

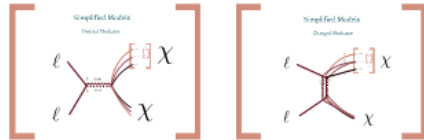
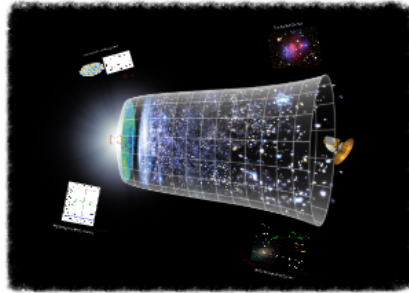
Conclusions

Leptophilic dark matter has novel phenomenology

Can potentially explain the anomalous magnetic moment of the muon

For chiral couplings, new physics scale is $O(100)$ GeV

Under tension from direct detection and collider constraints



Leptophilic dark matter

Tree-level interactions with SM leptons

Charged mediators Neutral mediators

Interaction with quarks at one-loop level

Why Leptophilic?

See also lectures for QED which is well understood (spin-1)

$\mu_{\text{anom}} = \frac{e^2}{4\pi^2} \frac{1}{m} \sum_f Q_f^2 g_{\ell f}^2$
 $\mu_{\text{anom}} = \frac{e^2}{4\pi^2} \frac{1}{m} \sum_f Q_f^2 g_{\ell f}^2 \approx 3 \times 10^{-11} \mu_B$
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(1) Both fermions are leptons (e.g. μ, τ, ν_μ)
 (2) Leptons with higher charges are preferred for direct detection
 (3) Leptons with higher charges are preferred for direct detection

Which Leptophilic?

Charged Mediators

Loop diagrams with charged leptons require additional diagrams

See also charged mediator couplings to right-handed leptons only

Neutral Mediators

See also neutral mediator couplings to charged leptons only

Which Leptophilic?

See also couplings to the neutrinos of charged mediators

Two broad possibilities

Current mediators: ν_μ, ν_τ
 Charged mediators: e, μ, τ

Checklist of model constraints can be found here: [https://arxiv.org/abs/1708.07457](#)

We assume flavor universal couplings

Direct Detection

g-2



Colliders

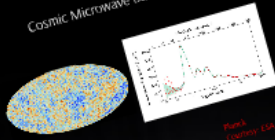
Dark Matter: The Lepton Connection

Prateek Agrawal
Fermilab

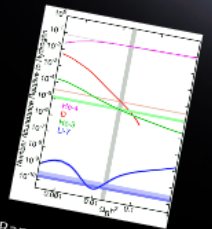
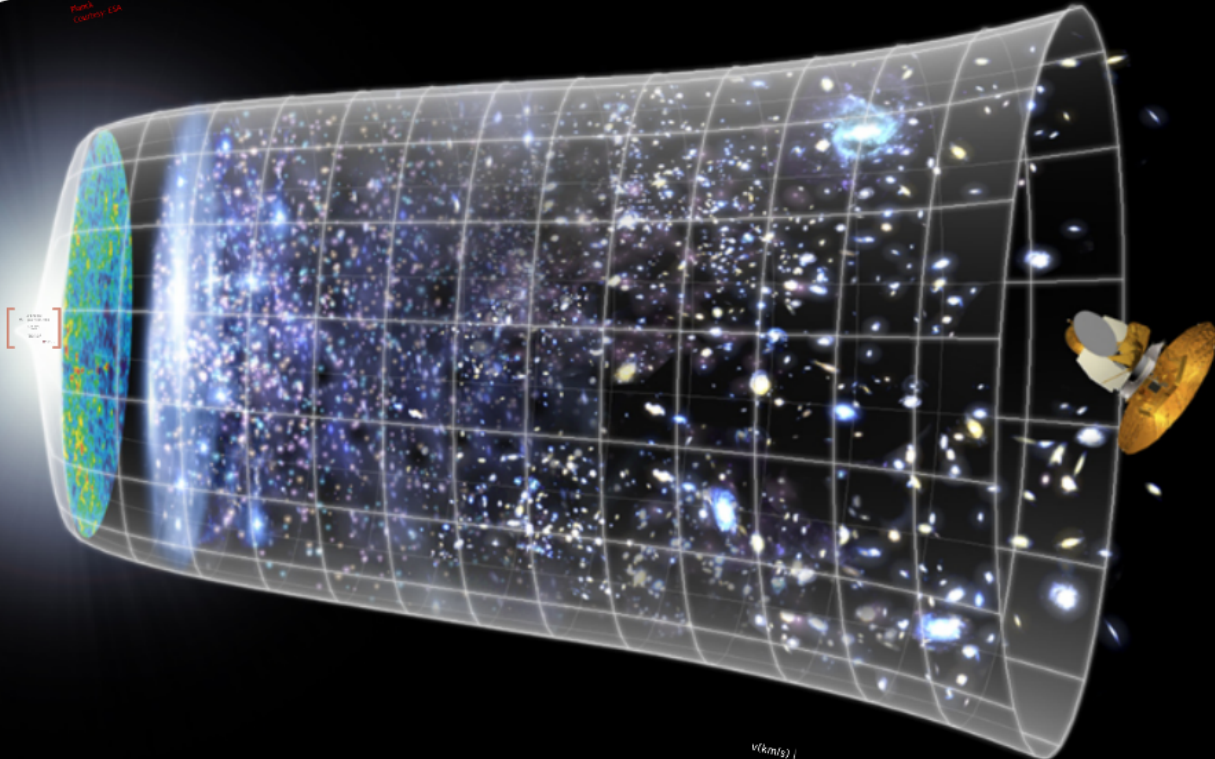
Rutgers University
April 29, 2014

| arXiv:1402.7369
| PA, Z. Chacko, C. Verhaaren

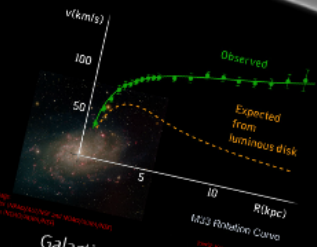
Cosmic Microwave Background



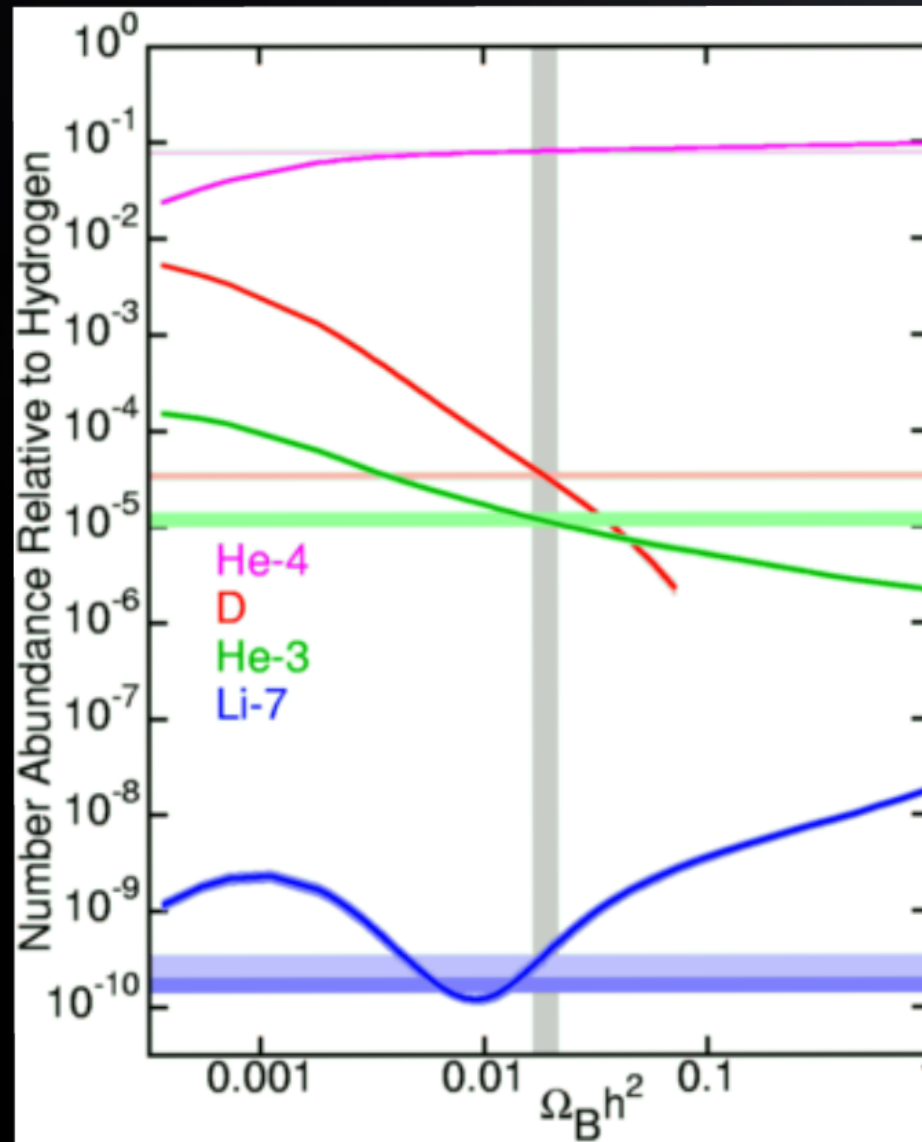
The Bullet Cluster



Big Bang Nucleosynthesis



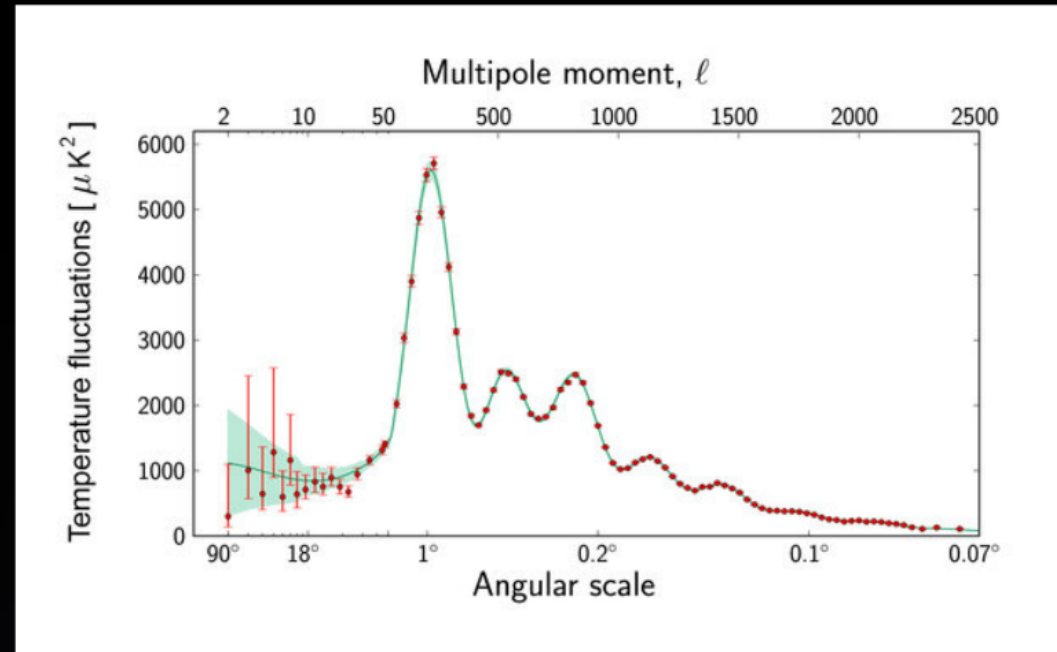
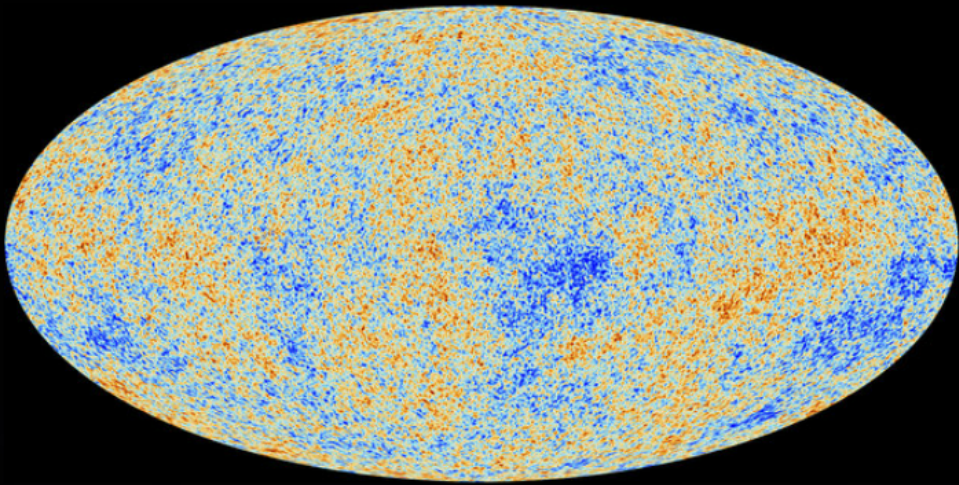
Galactic Rotation Curves



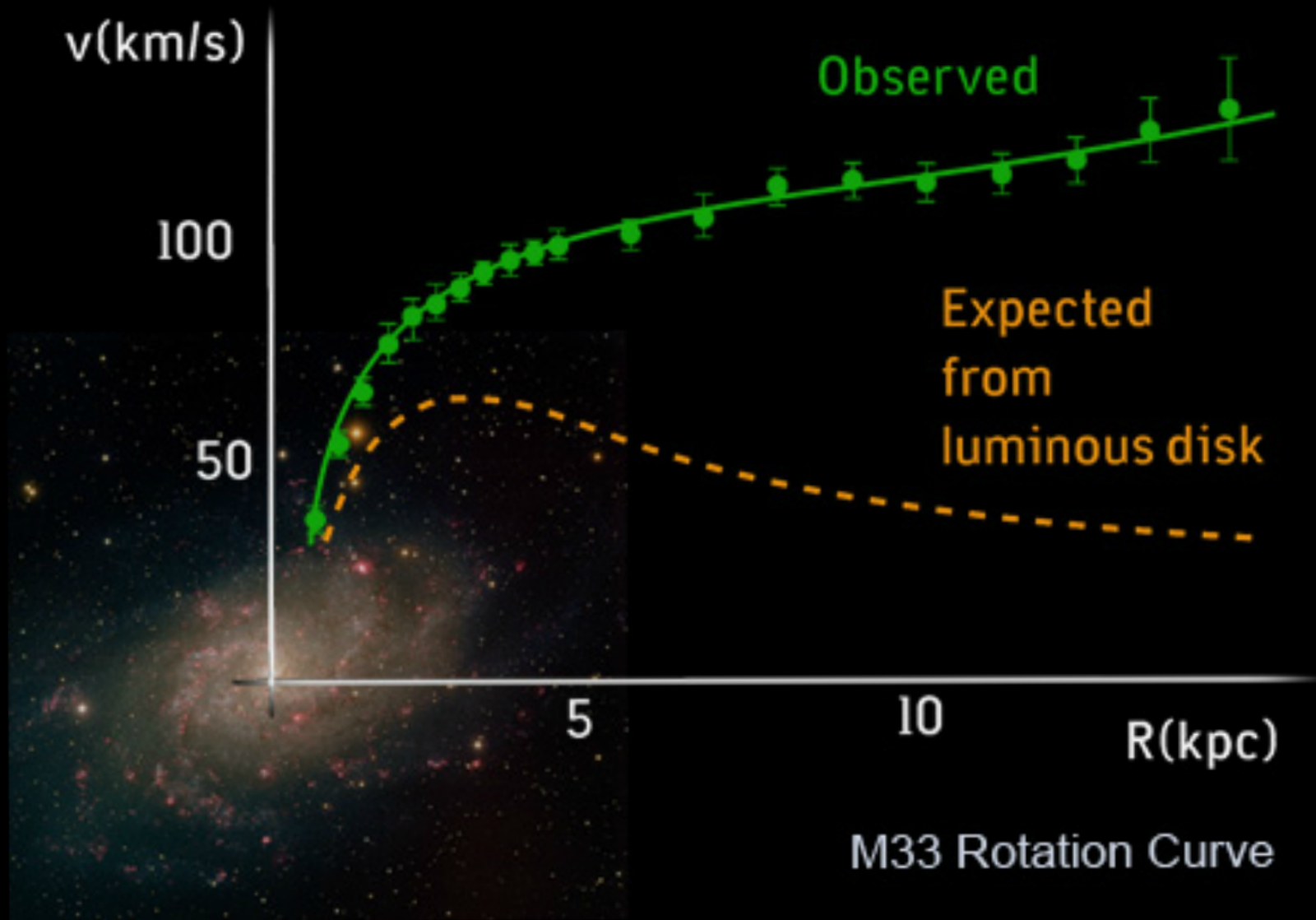
[credit: Edward L. Wright]

