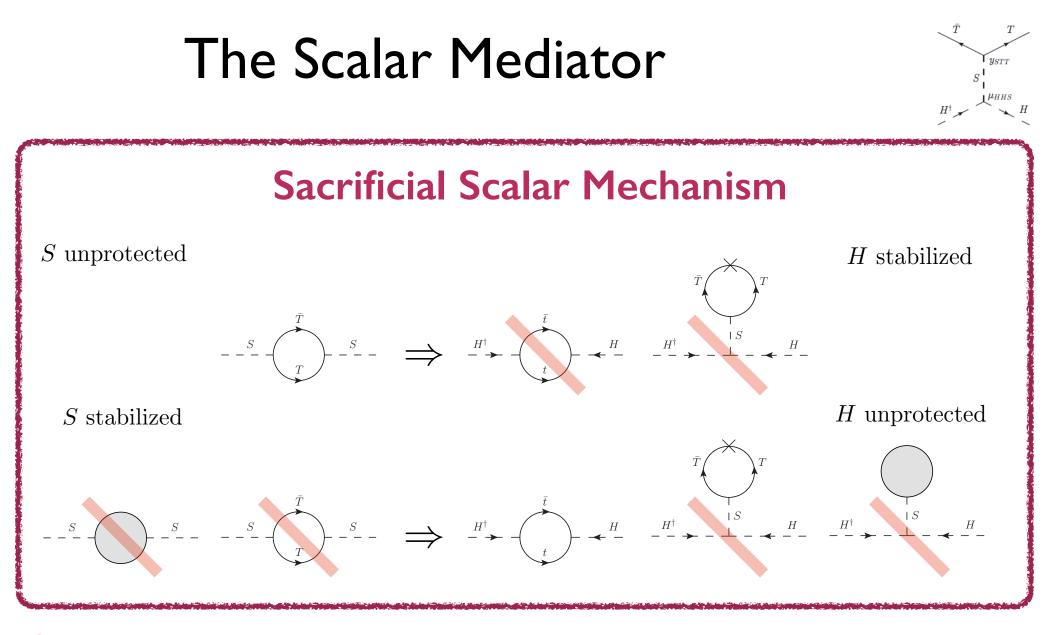


Before we can proceed, we have to know: How heavy is the scalar mediator?

Naive expectation: new scalars can't be light, otherwise we have another hierarchy problem! \Rightarrow m_s should be significantly above weak scale!

Naive counterargument: we know of many ways to solve the hierarchy problem! Dress up mediator sector with partners etc...

Nope!



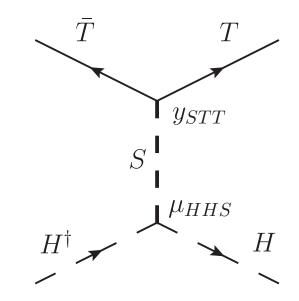
Consequences:

- I. Mass of scalar is tied to UV completion scale!
- 2. $m_s >> m_h$ makes it easy to compute experimental signals.

Higgs Mixing

Take one scalar mediator S

(generalizes simply)



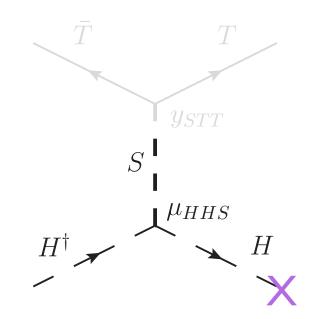
Higgs Mixing

Take one scalar mediator S

(generalizes simply)

In the $m_s >> m_h$ limit, mixing angle is simple:

$$s_{\theta} \approx -\frac{\mu_{HHS}}{m_S^2} v$$



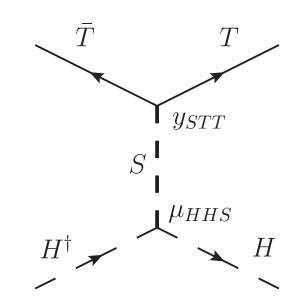
Computing Observables



(generalizes simply)

In the $m_s >> m_h$ limit, mixing angle is simple:

$$s_{\theta} \approx -\frac{\mu_{HHS}}{m_S^2} v$$



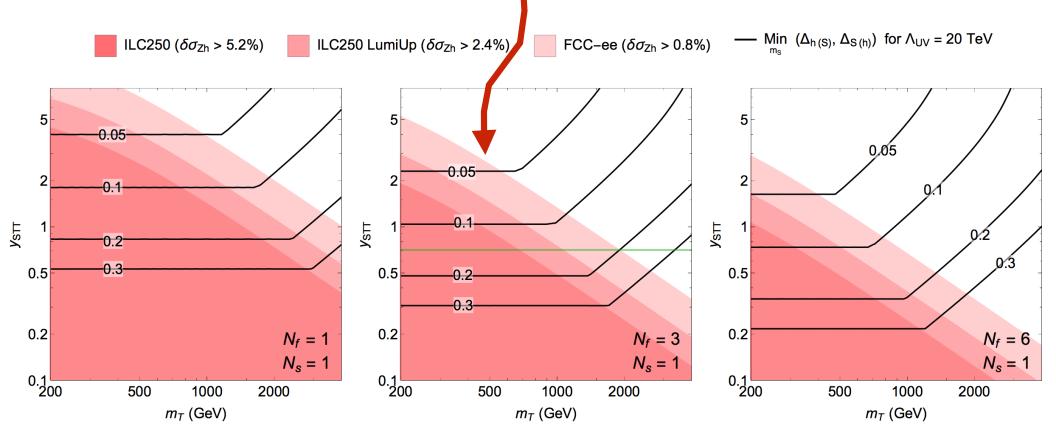
Naturalness condition:

$$\frac{\mu_{HHS}y_{SST}}{m_S^2} = \frac{3}{2N_f}\frac{y_t^2}{M_T} \longrightarrow s_\theta \approx -\frac{3}{2N_f}\frac{y_t^2}{y_{SST}}\frac{v}{M_T}$$

Mediator mass drops out! Only depends on (MT, YSST)

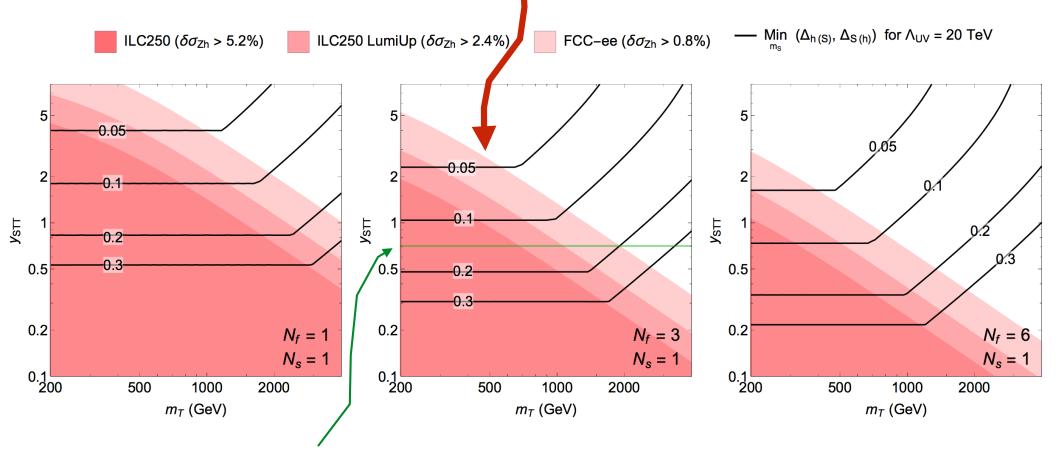
Higgs Mixing in (m_T, y_{STT}) Plane

Lepton colliders have great sensitivity in much of parameter space.



Higgs Mixing in (m_T, y_{STT}) Plane

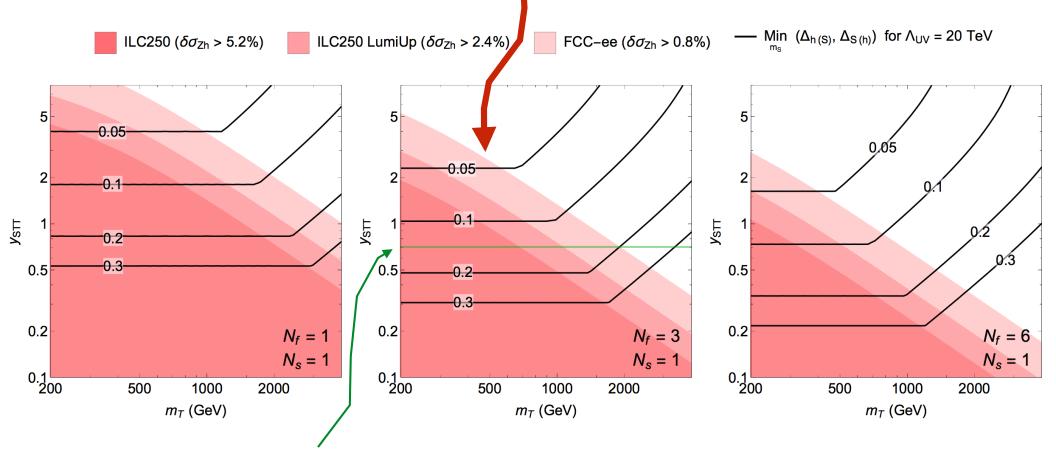
Lepton colliders have great sensitivity in much of parameter space.



Twin Higgs models are subspaces (lines) in this more general parameter space.

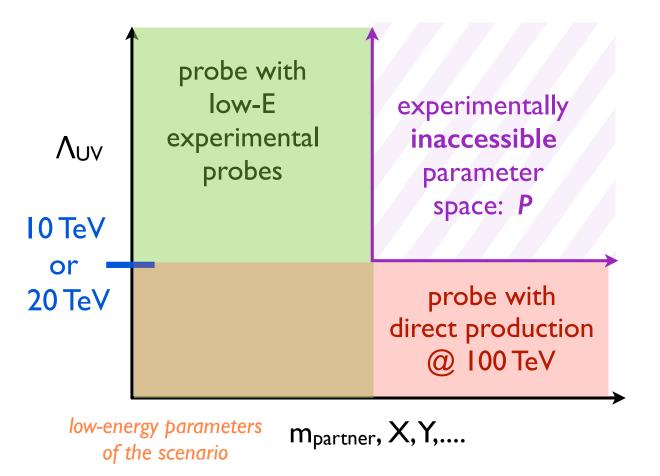
Higgs Mixing in (m_T, y_{STT}) Plane

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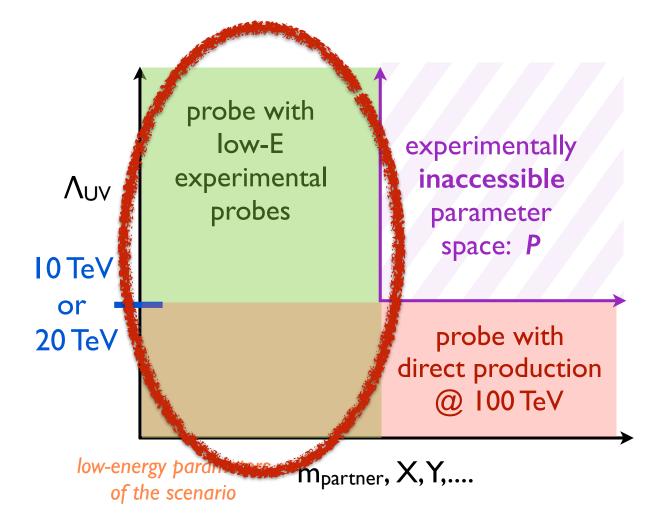


Twin Higgs models are subspaces (lines) in this more general parameter space.