Big Bang Nucleosynthesis

Cosmic Microwave Background



Planck
Courtesy: ESA

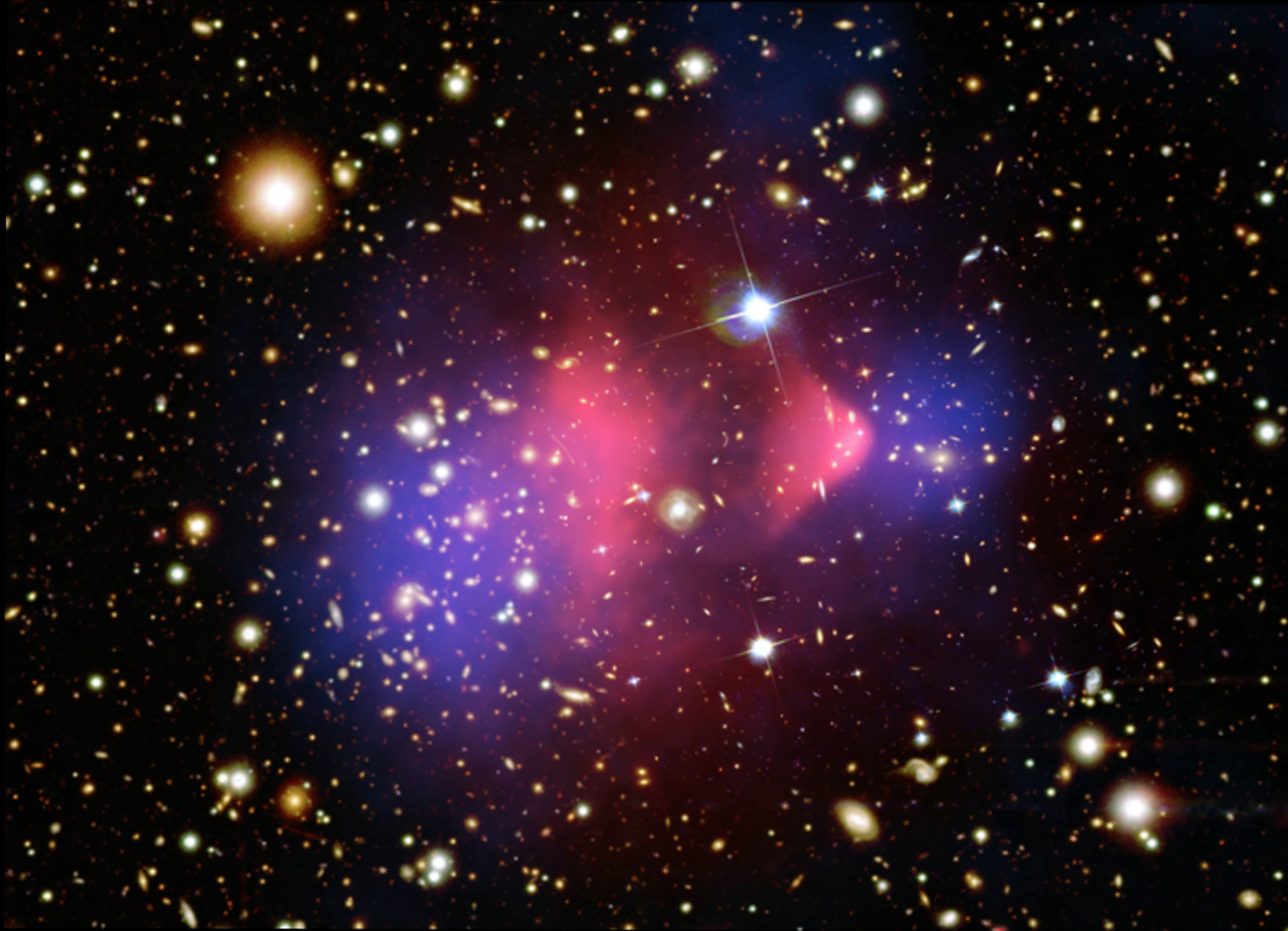


M33 image:
T.A.Rector (NRAO/AUI/NSF and NOAO/AURA/NSF)
M.Hanna (NOAO/AURA/NSF)

[credit: Harvard-Smithsonian Center for Astrophysics]

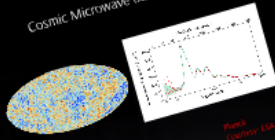
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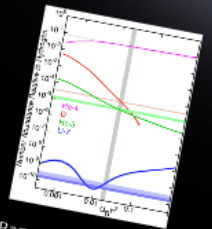
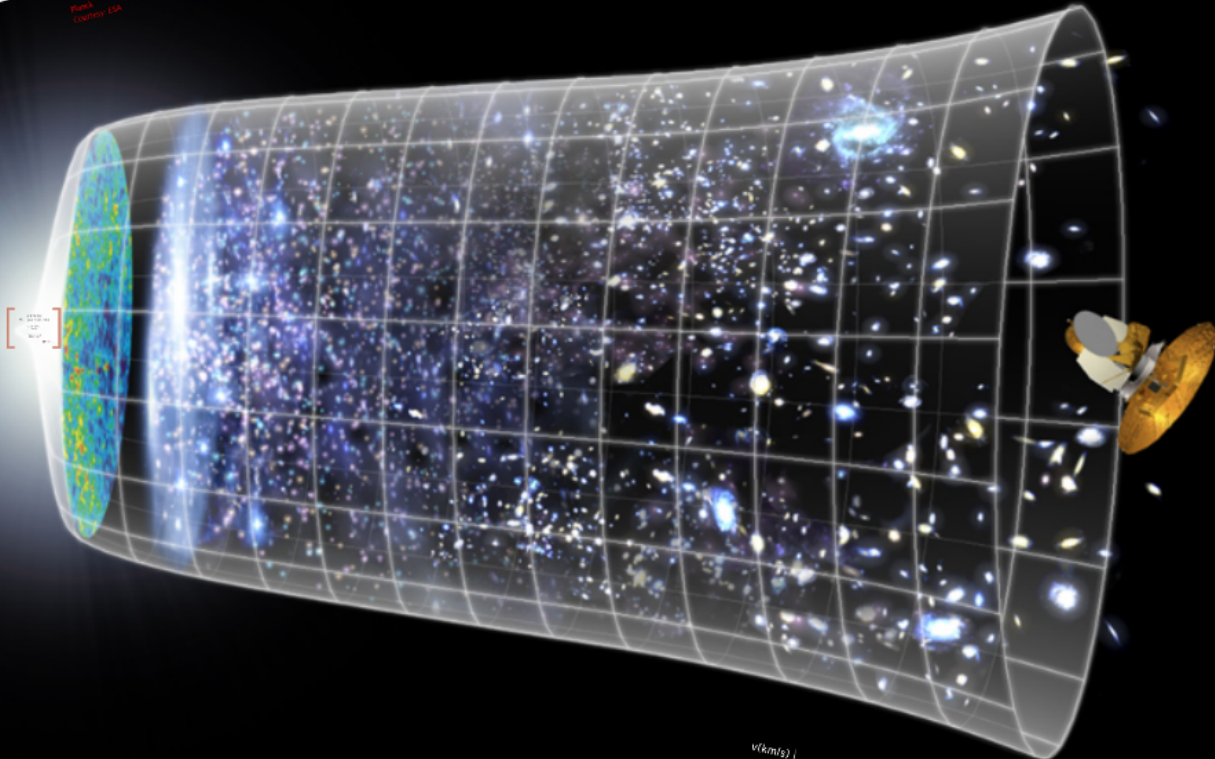


[Credit: X-ray: NASA/CXC/CfA/M.Markevitch et al.;
Optical: NASA/STScI; Magellan/U.Arizona/D.Clowe et al.]

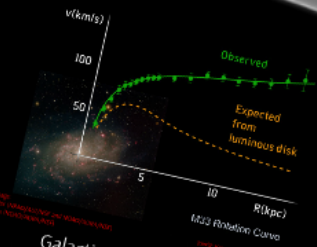
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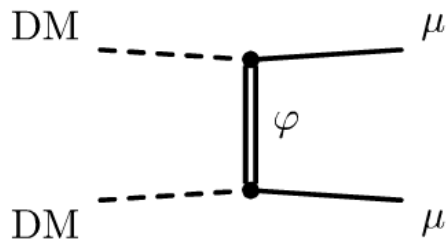
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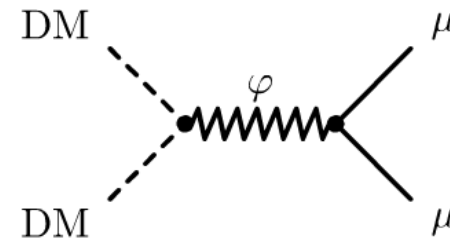
Galactic Rotation Curves

Leptophilic dark matter

Tree-level interactions with SM leptons



Charged mediators

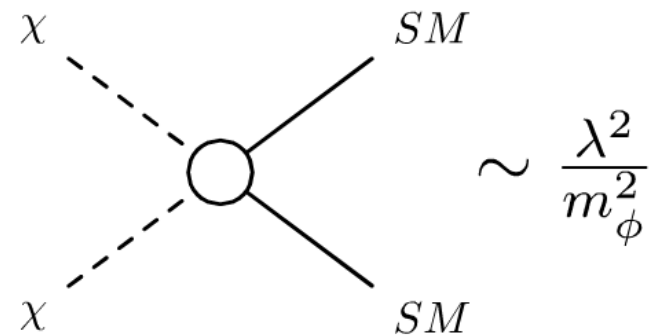


Neutral mediators

Interactions with quarks at one-loop level

Why Leptophilic?

Tension between the WIMP miracle and direct detection signatures



$$\langle \sigma_{Av} \rangle \sim \frac{\lambda^4 m_\chi^2}{32\pi m_\phi^4} \sim \frac{1}{2} \frac{(1.4)^4 (100 \text{ GeV})^2}{32\pi (500 \text{ GeV})^4} \sim 3 \times 10^{-26} \text{ cm}^3/\text{s}$$

$$\sigma^{(n)} \sim \frac{\lambda^4 m_n^2}{64\pi m_\phi^4} \sim \frac{(1.4)^4 (1 \text{ GeV})^2}{64\pi (500 \text{ GeV})^4} \sim 10^{-40} \text{ cm}^2$$

LUX limits on WIMP-nucleon cross section $\sim \text{zb}$ [10^{-45} cm^2]

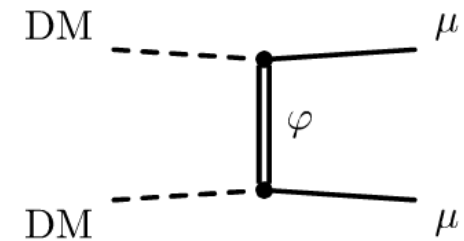
Leptophilic dark matter is loop suppressed for direct detection

Promising for upcoming direct detection experiments

Which Leptophilic?

Charged Mediators

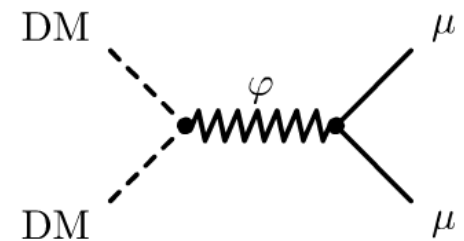
Couplings to left-handed leptons require additional structure



We assume charged mediator couplings to right-handed leptons only

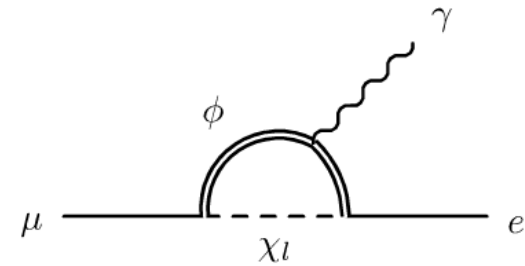
Neutral Mediators

Do not necessitate additional particles charged under $SU(2)_W$



Which Leptophilic?

Flavor constraints restrict the couplings of charged mediators



Two broad possibilities

Flavored mediator - sleptons

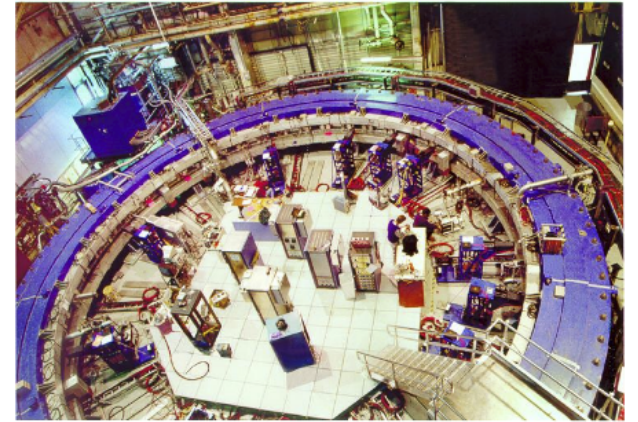
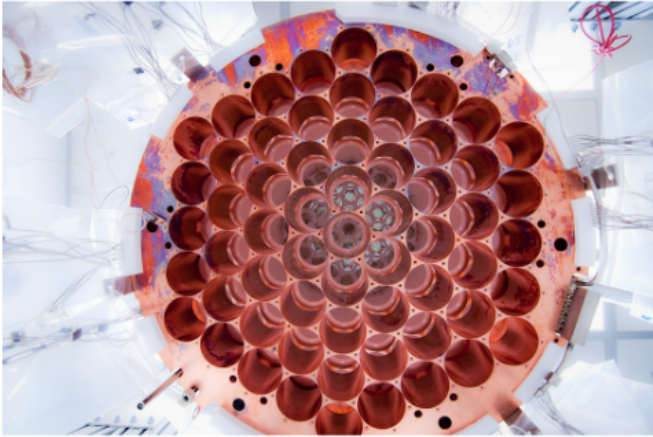
Flavored dark matter - sneutrino, KK neutrino

Couplings of neutral mediators can have more general flavor structure (e.g. gauged $U(1)_{\mu-\tau}$)

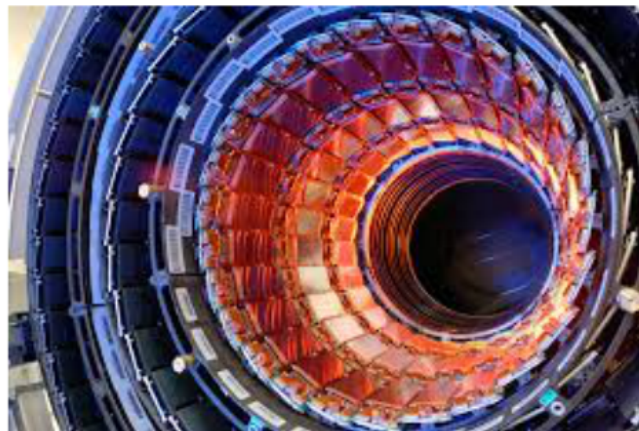
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Direct

g-2

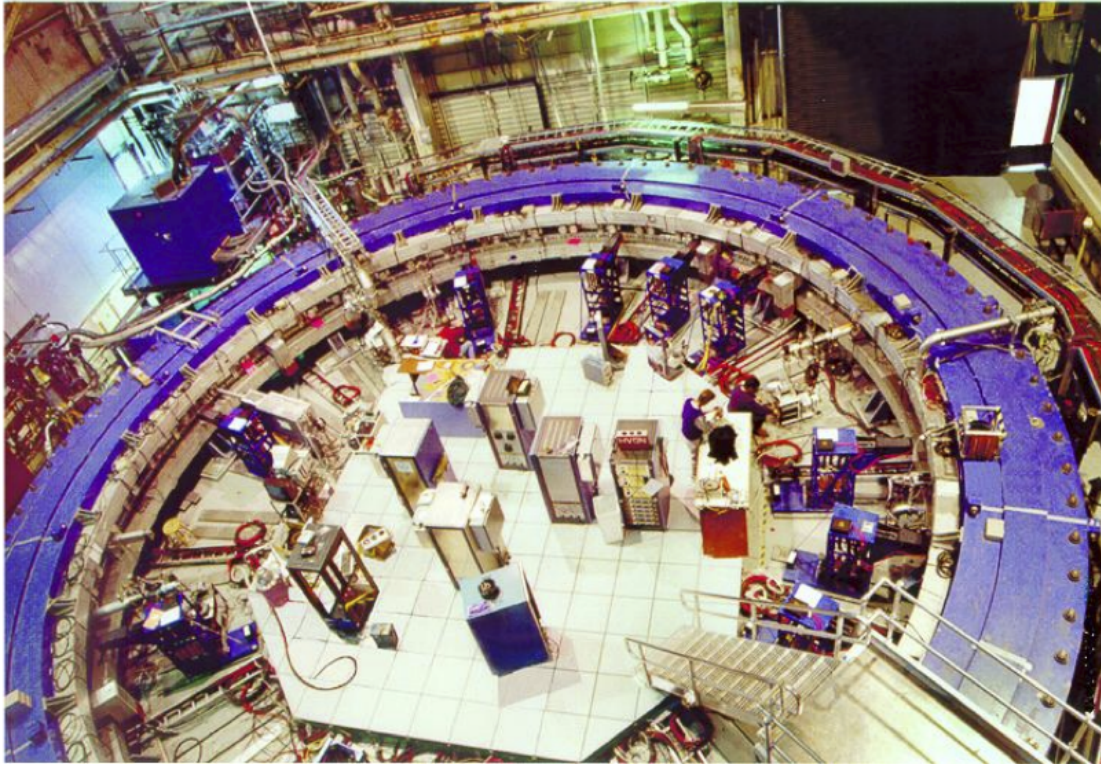


Detection



Colliders

g-2



g - 2

$$a_\mu = \frac{g-2}{2}$$

$$\Delta a_\mu \equiv a_\mu^{\text{exp}} - a_\mu^{\text{SM}} = 28.7(6.3)(4.9) \times 10^{-10}$$

$$\delta a_\mu \equiv 8.0 \times 10^{-10}$$

experimental
theoretical



Hadronic uncertainty

g - 2

Finite effect, cannot capture by an effective field theory

$$\bar{\psi} \sigma^{\mu\nu} \mu F_{\nu\omega}$$

Chiral symmetry breaking effect

$$\hat{a}_\mu \approx \left[\frac{m_\mu^2}{\Lambda^2} \right] \frac{\delta m_\mu}{m_\mu} \quad \frac{\delta m_\mu}{\Lambda^2} \bar{\psi} \sigma^{\mu\nu} \mu F_{\nu\omega} \rightarrow \delta m_\mu \bar{\psi} \mu$$

[Czarnecki, Marciano 2001]

Electron g-2 is a less sensitive probe of new physics

$$\frac{\delta a_e / \delta \delta_e}{\delta a_\mu / \delta \delta_\mu} \sim \frac{\delta_e}{\delta_\mu} \times \frac{m_\mu^2}{m_e^2} \sim 0.02$$

g - 2

$$\mathcal{L}_{\text{int}} = \bar{\psi} (a + b\gamma^5) F S + \text{h.c.} \quad r \equiv m_{\text{SM}}/m_{\text{SM}} \quad \tau \equiv m_\mu/m_{\text{SM}}$$

Charged scalar mediator

$$a_\mu^{CS} = -\frac{\tau}{16\pi^2} \left[\frac{\tau (|a|^2 + |b|^2)}{3\tau^2 - \tau^2} [1 - 6r^2 + 3r^4 + 2r^6 - 6r^4 \ln(r^2)] + \frac{\tau (|a|^2 - |b|^2)}{(1-r^2)^2} [-r^2 + 2r^2 \ln(r^2)] \right] + \mathcal{O}(\tau^6)$$



Charged fermion mediator

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g - 2

Neutral vector mediator

$$\mathcal{L}_{\text{int}} = \bar{\psi} \gamma^\mu (a + b\gamma^5) \mu V_\mu + \text{h.c.} \quad z \equiv m_\mu/m_{\text{SM}}$$

$$a_\mu^N = \frac{z^2}{4\pi^2} \left[\frac{2}{3} (|a|^2 + |b|^2) + (|a|^2 - |b|^2) \right] + \mathcal{O}(z^6)$$

Independent of dark matter couplings

For $m_V \ll 100$ MeV

$$(g-2)_\mu \rightarrow \text{const.}$$

$(g-2)_e$ constraints start being relevant



$g - 2$

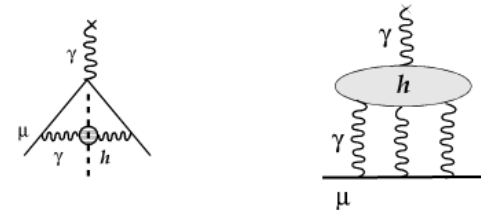
$$a_\mu = \frac{g - 2}{2}$$

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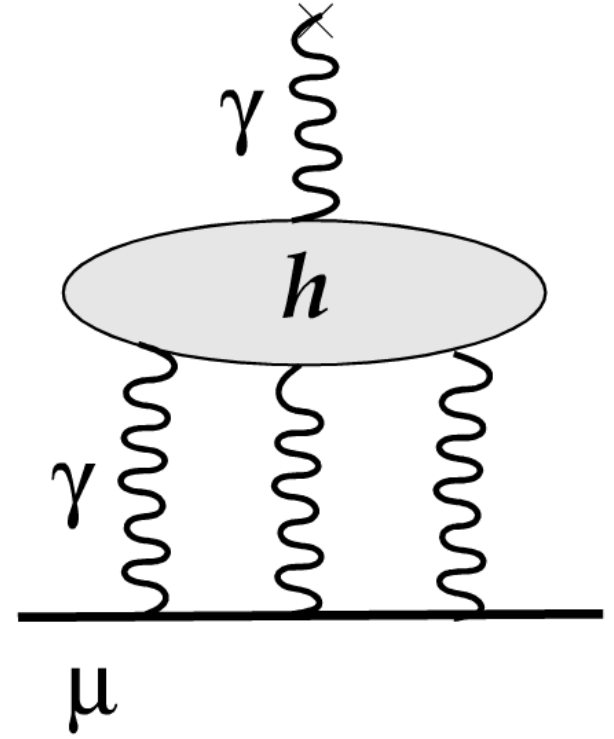
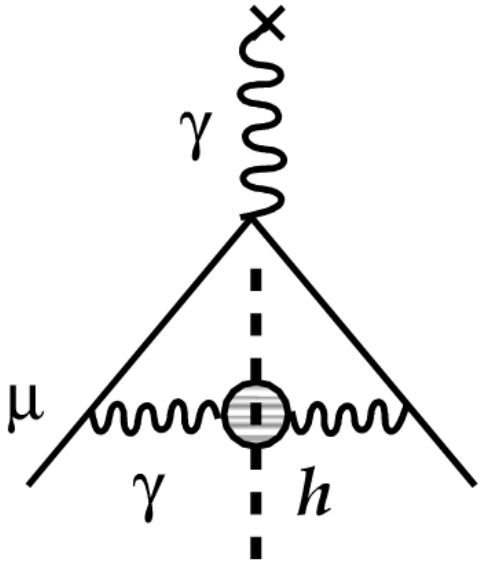
experimental

theoretical

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Hadronic uncertainty



Hadronic uncertainty

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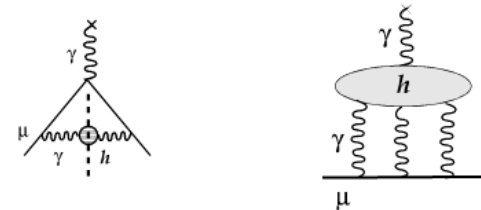
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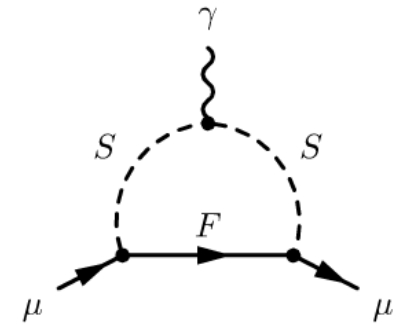
g - 2

$$\mathcal{L}_{\text{scalar}} = \bar{\mu} (a + b\gamma^5) FS + \text{h.c.} \quad r \equiv m_{\text{DM}}/m_{\text{Med}}$$

$$\varepsilon \equiv m_{\mu}/m_{\text{Med}}$$

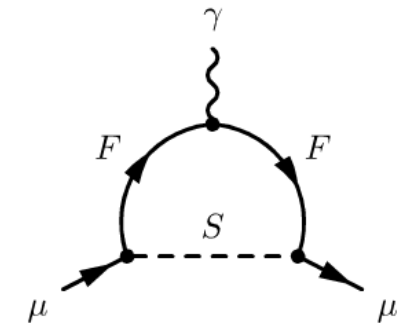
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$g - 2$

Neutral vector mediator

$$\mathcal{L}_{\text{vector}} = \bar{\mu} \gamma^\nu (a + b\gamma^5) \mu V_\nu + \text{h.c.} \quad \varepsilon \equiv m_\mu / m_{\text{Med}}$$

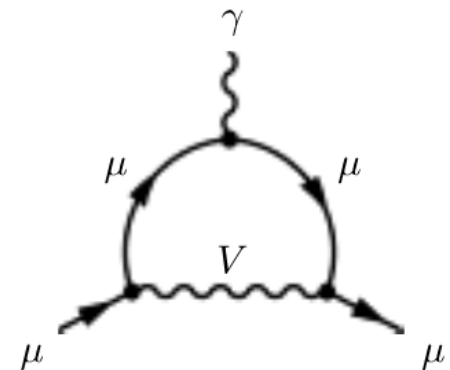
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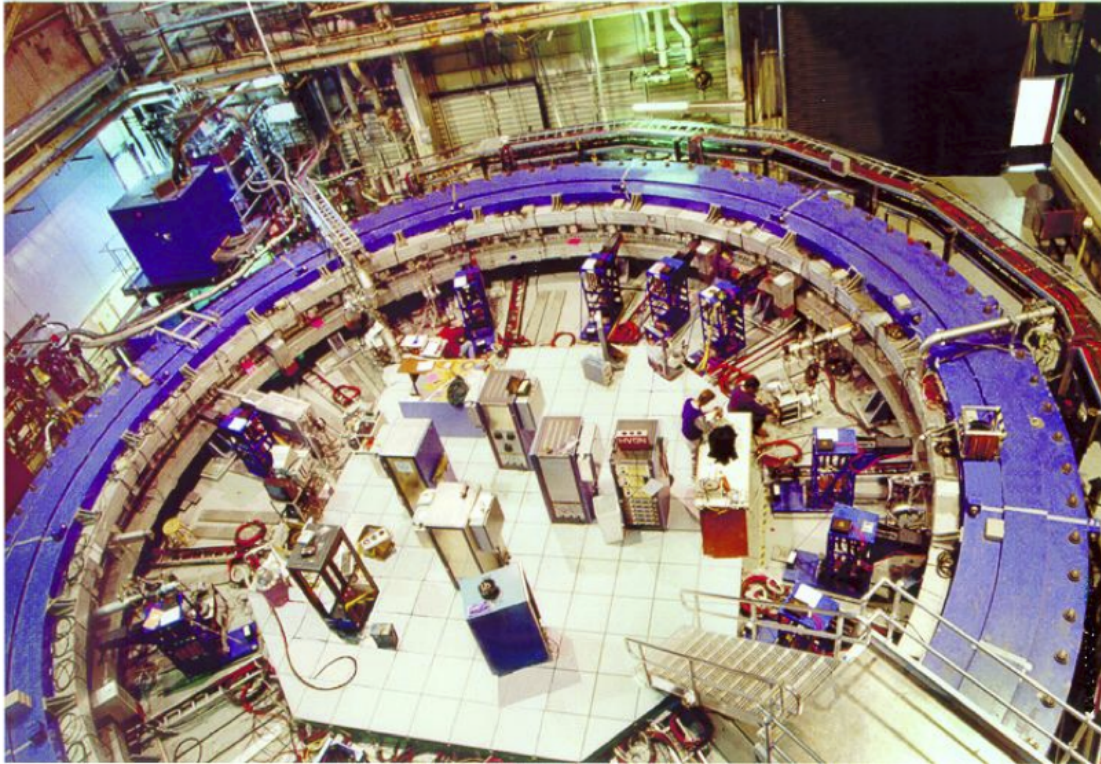
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g - 2

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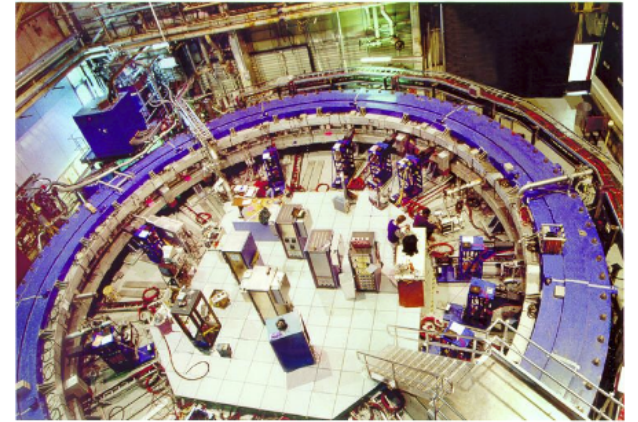
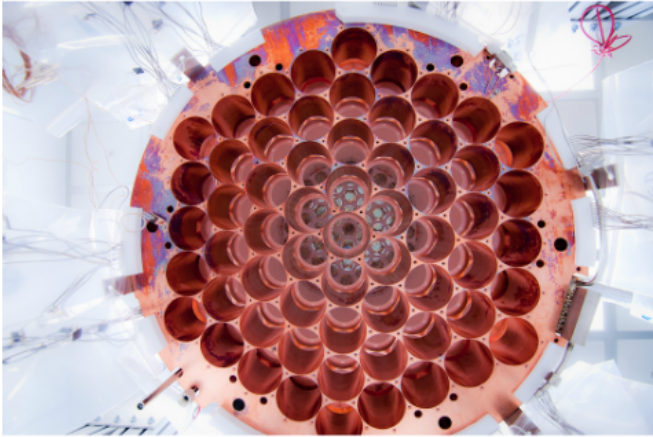
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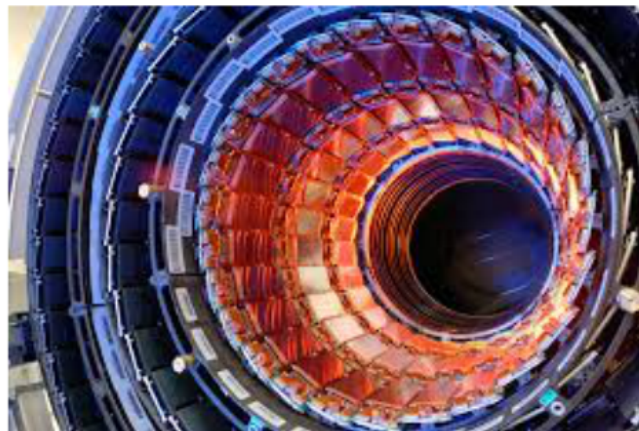


Direct

g-2



Detection



Colliders



Colliders

Collider constraints
Monophoton

$e^+e^- \rightarrow \chi\chi\gamma$

Model independent EFT approach

Not suitable for light neutral mediators

Limits exist for fermionic dark matter

Need simplified models for LHC bounds

(Fox, Hamik, Kopp, Pal 2011)

Collider constraints
Charged Mediator

Two leptons + MET

LEP direct production

LHC limits on right-handed sleptons

Fermionic mediators have much larger production cross section

(Diehl, Haber 2012)

Collider constraints
Neutral mediator

Compositeness bounds

$$\mathcal{L} = \frac{4\pi}{(1+\lambda)^2} [a_{11}e^+e^+e^+e^- + a_{21}e^+e^+e^+e^- + a_{31}e^+e^+e^+e^- + a_{41}e^+e^+e^+e^-]$$

Only valid for $m_{Med} > 208$ GeV

Bound	Operators	Limit
	$W.C. \times \text{Mass}^2 \times \text{Unit} \times \Lambda^2$ (TeV)	
VV	-1 -1 -1 -1	20.0
AA	-1 -1 1 1	18.1
LL+RL	0 0 -1 -1	14.5

Collider constraints
Light Neutral mediator

Resonant production at LEP

- Exclusions in the context of specific models
- Excludes coupling strengths of $O(0.01)$

BaBar

$e^+e^- \rightarrow e\gamma$
 $0 \rightarrow \mu^+\mu^-$

Beam dumps

Supernovae cooling

(Bjorken, Chung, Schrock, Tan 2010)

Collider constraints

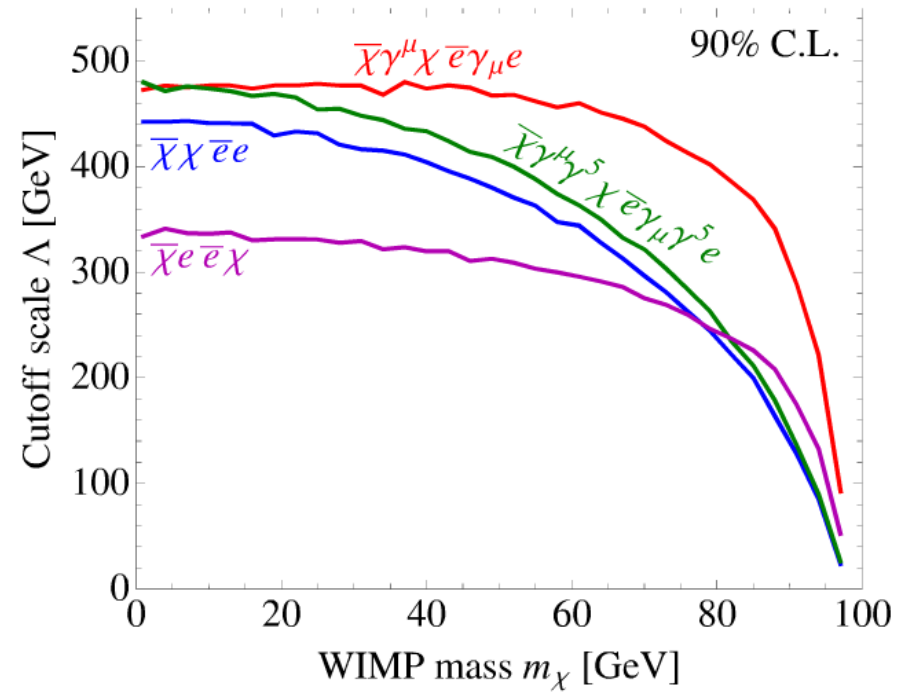
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[Fox, Harnik, Kopp, Tsai 2011]