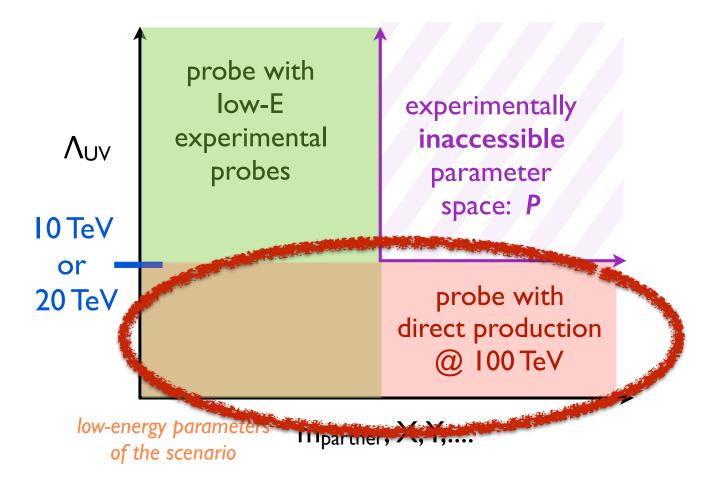
But what if y_{STT} is large??



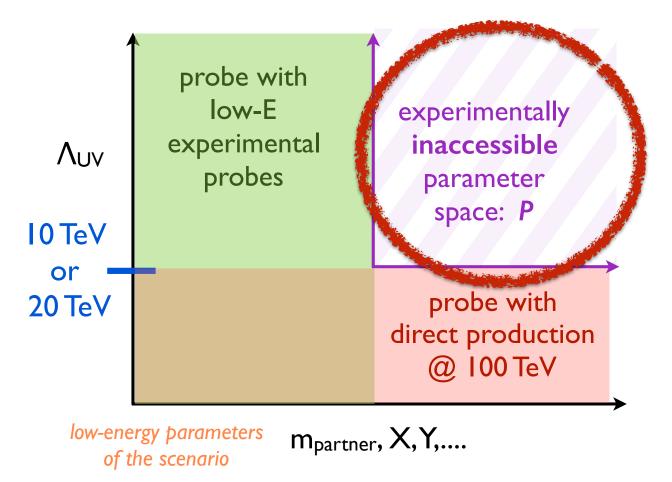
We've determined the reach of low-energy observables (higgs mixing).



Now we exploit the 100 TeV collider's ability to probe the UV scale.



Assuming 10 or 20 TeV can be probed, what unavoidable tuning are we stuck with?



Tunings (I)

 $\Delta_{h(S)}$ = log tuning of m_h from mediator loops.

(have to differentiate case where Higgs = PNGB from case without such symmetries....)

Gets worse with large ms!

 $\Delta_{S(T)} = \text{tuning from quadratic sensitivity of } m_S \text{ to } T \text{ loops}$ (required by Sacrificial Scalar Mechanism!) **Gets better with large ms!** Can find conservative tuning estimate by maximizing over (unknown) mediator mass!

$$\Rightarrow \Delta_{H,S} = \max_{m_S} f(\Delta_{h(S)}, \Delta_{S(T)})$$

Tunings (I)

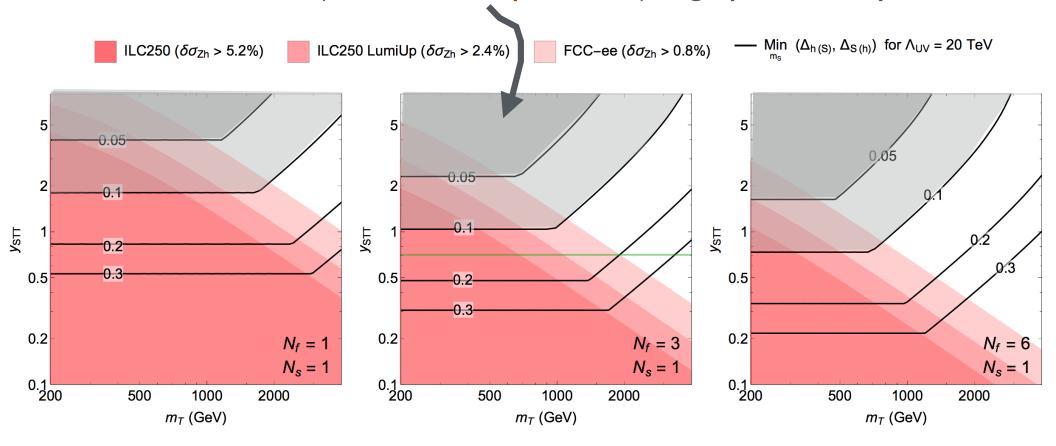
 $\Delta_{h(S)} = \log \text{ tuning of } m_h \text{ from mediator loops.}$ (have to differentiate case where Higgs = PNGB from case without such symmetries....) $Gets \text{ worse with large } m_s!$ $\Delta_{S(T)} = \text{ tuning from quadratic sensitivity of } m_s \text{ to } T \text{ loops}$ (required by Sacrificial Scalar Mechanism!) $Gets \text{ better with large } m_s!$ $\Rightarrow \quad \Delta_{H,S} = Max f(\Delta_{h(S)}, \Delta_{S(T)})$

 $\Rightarrow \Delta_{H,S} = \underset{m_{S}}{\text{Max}} f(\Delta_{h(S)}, \Delta_{S(T)})$

Since we marginalize over m_S , $\Delta_{H,S}$ is uniquely defined in the (m_T, y_{STT}) plane as the tuning from the mediator sector.