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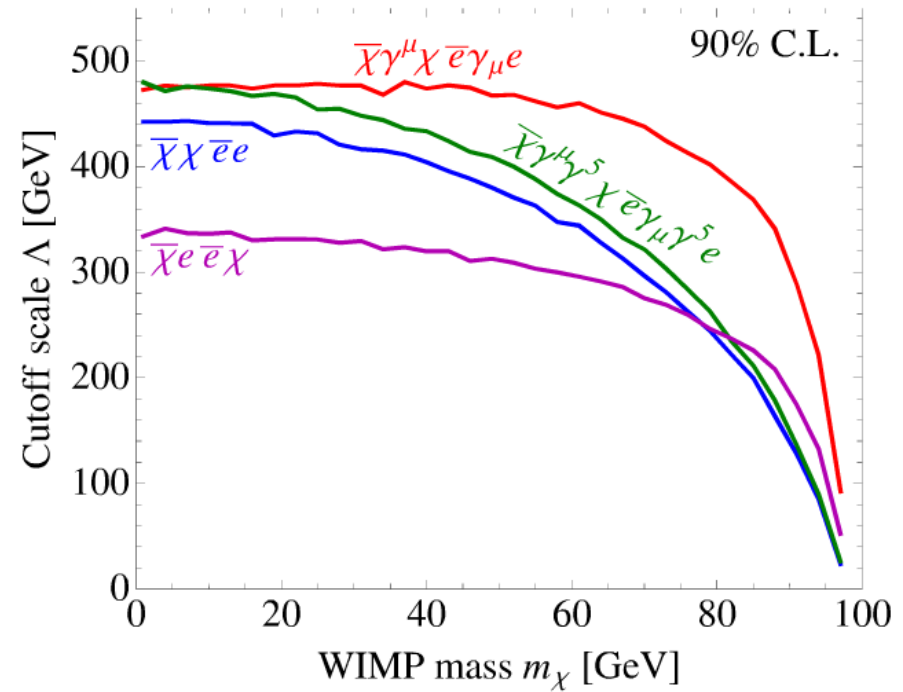
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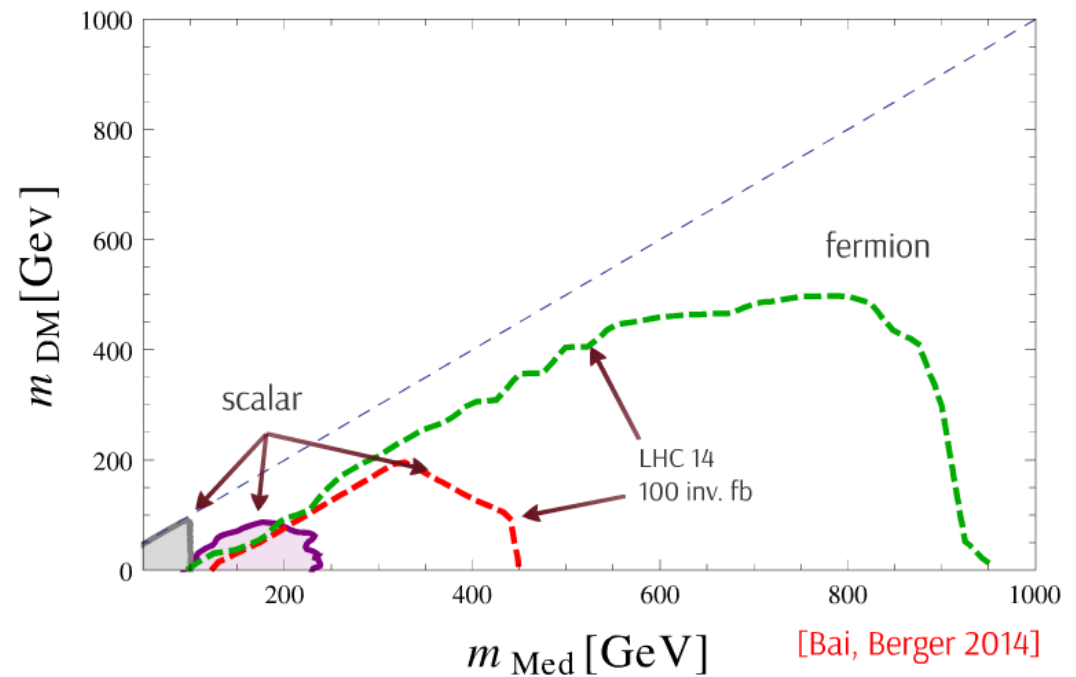
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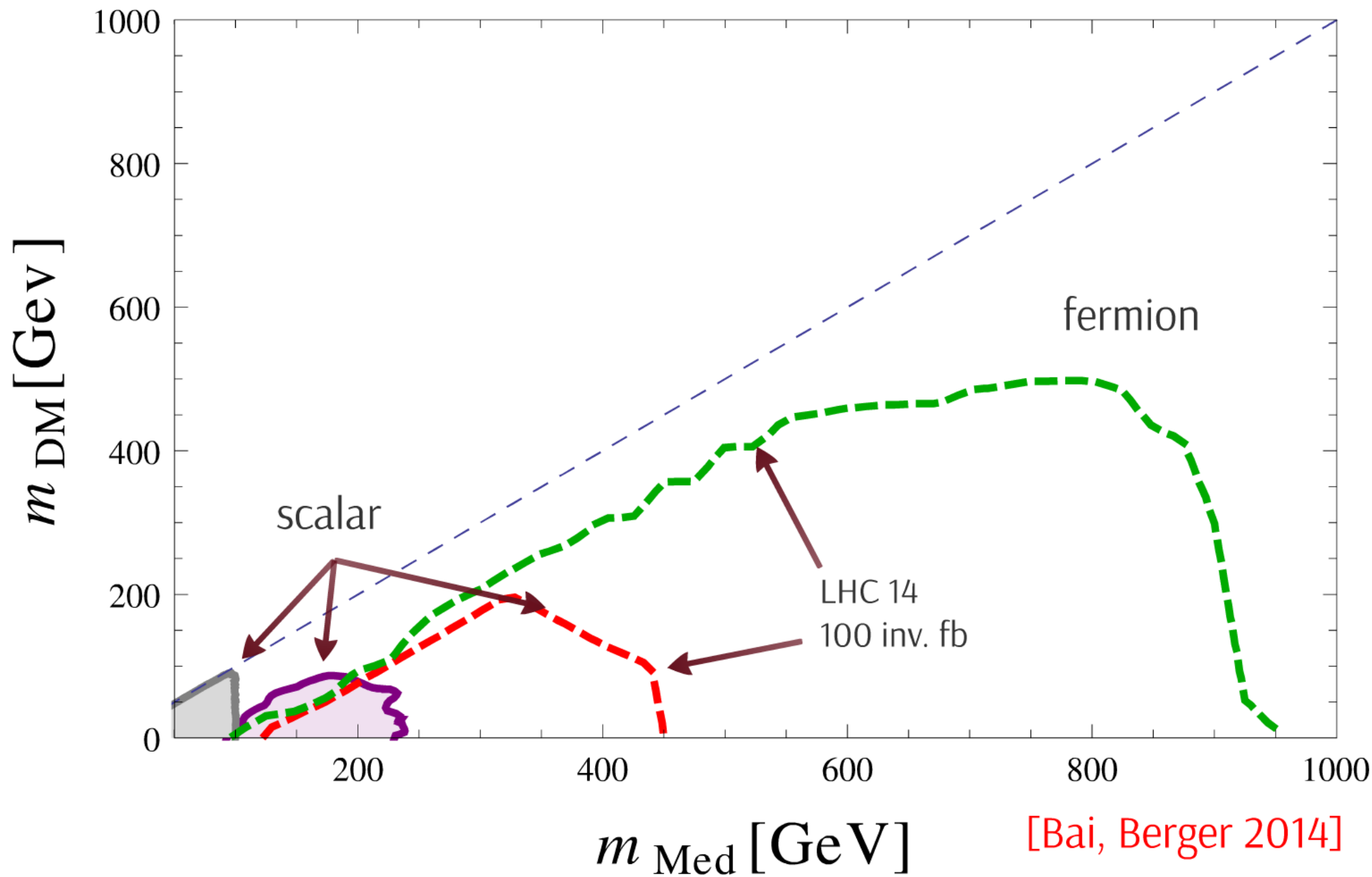
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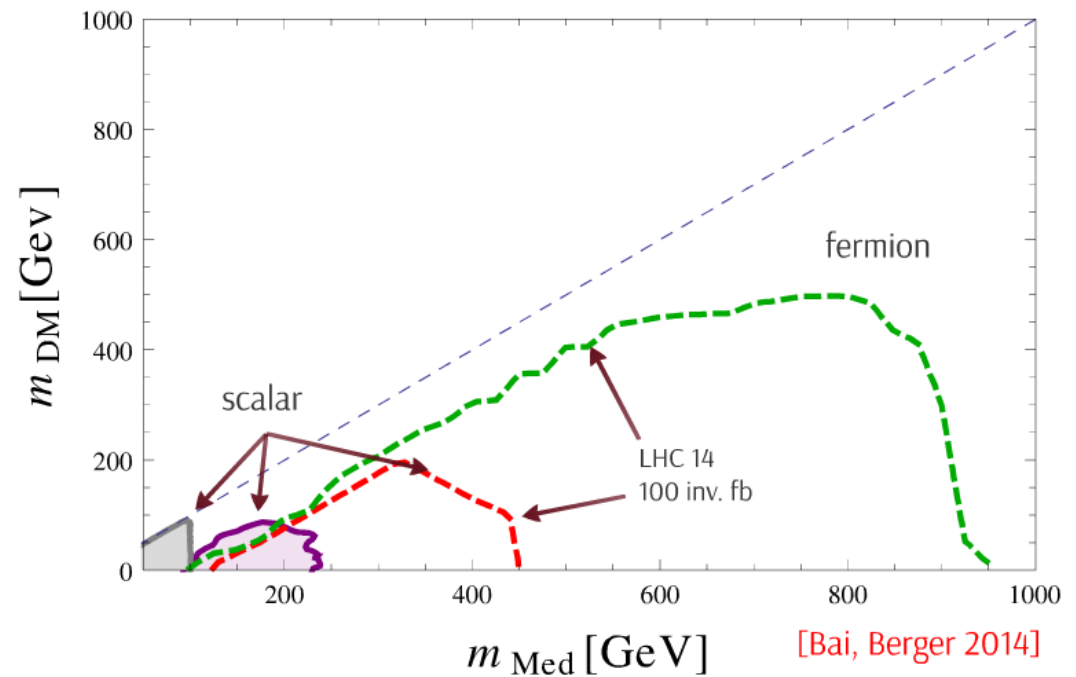
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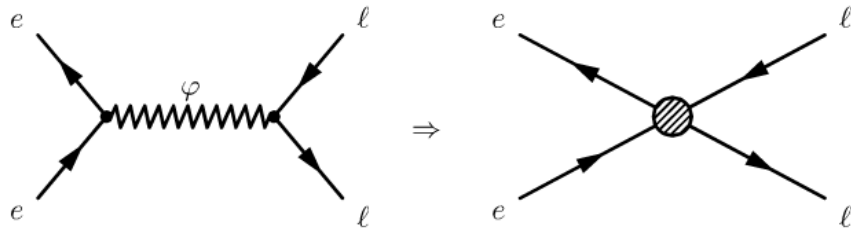
Fermionic mediators have much larger production cross section

Collider constraints

Neutral mediator

Compositeness bounds

$$\mathcal{L} = \frac{4\pi}{(1 + \delta)\Lambda^2} [\eta_{LL}\bar{e}_L\gamma^\mu e_L\bar{\ell}_L\gamma_\mu\ell_L + \eta_{RR}\bar{e}_R\gamma^\mu e_R\bar{\ell}_R\gamma_\mu\ell_R + \eta_{LR}\bar{e}_L\gamma^\mu e_L\bar{\ell}_R\gamma_\mu\ell_R + \eta_{RL}\bar{e}_R\gamma^\mu e_R\bar{\ell}_L\gamma_\mu\ell_L]$$



Only valid for $m_{Med} > 208$ GeV

Bound	Operators				Limit
	η_{LL}	η_{RR}	η_{LR}	η_{RL}	Λ (TeV)
VV	-1	-1	-1	-1	20.0
AA	-1	-1	1	1	18.1
LR+RL	0	0	-1	-1	14.5

Collider constraints

Light Neutral mediator

Resonant production at LEP

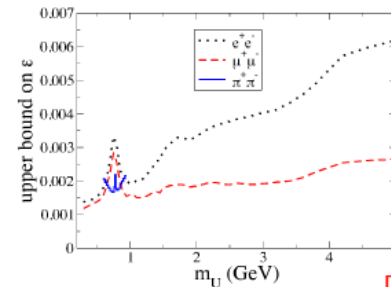
- Exclusions in the context of specific models
- Excludes coupling strengths of $O(0.01)$

BaBar

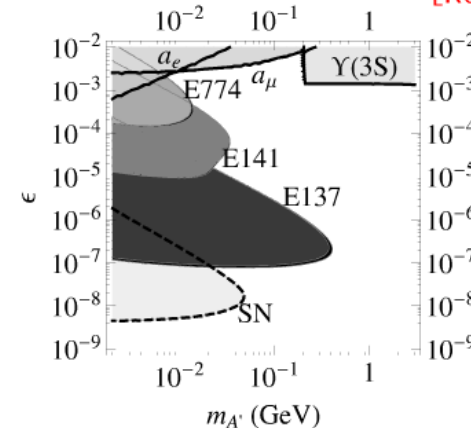
$$e^+e^- \rightarrow \phi\gamma$$
$$\phi \rightarrow \mu^+\mu^-$$

Beam dumps

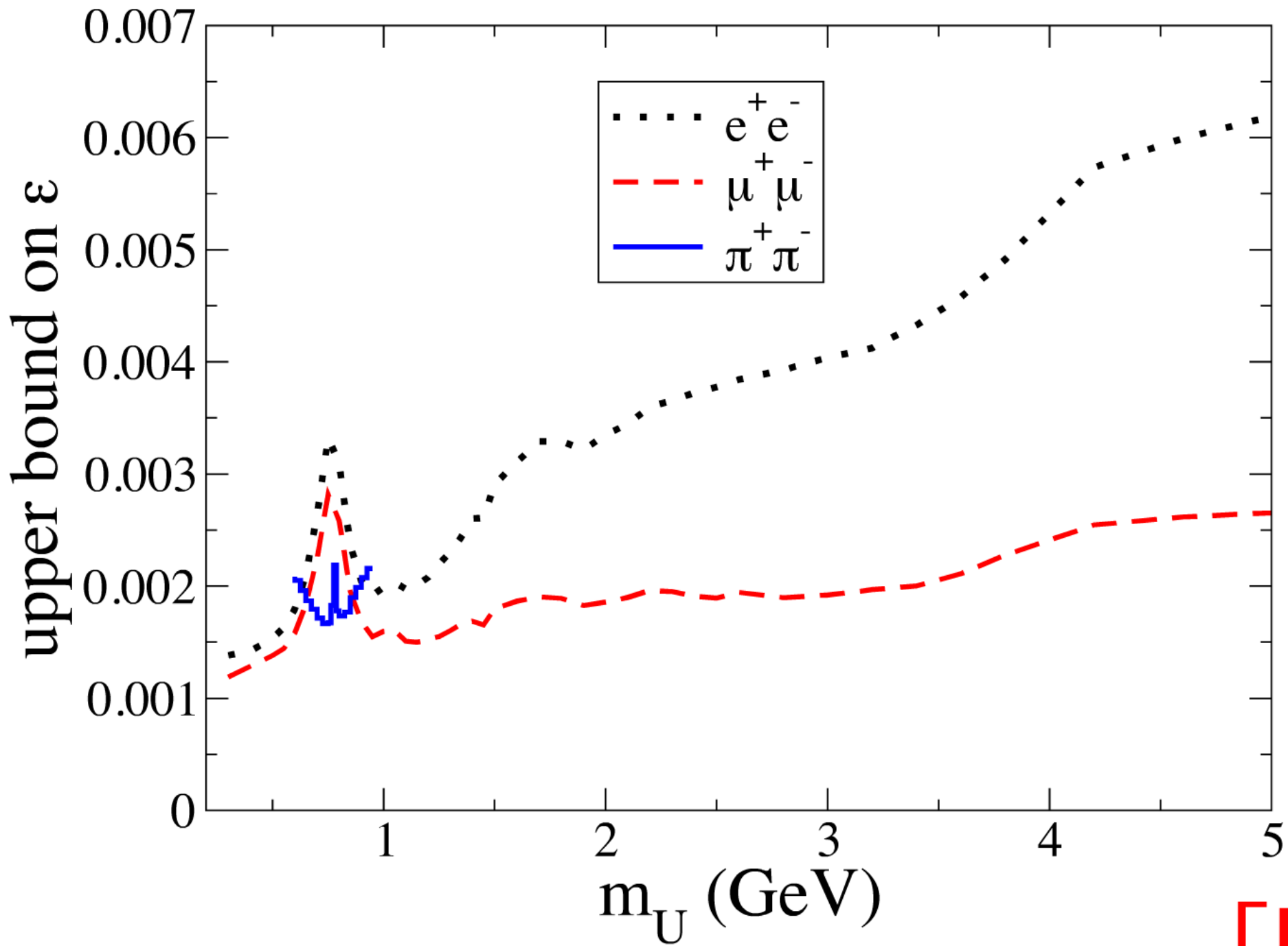
Supernovae cooling



[Reece, Wang 2010]



[Bjorken, Essig, Schuster, Toro 2009]



[Re

Collider constraints

Light Neutral mediator

Resonant production at LEP

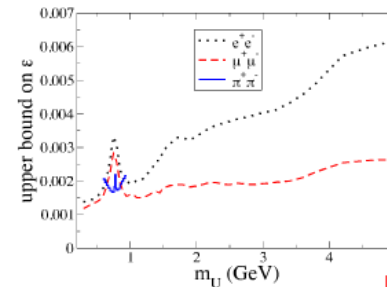
- Exclusions in the context of specific models
- Excludes coupling strengths of $O(0.01)$