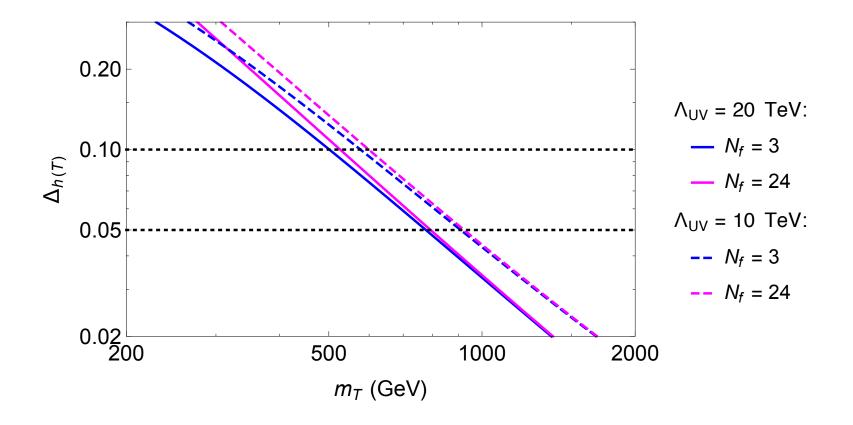
Tuning from Mediator in (m_T, y_{STT}) Plane

For $\Lambda_{UV} \ge 20 \text{ TeV}$ (undetectable by 100 TeV), high y_{STT} is badly tuned!



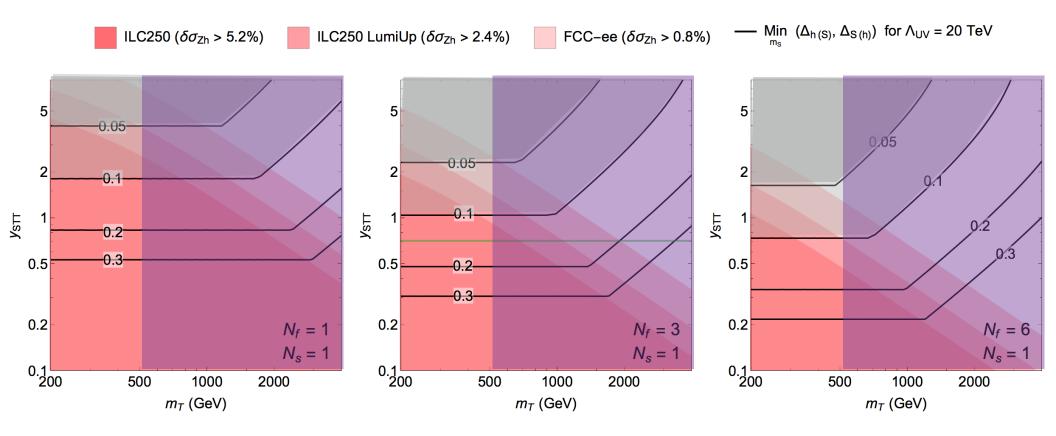
Tunings (2)

For $\Lambda_{UV} \ge 20 \text{ TeV}$ (undetectable by 100 TeV), top partners heavier than ~500 GeV give log-tuning to Higgs mass worse than 10%



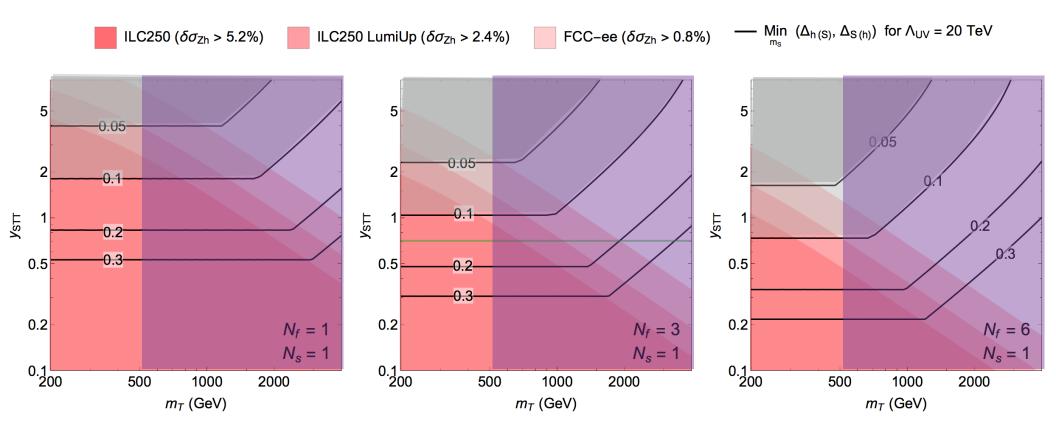
Log tuning from t vs T in (m_T, y_{STT}) Plane

For $\Lambda_{UV} \ge 20 \text{ TeV}$ (undetectable by 100 TeV), top partners heavier than ~500 GeV give log-tuning to Higgs mass worse than 10%



Log tuning from t vs T in (m_T, y_{STT}) Plane

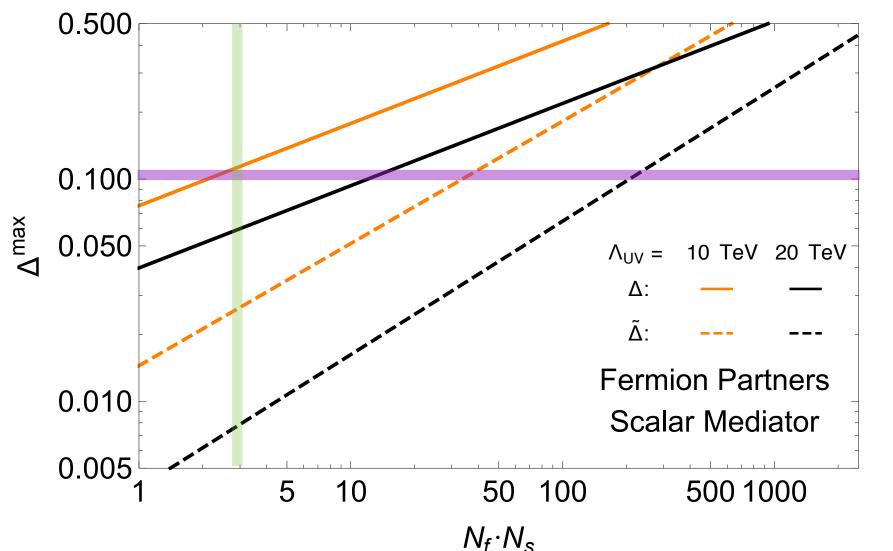
For $\Lambda_{UV} \ge 20 \text{ TeV}$ (undetectable by 100 TeV), top partners heavier than ~500 GeV give log-tuning to Higgs mass worse than 10%



No untuned parameter space left for $N_f \times N_s \sim O(SM)!$

Fermion Partner - Scalar Mediator

USTT



A natural theory needs to have VERY MANY fermion partners/scalar mediators to possibly escape detection.

Need both colliders for full coverage!

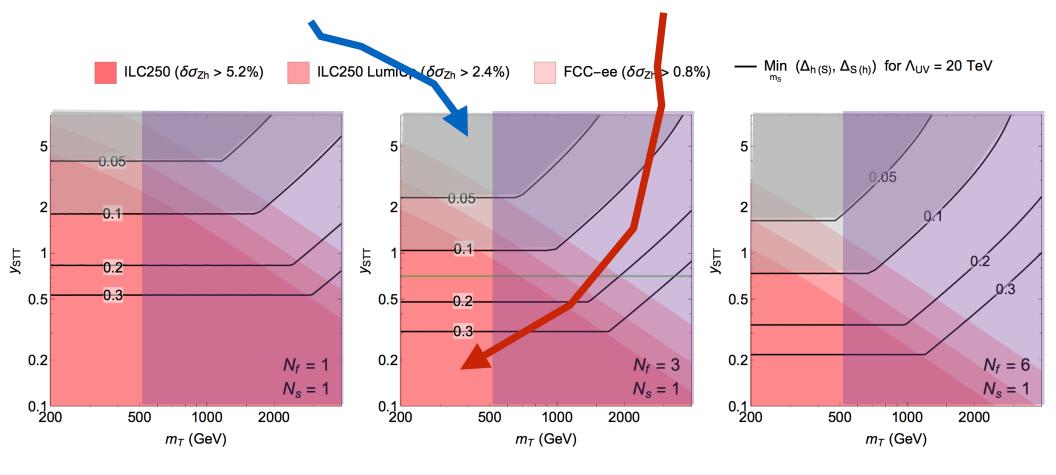
Large hidden sector coupling: Higgs mixing is tiny, but need low Λ_{UV} .

No guaranteed signal at lepton collider, but slam dunk at 100 TeV!

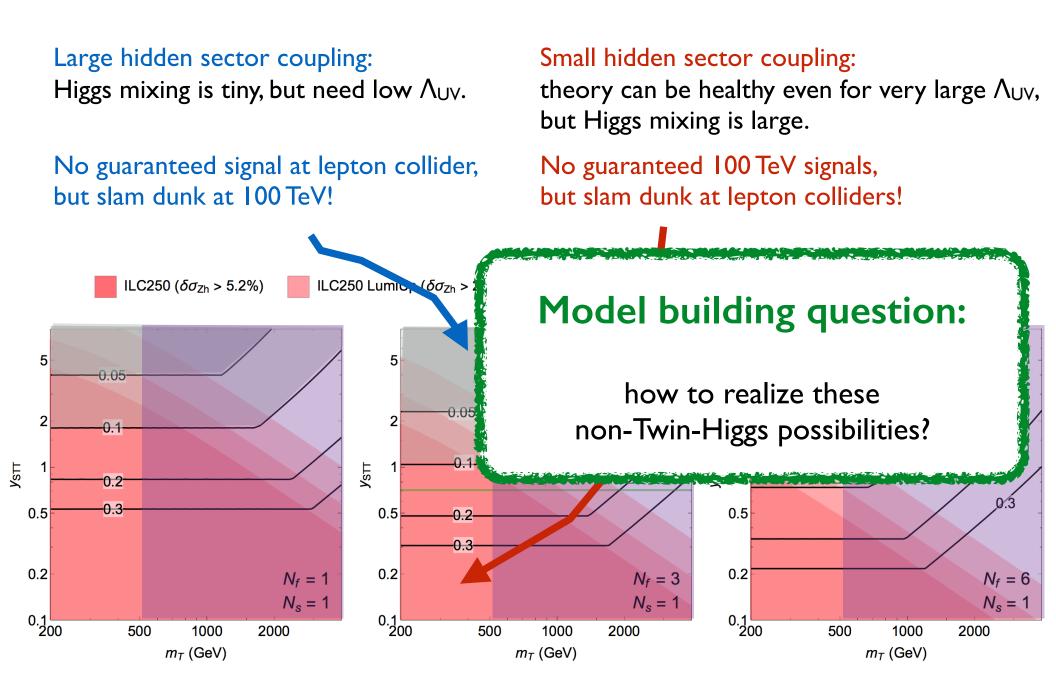
Small hidden sector coupling:

theory can be healthy even for very large $\Lambda_{UV},$ but Higgs mixing is large.