BaBar

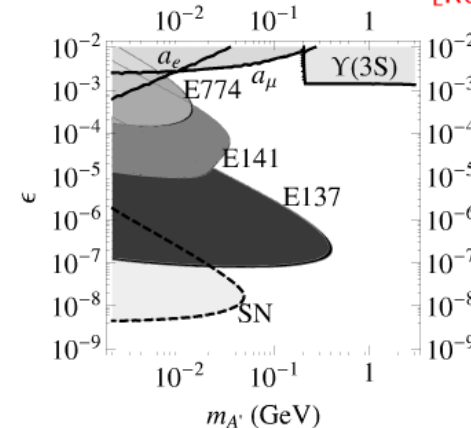
$$e^+e^- \rightarrow \phi\gamma$$
$$\phi \rightarrow \mu^+\mu^-$$

Beam dumps

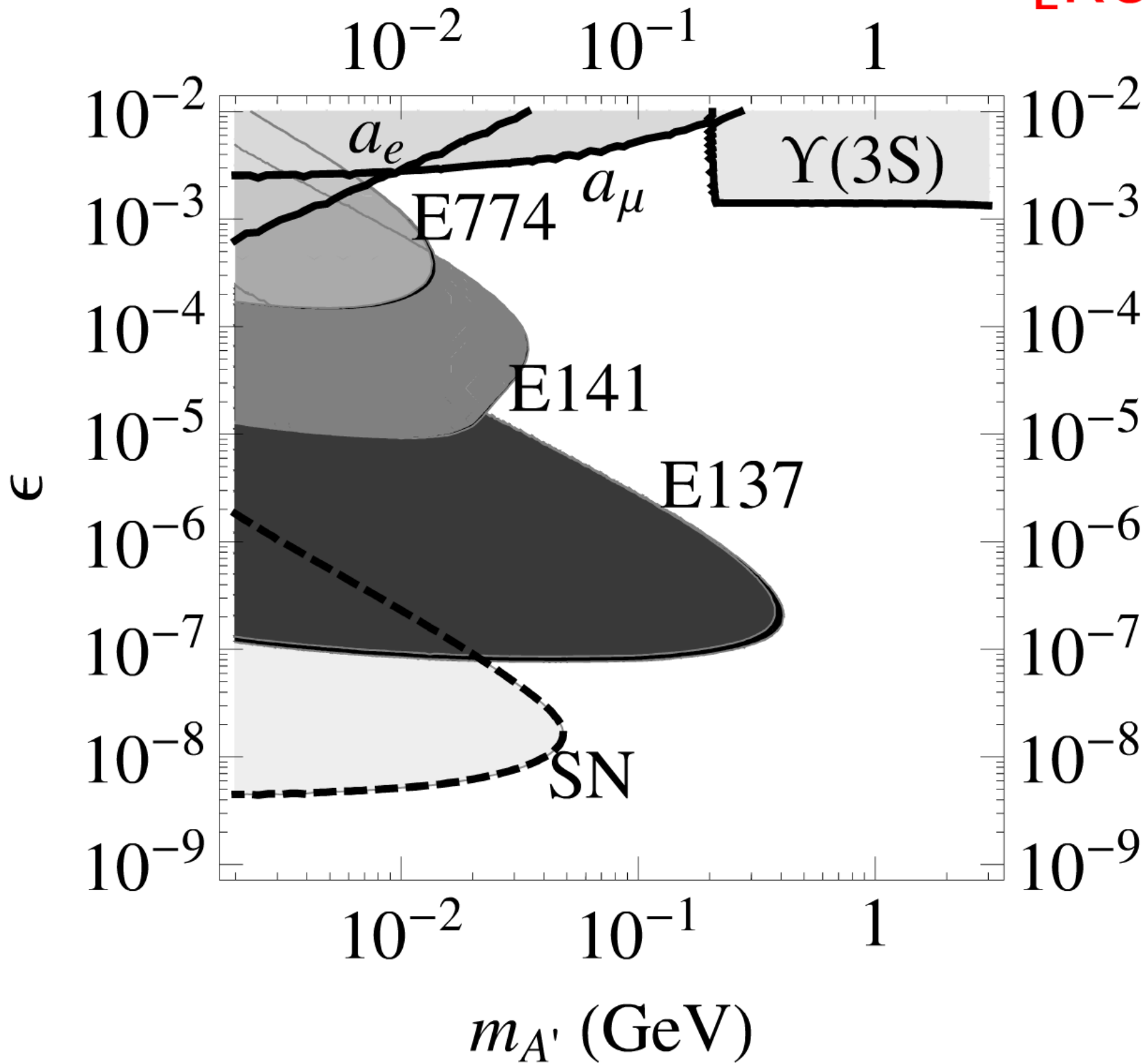
Supernovae cooling



[Reece, Wang 2010]



[Bjorken, Essig, Schuster, Toro 2009]



Collider constraints

Light Neutral mediator

Resonant production at LEP

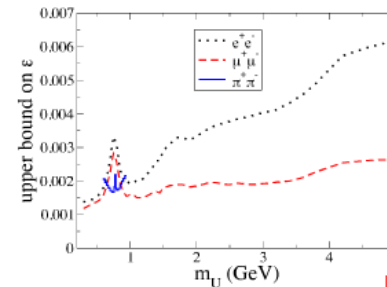
- Exclusions in the context of specific models
- Excludes coupling strengths of $O(0.01)$

BaBar

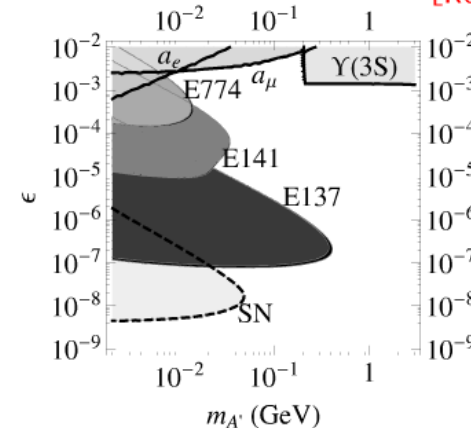
$$e^+e^- \rightarrow \phi\gamma$$
$$\phi \rightarrow \mu^+\mu^-$$

Beam dumps

Supernovae cooling



[Reece, Wang 2010]



[Bjorken, Essig, Schuster, Toro 2009]



Colliders

Collider constraints
Monophoton

$e^+e^- \rightarrow \chi\chi\gamma$

Model independent EFT approach

Not suitable for light neutral mediators

Limits exist for fermionic dark matter

Need simplified models for LHC bounds

(Fox, Hamik, Kopp, Pal 2011)

Collider constraints
Charged Mediator

Two leptons + MET

LEP direct production

LHC limits on right-handed sleptons

Fermionic mediators have much larger production cross section

(Diehl, Haber 2012)

Collider constraints
Neutral mediator

Compositeness bounds

$$\mathcal{L} = \frac{4\pi}{(1+\lambda)^2} [a_{11}\psi_1^* \psi_1 \psi_1^* \psi_1 + a_{22}\psi_2^* \psi_2 \psi_2^* \psi_2 + a_{33}\psi_3^* \psi_3 \psi_3^* \psi_3 + a_{44}\psi_4^* \psi_4 \psi_4^* \psi_4]$$

Only valid for $m_{Med} > 208$ GeV

Bound	Operators	Limit
	$W.C. \times \text{Mass}^2 \times \text{Unit} \times \Lambda^2$ (TeV)	
VV	-1 -1 -1 -1	20.0
AA	-1 -1 1 1	18.1
LL+RL	0 0 -1 -1	14.5

Collider constraints
Light Neutral mediator

Resonant production at LEP

- Exclusions in the context of specific models
- Excludes coupling strengths of $O(0.01)$

BaBar

$e^+e^- \rightarrow \sigma\gamma$
 $0 \rightarrow \mu^+\mu^-$

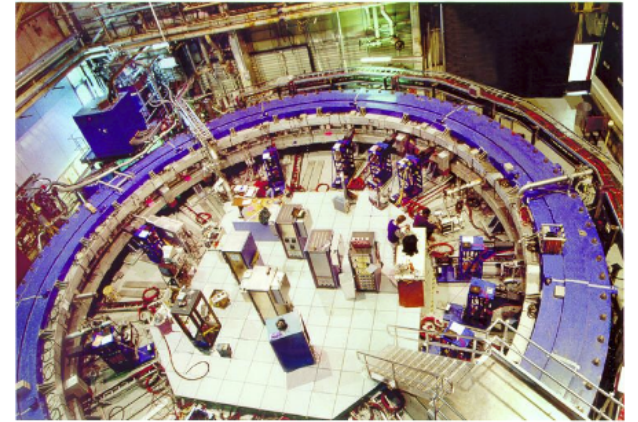
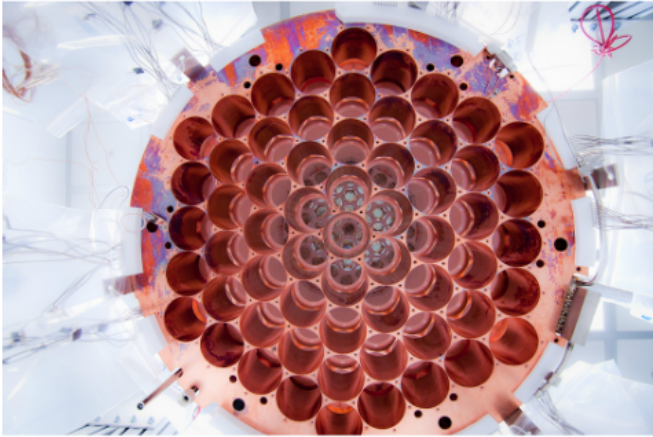
Beam dumps

Supernovae cooling

(Shen, Wang 2010)
(Griest, Goldberg, Schwartz, Tan 2000)

Direct

g-2



Detection



Colliders

Direct

Direct detection

More loop through photon exchange
 Can analyze form factors for coupling to the DM current
 Example Fermion dark matter

$$T^{\mu\nu} = \gamma^{\mu} \gamma^{\nu} f(p') + \frac{1}{2m^2} \epsilon^{\mu\nu\alpha\beta} p_{\alpha} p'_{\beta} + \frac{1}{2m^2} \epsilon^{\mu\nu\alpha\beta} p_{\alpha} p'_{\beta} + \dots$$

charge form factor magnetic dipole moment electric dipole moment anapole moment

Direct detection

Self-conjugate particles
 Real scalar, Majorana fermions, real vector

Example Majorana fermion dark matter

$$T^{\mu\nu} = \frac{1}{2} \gamma^{\mu} \gamma^{\nu} f(p') + \frac{1}{2m^2} \epsilon^{\mu\nu\alpha\beta} p_{\alpha} p'_{\beta} + \dots$$

Only non-zero form factor is P-odd Anapole moment
 Leads to velocity-suppressed scattering amplitudes

$$\chi^{\dagger} \gamma^{\mu} \gamma^{\nu} \chi$$

Direct detection

Least significant only for complex scalar and Dirac fermions

$$\mathcal{M} = \frac{1}{2} \gamma^{\mu} \gamma^{\nu} f(p') + \frac{1}{2m^2} \epsilon^{\mu\nu\alpha\beta} p_{\alpha} p'_{\beta} + \dots$$

$$M = \frac{1}{2} \sum_{\lambda} \langle N_{\lambda} | \mathcal{O} | N_{\lambda} \rangle \langle N_{\lambda} | \mathcal{O} | N_{\lambda} \rangle$$

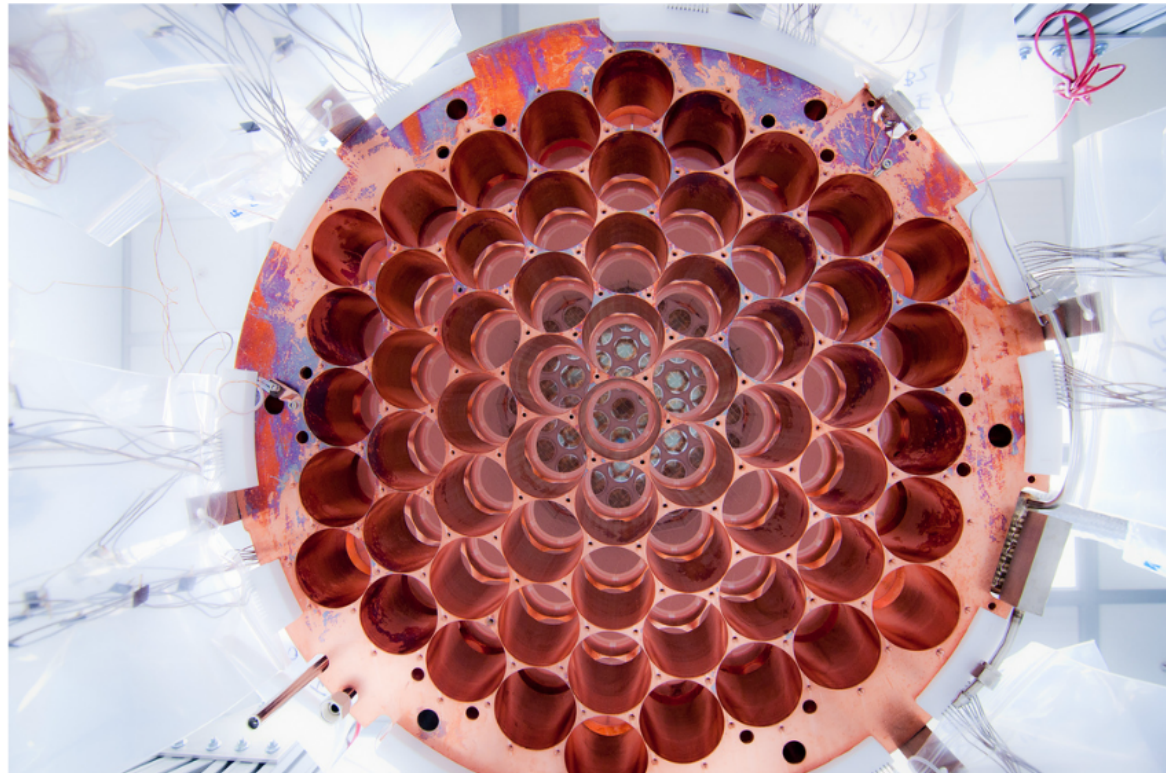
Direct detection

$$M = \frac{1}{2} \sum_{\lambda} \langle N_{\lambda} | \mathcal{O} | N_{\lambda} \rangle \langle N_{\lambda} | \mathcal{O} | N_{\lambda} \rangle$$

$$\langle N_{\lambda} | \mathcal{O} | N_{\lambda} \rangle = \frac{1}{2} \text{Tr} \left[\gamma^{\mu} \gamma^{\nu} f(p') + \frac{1}{2m^2} \epsilon^{\mu\nu\alpha\beta} p_{\alpha} p'_{\beta} + \dots \right]$$

$$\langle N_{\lambda} | \mathcal{O} | N_{\lambda} \rangle = \frac{1}{2} \text{Tr} \left[\gamma^{\mu} \gamma^{\nu} f(p') + \frac{1}{2m^2} \epsilon^{\mu\nu\alpha\beta} p_{\alpha} p'_{\beta} + \dots \right]$$

Leads to velocity-suppressed scattering amplitudes
 Leads to velocity-suppressed scattering amplitudes
 Leads to velocity-suppressed scattering amplitudes



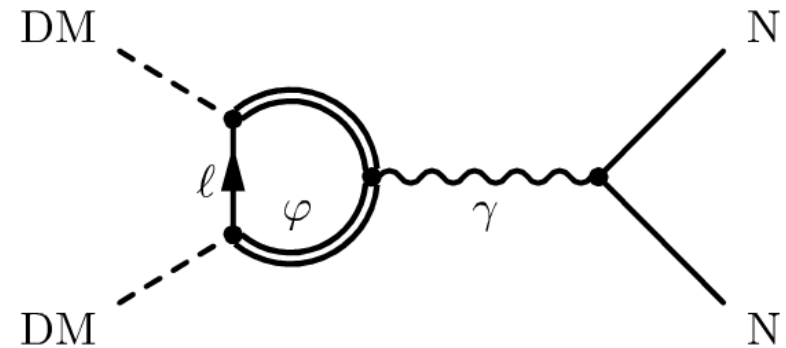
Detection

Direct detection

At one loop through photon exchange

Can analyze form-factors for coupling to the EM current

Example: Fermion dark matter



anapole moment

$$\Gamma^\mu = \gamma^\mu F_1(q^2) + \frac{i\sigma^{\mu\nu}q_\nu}{2m} F_2(q^2) + \frac{i\sigma^{\mu\nu}q_\nu\gamma^5}{2m} F_3(q^2) + (\gamma^\mu q^2 - \not{q}q^\mu)\gamma^5 F_A(q^2)$$

charge form-factor

magnetic dipole moment

electric dipole moment



Direct detection

Self-conjugate particles

Real scalar, Majorana fermions, real vector

Example: Majorana fermion dark matter

$$\Gamma^\mu = \cancel{\gamma^\mu F_1(q^2)} + \cancel{\frac{i\sigma^{\mu\nu}q_\nu}{2m} F_2(q^2)} + \cancel{\frac{i\sigma^{\mu\nu}q_\nu\gamma^5}{2m} F_3(q^2)} + (\gamma^\mu q^2 - \not{q}q^\mu)\gamma^5 F_A(q^2)$$