No guaranteed 100 TeV signals, but slam dunk at lepton colliders!



Need both colliders for full coverage!



... go through corresponding derivations for the other scenarios, with similar conclusions....

What's the upshot?

Great discovery potential TODAY

Long-lived hidden sector states (mirror glueballs, quarkonia) generate spectacular displaced signals that allow the LHC to probe TeV uncolored top partners

2. Implications for LHC searches

Displaced Vertex searches with just one DV + VBF or lepton are required. Also, need sub-mm decay length reconstruction.

3. No-Lose Theorem:

Any theory of ~10% naturalness with O(SM) top partners will be discovered at a planned lepton collider and/or 100 TeV

→ Model-independent (bottom-up) and very conservative (only top loop etc)

How to avoid this theorem?

Could have top partner swarms, or neutral top partners without SM charges in UV completion.

There might also be weird non-perturbative or stringy constructions that don't need top partners?

4. Implications for future colliders

Both lepton collider and 100 TeV have to work in tandem for full coverage of general naturalness

Without lepton collider:

could miss theory with large-ish Higgs mixing but small hidden sector couplings → very high UV completion scale out of 100 TeV collider reach

Without 100 TeV:

several scenarios give small IR signatures, need to probe UV

5. For full coverage, need to probe UV completion!

Central assumption of SM-charged BSM states at Λ_{UV} allows us to make these very powerful conclusions.

This seems very reasonable, and is certainly the case in all currenty proposed UV completions.

Can we formally prove this always has to be the case, or construct counter-examples?

Summary

I. Discovery potential TODAY

Neutral naturalness motivates spectacular displaced signatures that give the LHC TeV-reach for uncolored top partners.

2. Implications for LHC searches

Need searches with just one DV + lepton or VBF, and sub-mm decaylength reconstruction for full coverage

3. No-Lose Theorem

Any theory of ~10% naturalness with O(SM) top partners will be discovered at a planned lepton collider and/or 100 TeV

4. Implications for future colliders Both lepton collider and 100 TeV have to work in tandem for full coverage of general naturalness

5. Probing UV completion is vital!

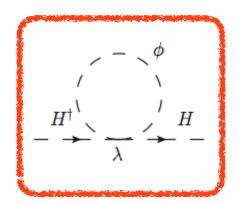
Can we formally prove that full that SM-charged BSM states appear at Λ_{UV} in full symmetry-based theories?

Thank you!

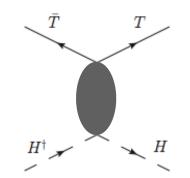
Backup Slides

Neutral Naturalness Scenarios

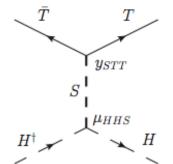
Scalar Partners



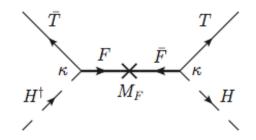
Fermion Partners (strong coupling)



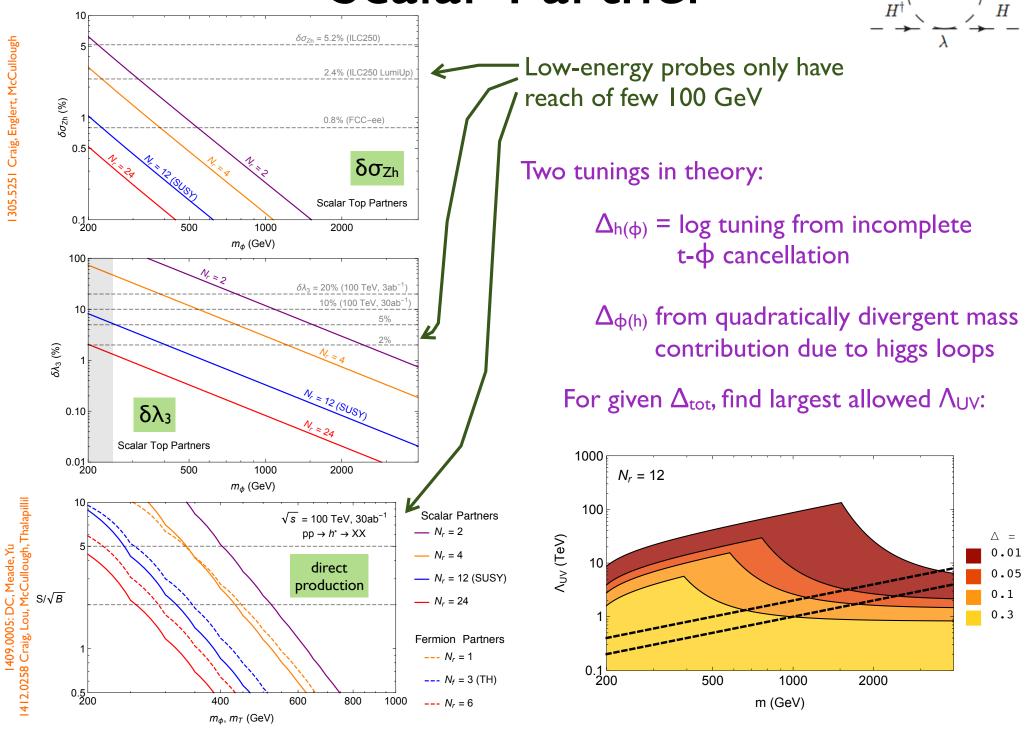
Fermion Partners (scalar mediator)

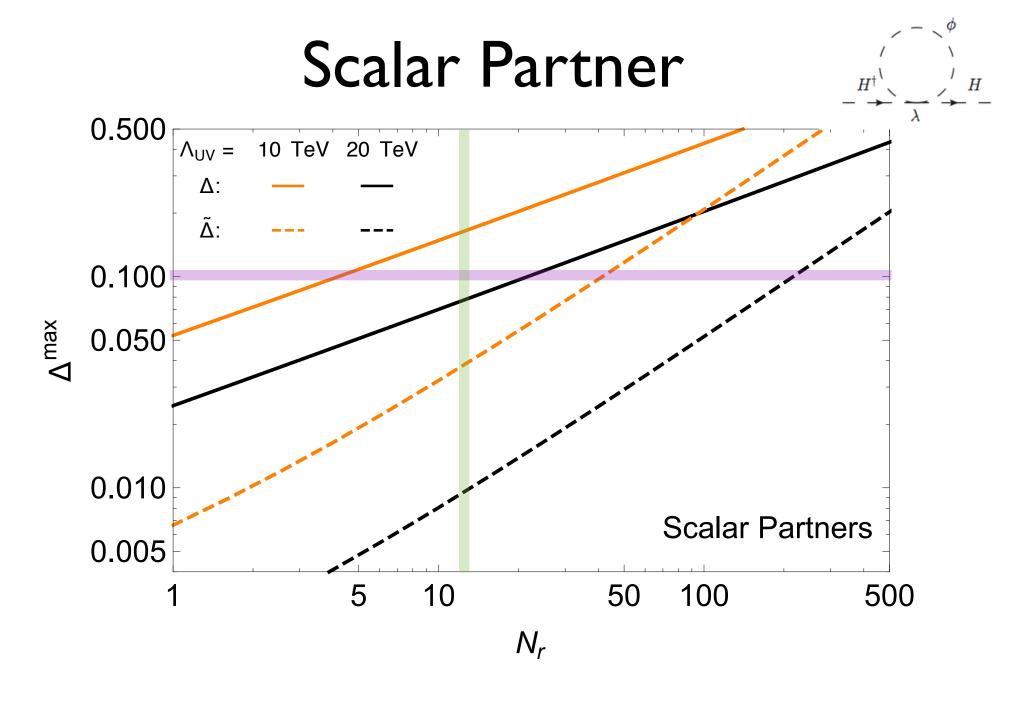


Fermion Partners (fermion mediator)



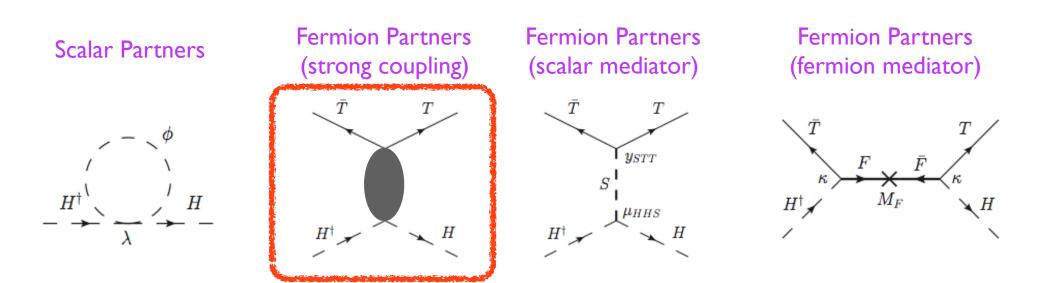
Scalar Partner



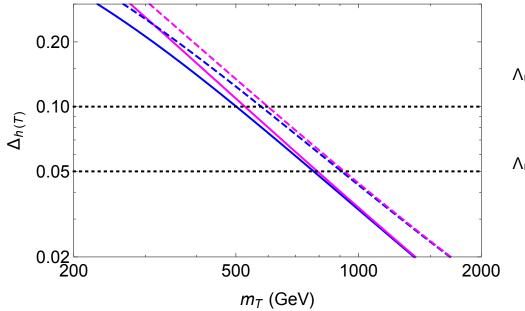


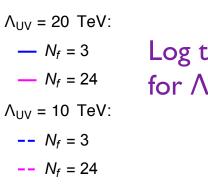
A natural theory needs to have VERY MANY scalar partners to possibly escape detection.

Neutral Naturalness Scenarios



Fermion Partner - Strong Coupling



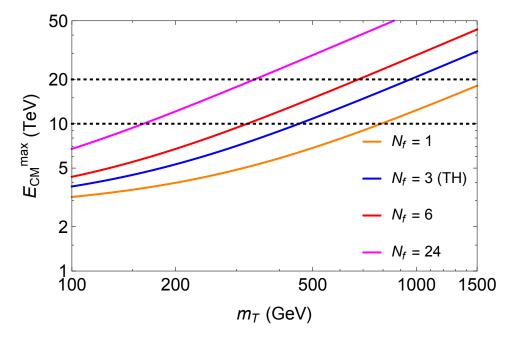


Log tuning of higgs mass: for $\Lambda_{UV} < 10 - 20$ TeV, $m_T \lesssim 500$ GeV OR tuning worse than 10%.