Only non-zero form factor is P-odd Anapole moment

Leads to velocity-suppressed scattering amplitudes

$$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu q$$

Direct detection

Limit significant only for complex scalar and Dirac fermion

Exception:

Dark matter and mediator are highly degenerate

[Kopp, Michaels, Smirnov 2014]

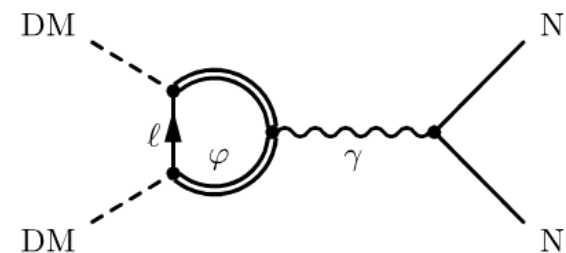
$$\mathcal{O}_1 = [\bar{\chi}\gamma^\mu(c + d\gamma^5)\partial^\nu\chi + \text{h.c.}] F_{\mu\nu},$$

$$\mathcal{O}_2 = [i\bar{\chi}\gamma^\mu(c + d\gamma^5)\partial^\nu\chi + \text{h.c.}] F^{\sigma\rho}\epsilon_{\mu\nu\sigma\rho}$$

$$\Gamma^\mu = \gamma^\mu F_1(q^2) + \frac{i\sigma^{\mu\nu}q_\nu}{2m} F_2(q^2) + \frac{i\sigma^{\mu\nu}q_\nu\gamma^5}{2m} F_3(q^2) + (\gamma^\mu q^2 - \not{q}q^\mu)\gamma^5 F_A(q^2)$$

$$\mathcal{M} = \tilde{\lambda} \sum_q \langle N_f | Q \bar{q} \gamma_\alpha q | N_i \rangle \bar{u}(p_2) \gamma^\alpha u(p_1)$$

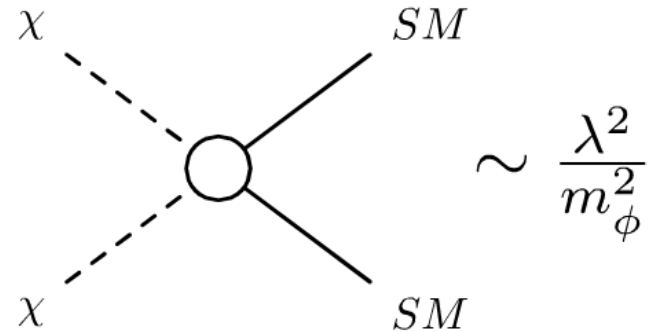
$$\tilde{\lambda} \sim \frac{\alpha}{4\pi} \frac{\lambda^2}{m_\phi^2} \log \left[\frac{m_l^2}{m_\phi^2} \right]$$



Direct detection

$$\mathcal{M} = \tilde{\lambda} \sum_q \langle N_f | Q \bar{q} \gamma_\alpha q | N_i \rangle \bar{u}(p_2) \gamma^\alpha u(p_1)$$

$$\tilde{\lambda} \sim \frac{\alpha}{4\pi} \frac{\lambda^2}{m_\phi^2} \log \left[\frac{m_l^2}{m_\phi^2} \right]$$



$$\langle \sigma_{Av} \rangle \sim \frac{\lambda^4 m_\chi^2}{32\pi m_\phi^4} \sim \frac{1}{2} \frac{(1.4)^4 (100 \text{ GeV})^2}{32\pi (500 \text{ GeV})^4} \sim 3 \times 10^{-26} \text{ cm}^3/\text{s}$$

$$\sigma^{(n)} \sim \frac{\tilde{\lambda}^2 m_n^2}{4\pi} \sim \left(\frac{\alpha}{4\pi} \log \left[\frac{m_l^2}{m_\phi^2} \right] \right)^2 \frac{(1.4)^4 (1 \text{ GeV})^2}{64\pi (500 \text{ GeV})^4} \sim 10^{-45} \text{ cm}^2$$

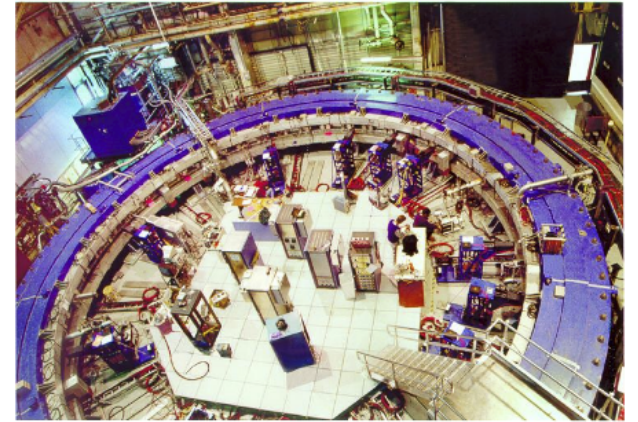
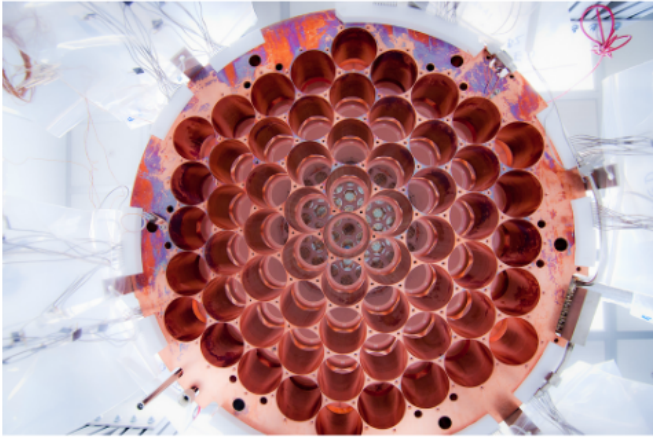
LUX limits on WIMP-nucleon cross section $\sim \text{zb}$ [10^{-45} cm^2]

Leptophilic dark matter is loop suppressed for direct detection

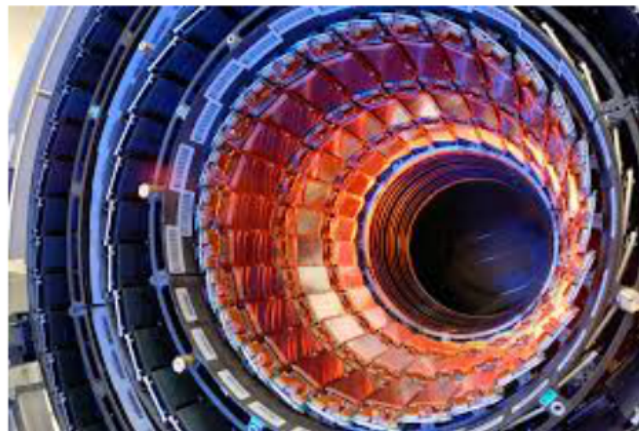
Promising for upcoming direct detection experiments