Т

H

 H^{\dagger}



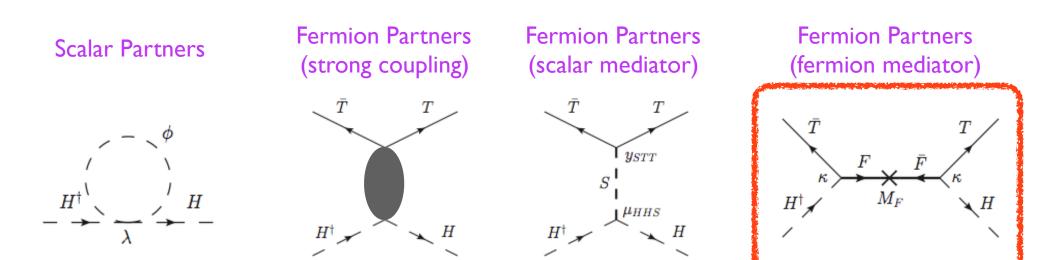
Unitarity constraints place strict upper bound on Λ_{UV} where new physics must get resolved.

TFermion Partner - Strong Coupling H^{\dagger} H0.500 0.100 Δ^{max} 10 TeV 20 TeV $\Lambda_{UV} =$ 0.050 Δ: Ã: **Fermion Partners** 0.010 Strong Coupling 0.005 100 5 10 50

A natural theory needs to have VERY MANY fermion partners to possibly escape detection.

 N_{f}

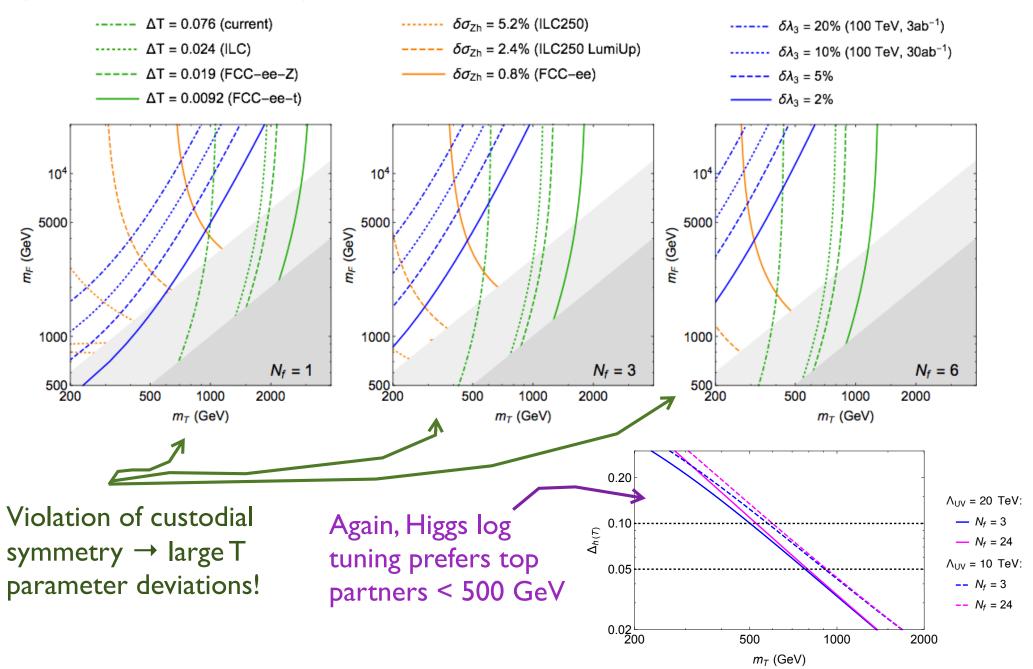
Neutral Naturalness Scenarios



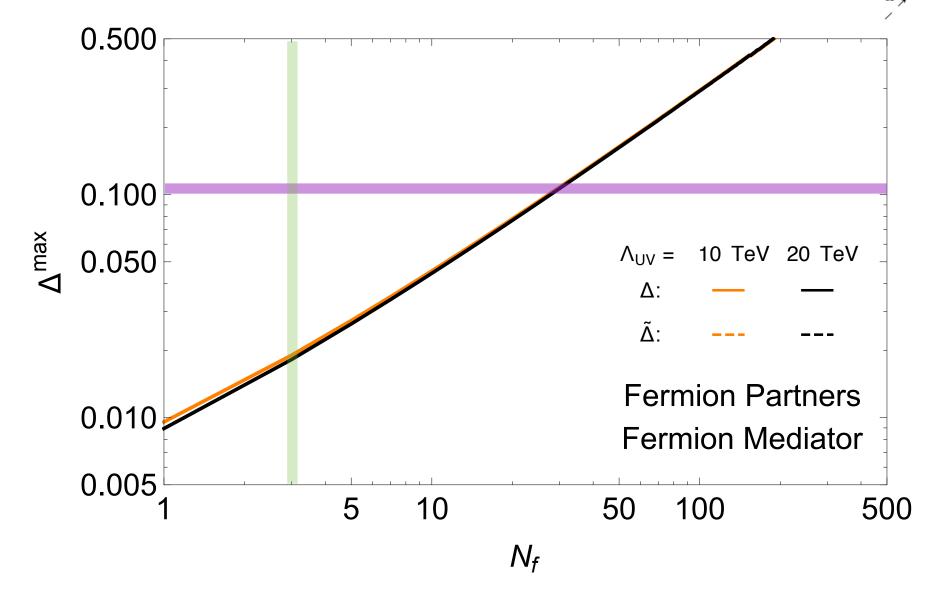
Fermion Partner - Fermion Mediator

H

using results from 1506.0546 Fedderke, Lin, Wang



Fermion Partner - Fermion Mediator



A natural theory needs to have VERY MANY fermion partners to possibly escape detection.