Direct

g-2



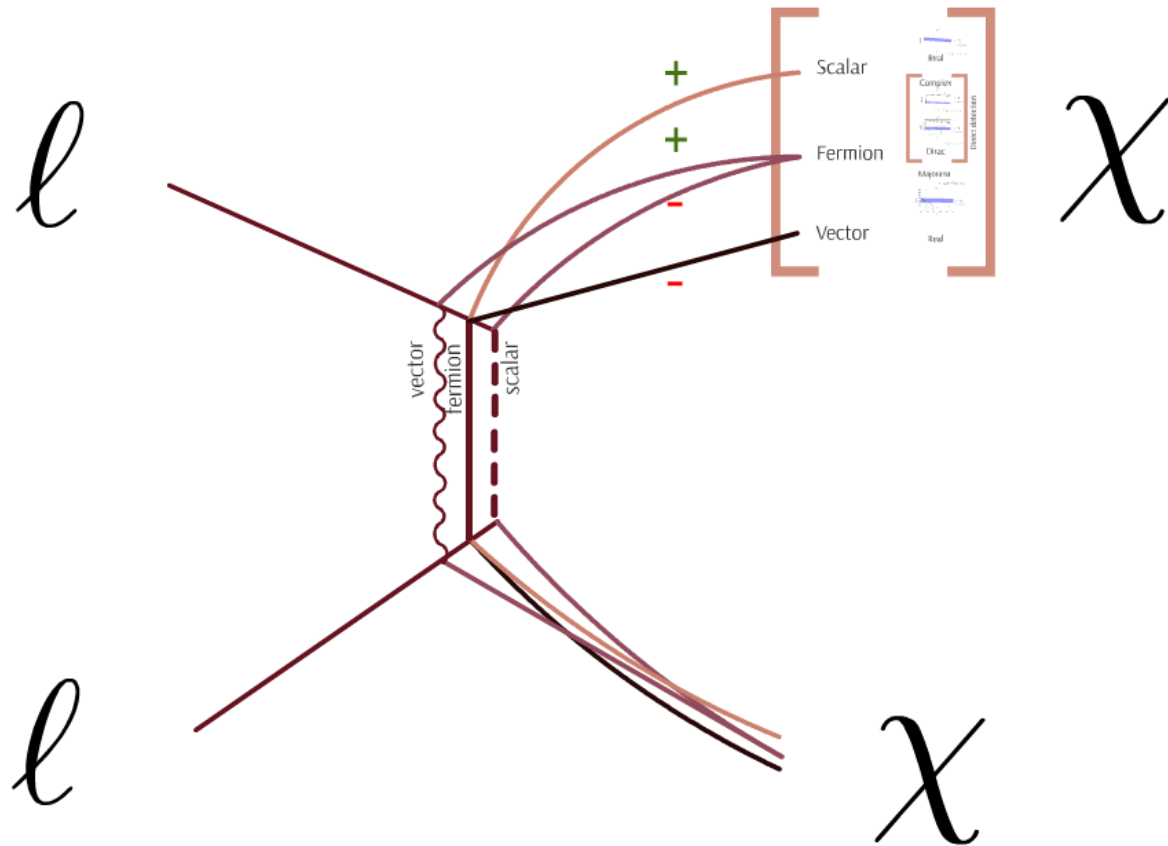
Detection

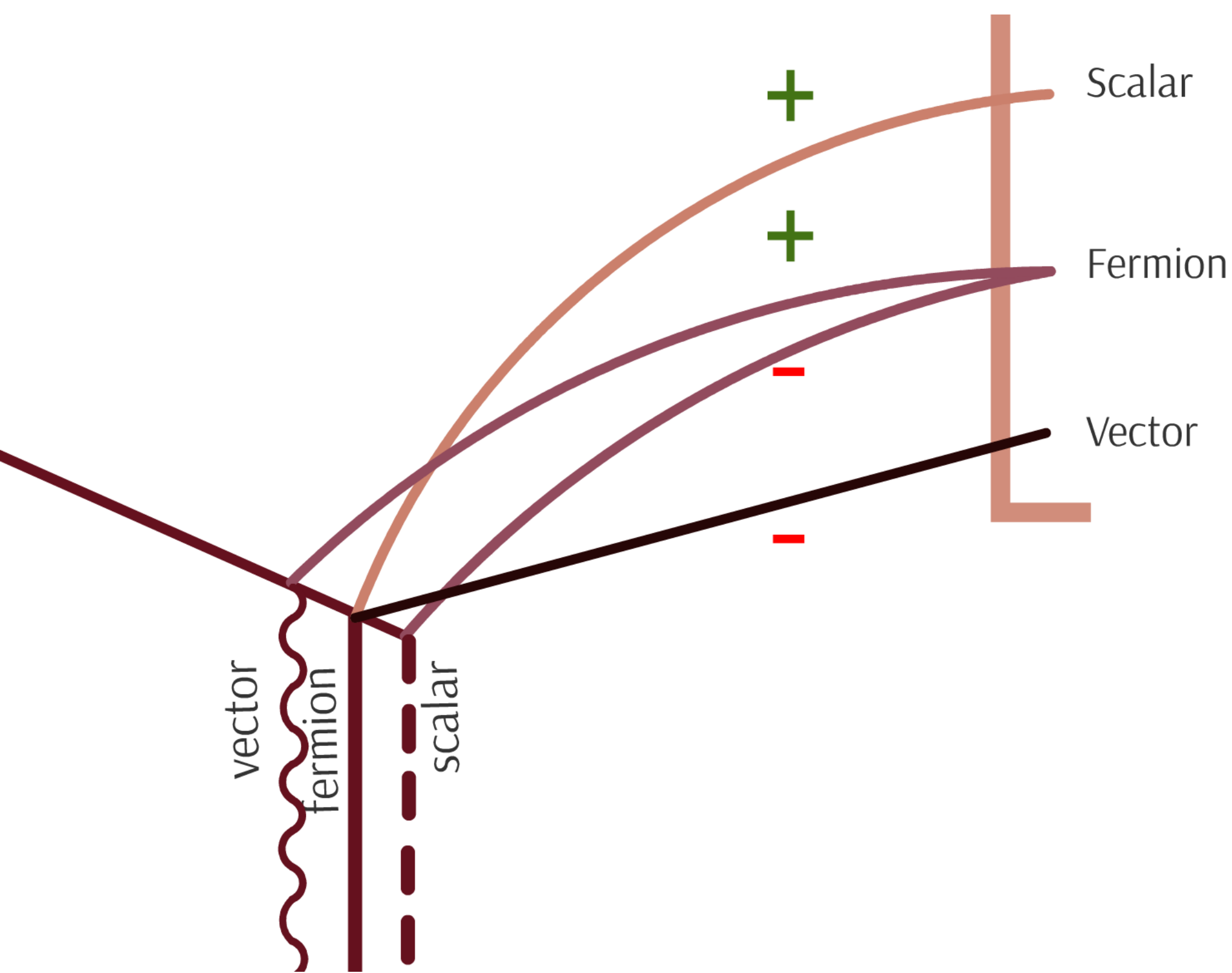


Colliders

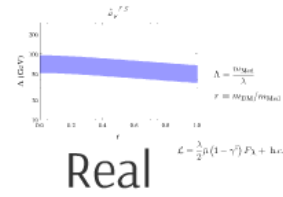
Simplified Models

Charged Mediator

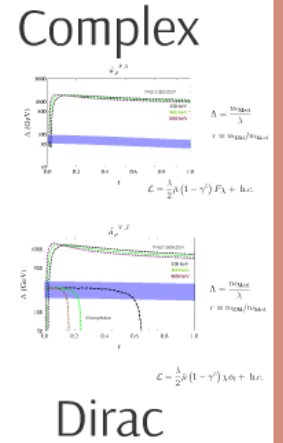




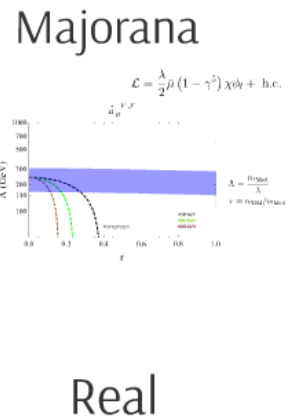
Scalar



Fermion

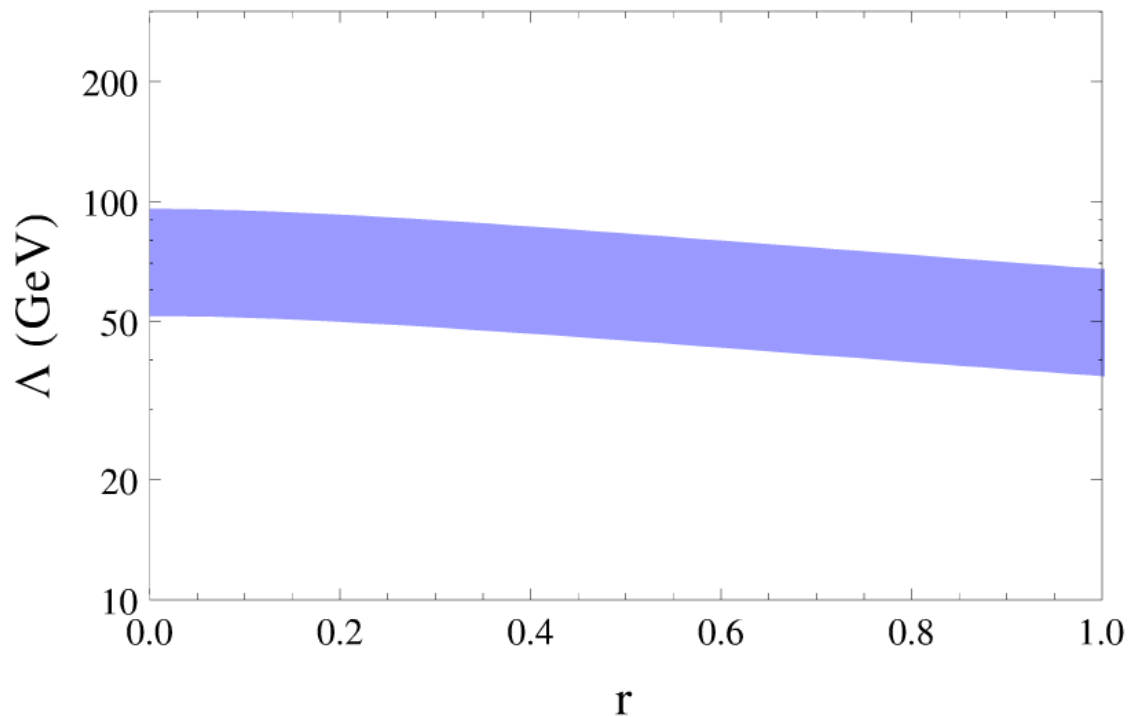


Direct detection



Vector

$$\hat{a}_\mu^{F,S}$$



$$\Lambda = \frac{m_{\text{Med}}}{\lambda}$$

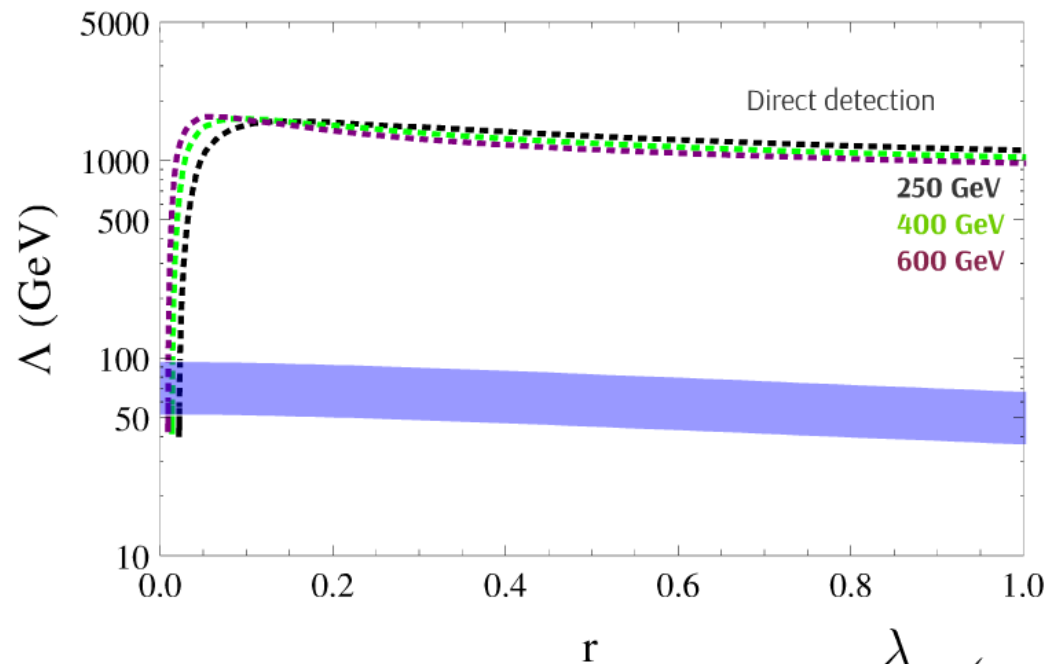
$$r \equiv m_{\text{DM}}/m_{\text{Med}}$$

Real

$$\mathcal{L} = \frac{\lambda}{2} \bar{\mu} (1 - \gamma^5) F \chi + \text{h.c.}$$

Complex

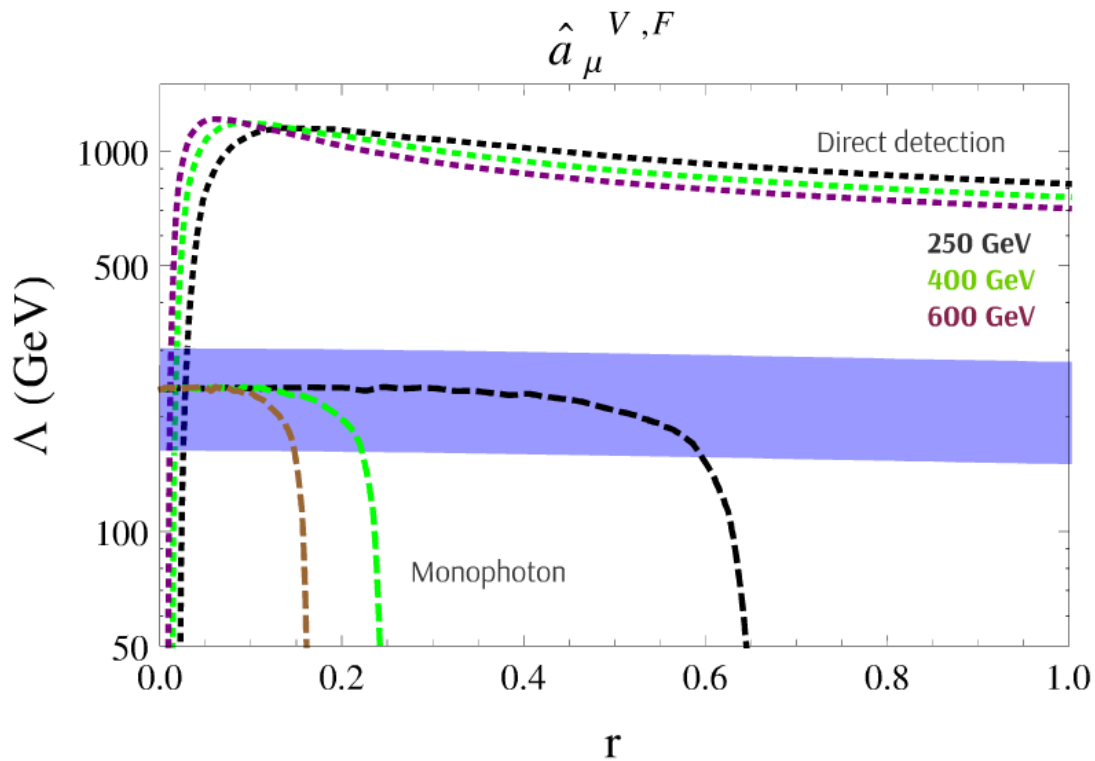
$$\hat{a}_\mu^{F,S}$$



$$\Lambda = \frac{m_{\text{Med}}}{\lambda}$$

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$$\Lambda = \frac{m_{\text{Med}}}{\lambda}$$

$$r \equiv m_{\text{DM}}/m_{\text{Med}}$$

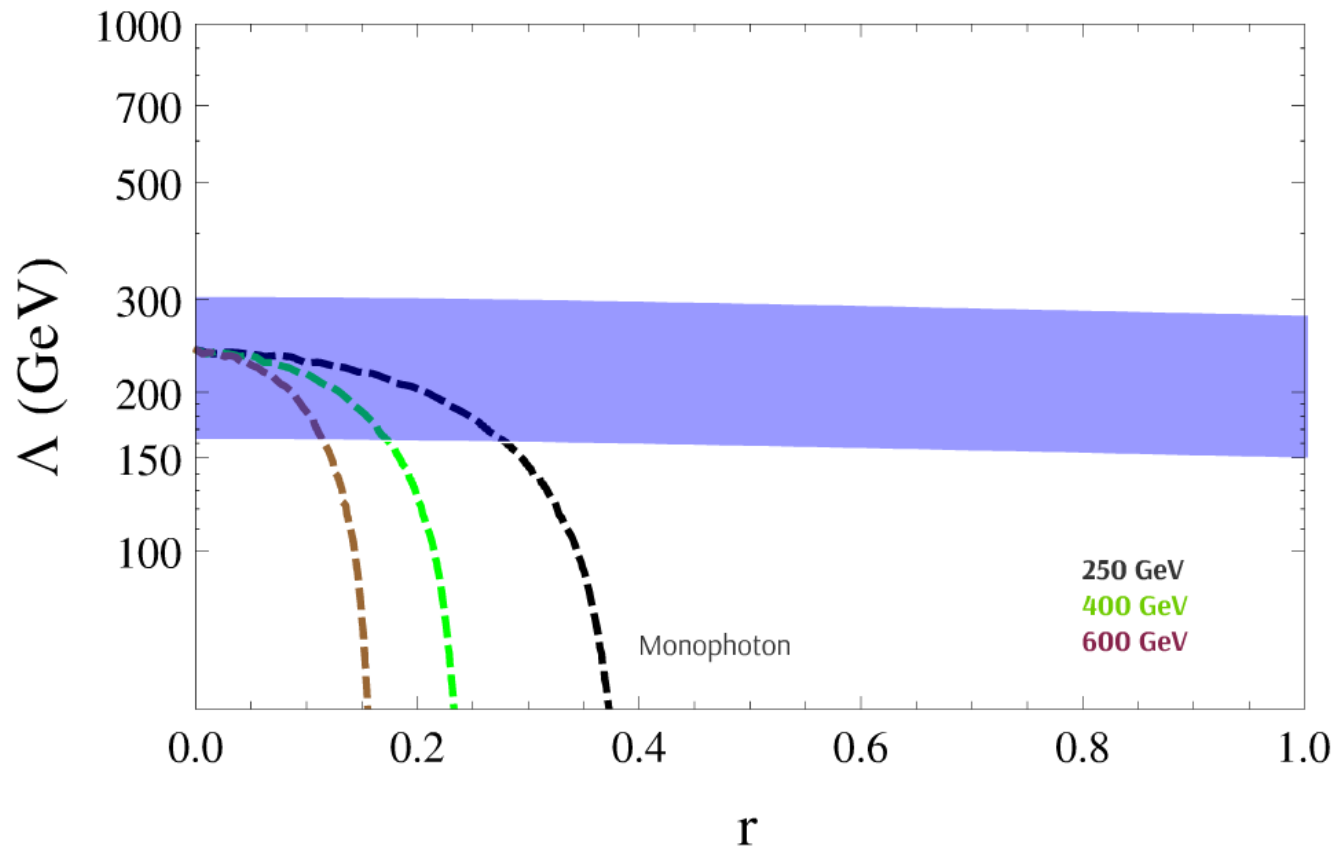
$$\mathcal{L} = \frac{\lambda}{2} \bar{\mu} (1 - \gamma^5) \chi \phi_l + \text{h.c.}$$

Dirac

Majorana

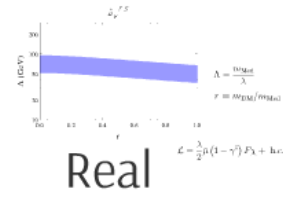
$$\mathcal{L} = \frac{\lambda}{2} \bar{\mu} (1 - \gamma^5) \chi \phi_l + \text{h.c.}$$

$$\hat{a}_\mu^{V,F}$$

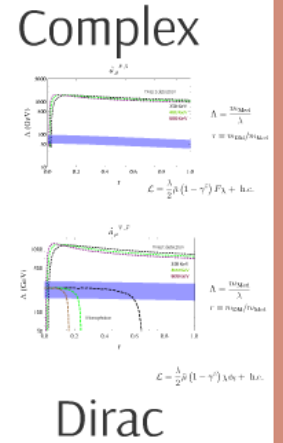


$$\Lambda = \frac{m_{\text{Med}}}{\lambda}$$
$$r \equiv m_{\text{DM}}/m_{\text{Med}}$$

Scalar

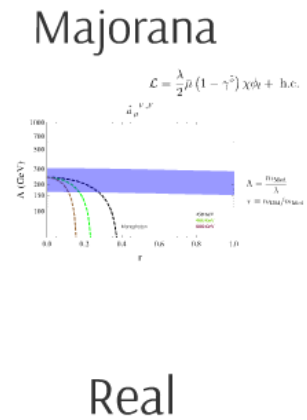


Fermion



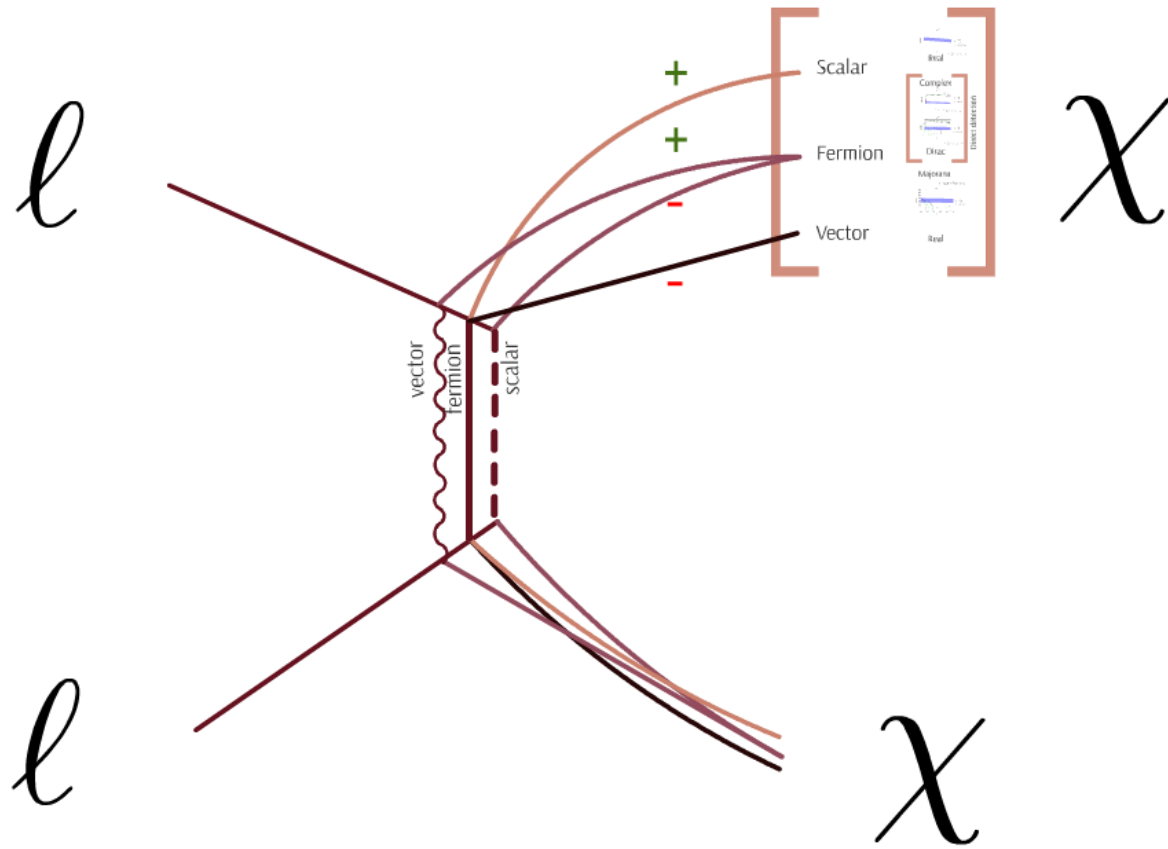
Direct detection

Vector



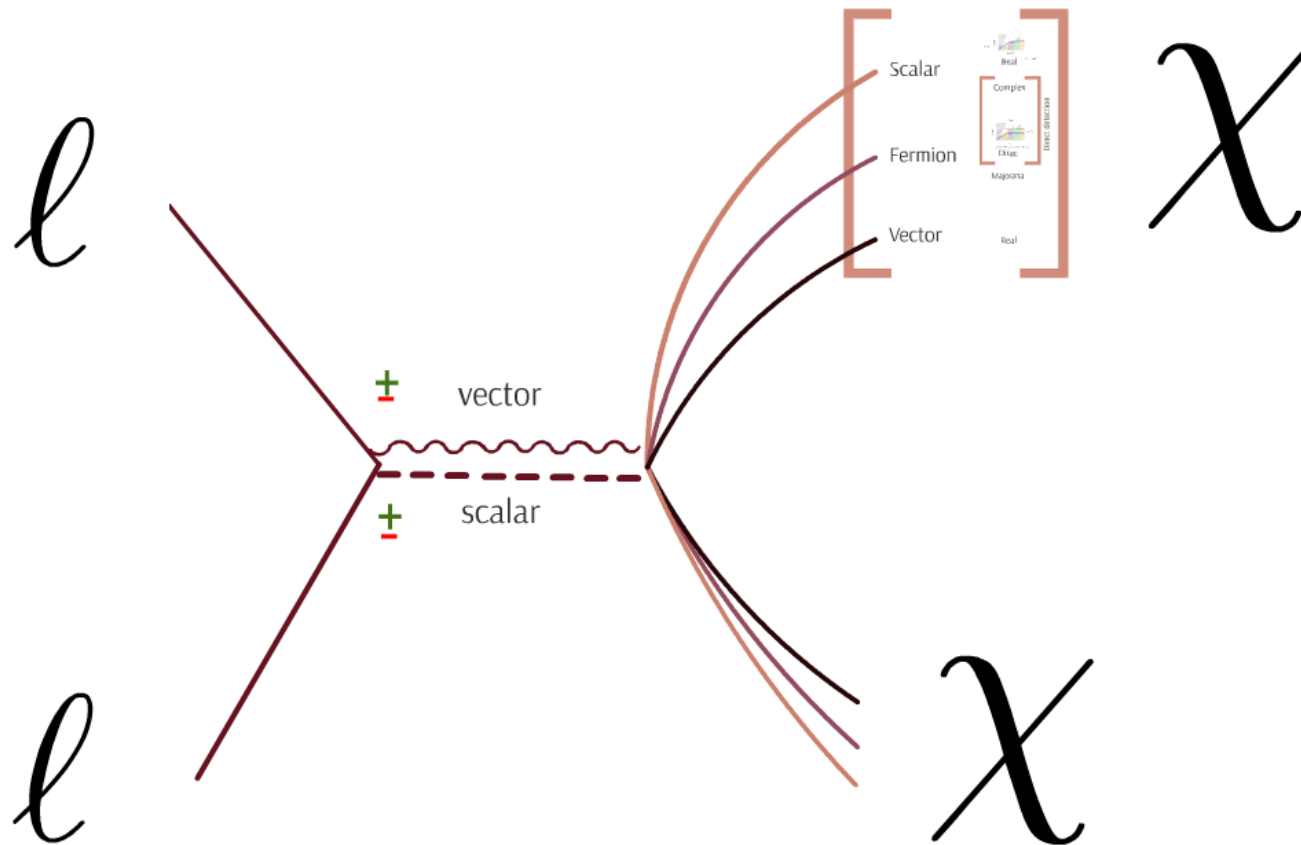
Simplified Models

Charged Mediator



Simplified Models

Neutral Mediator



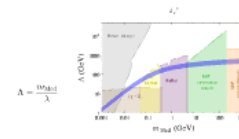


vector



scalar

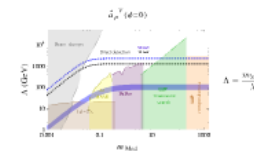
Scalar



Real

$$\mathcal{L}_S = \lambda \bar{\psi} \psi S$$

Complex



Dirac

$$C = \frac{1}{2} \bar{\psi} \psi + i \gamma_5 \bar{\psi} \psi \quad (C_1 + \frac{1}{2} \gamma_5) \psi \quad (i \gamma_5 C_1 + \frac{1}{2} \gamma_5) \psi$$

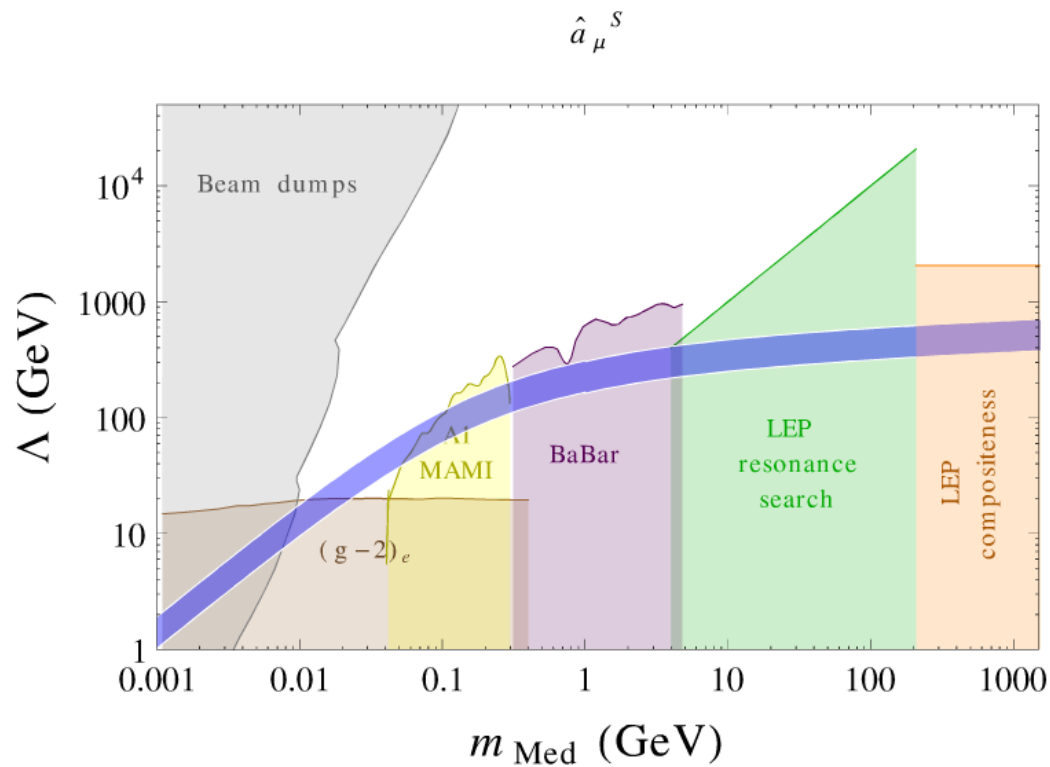
Direct detection

Majorana

Vector

Real

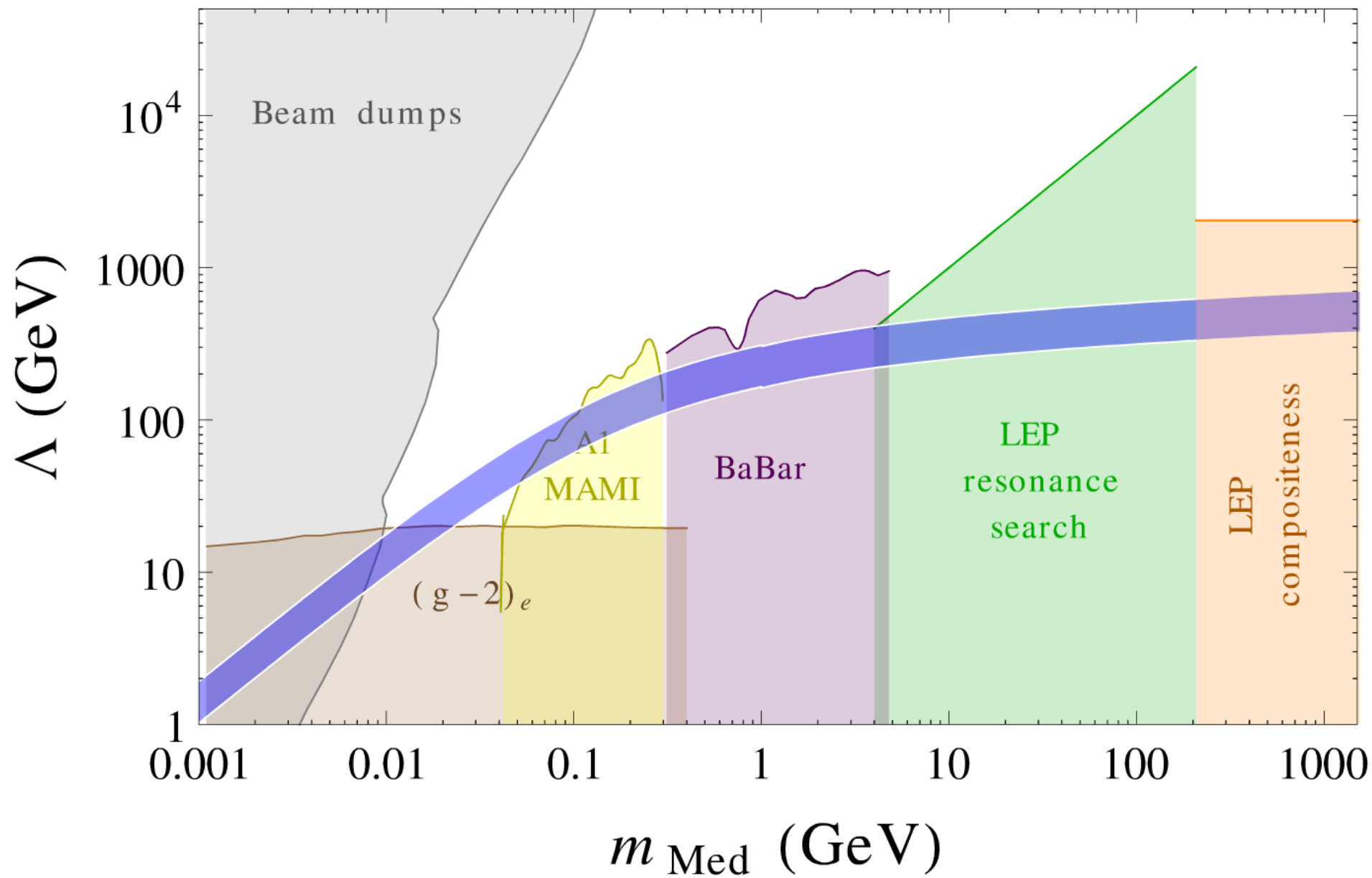
$$\Lambda = \frac{m_{\text{Med}}}{\lambda}$$



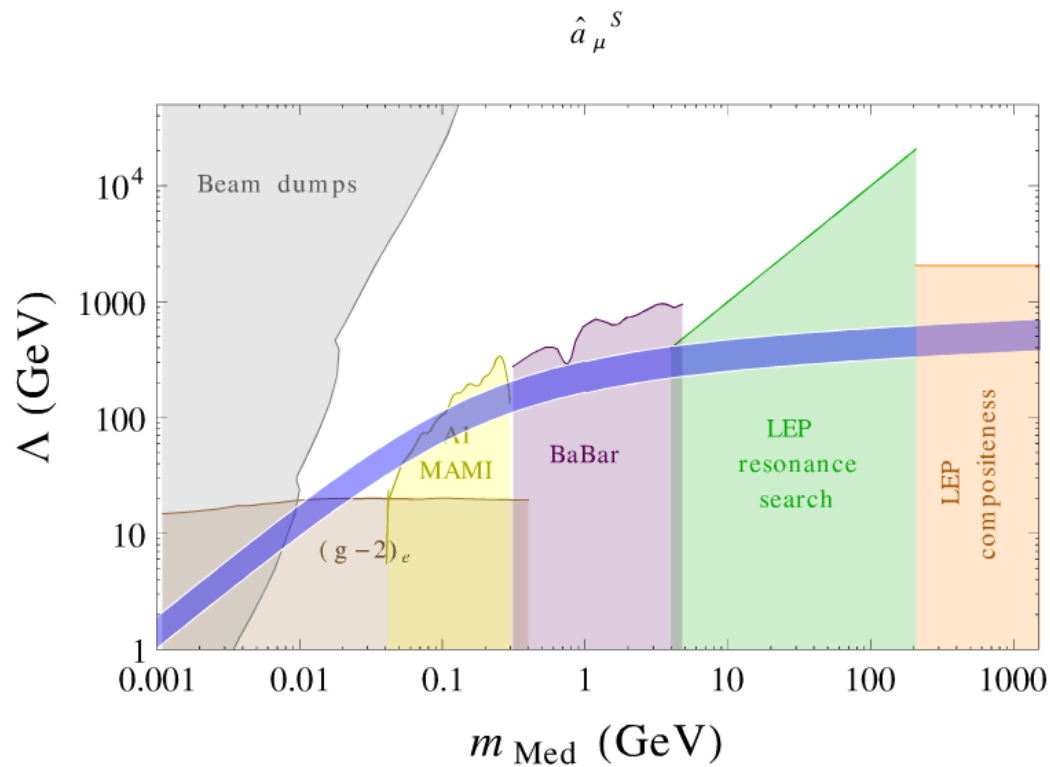
Real

$$\mathcal{L}_S = \lambda \bar{\mu} \mu S$$

$$\hat{a}_\mu^S$$



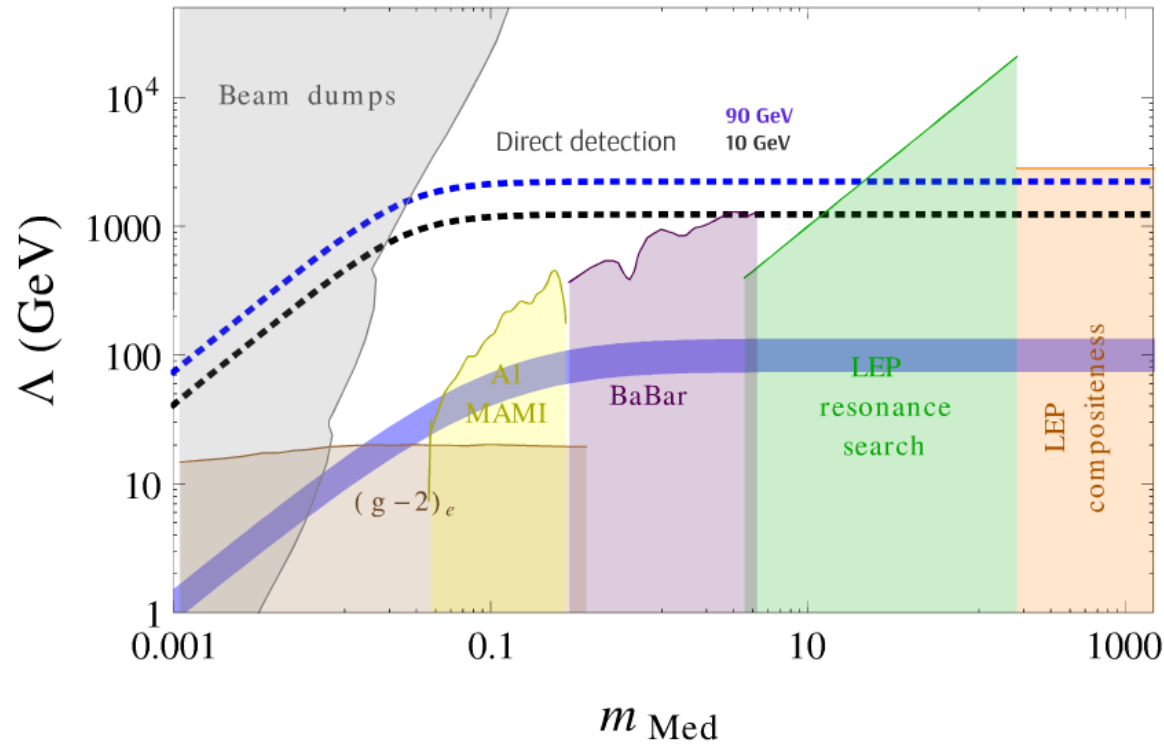
$$\Lambda = \frac{m_{\text{Med}}}{\lambda}$$



Real

$$\mathcal{L}_S = \lambda \bar{\mu} \mu S$$

$$\hat{a}_\mu^V (\phi=0)$$

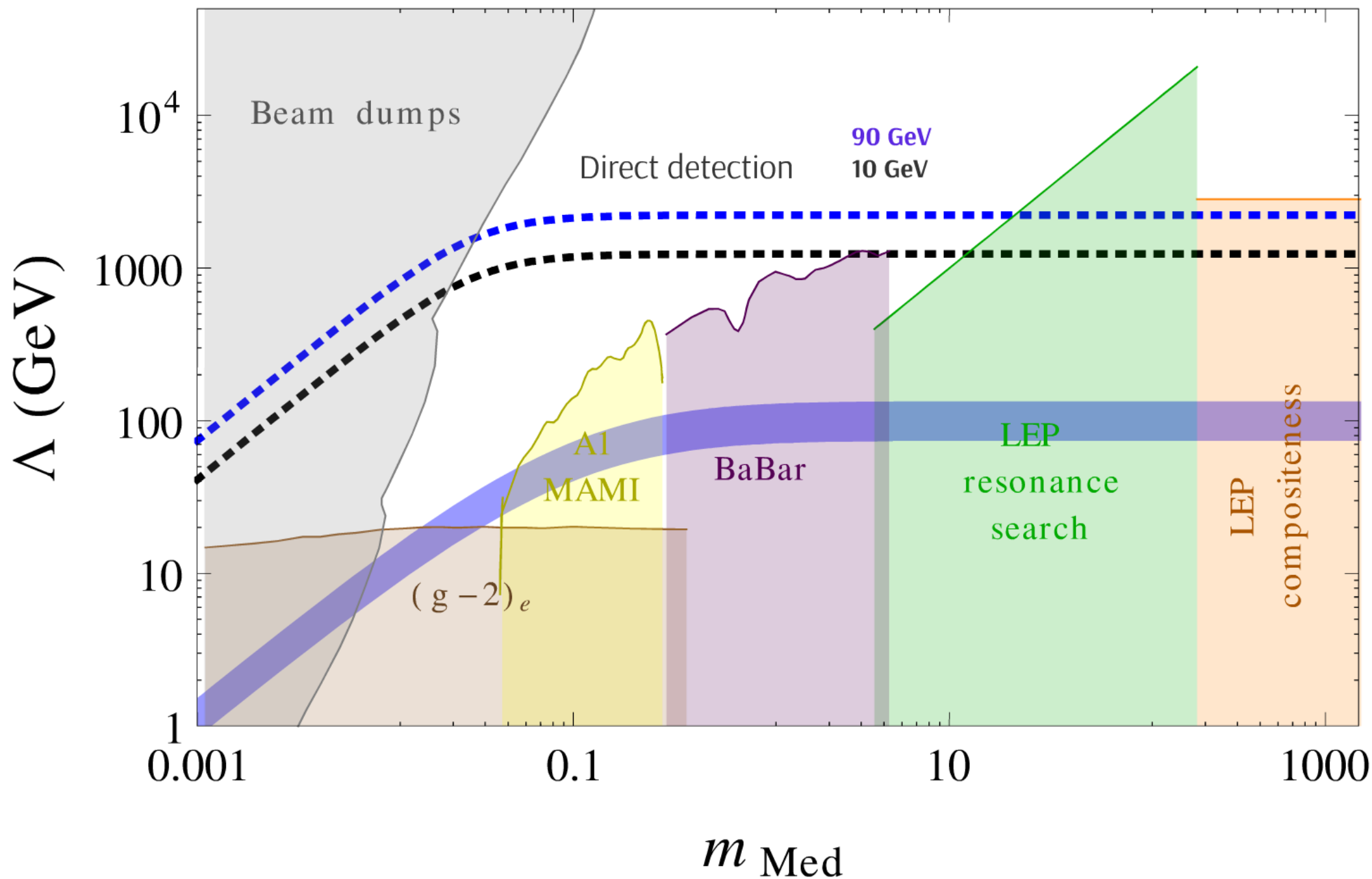


$$\Lambda = \frac{m_{\text{Med}}}{\lambda}$$

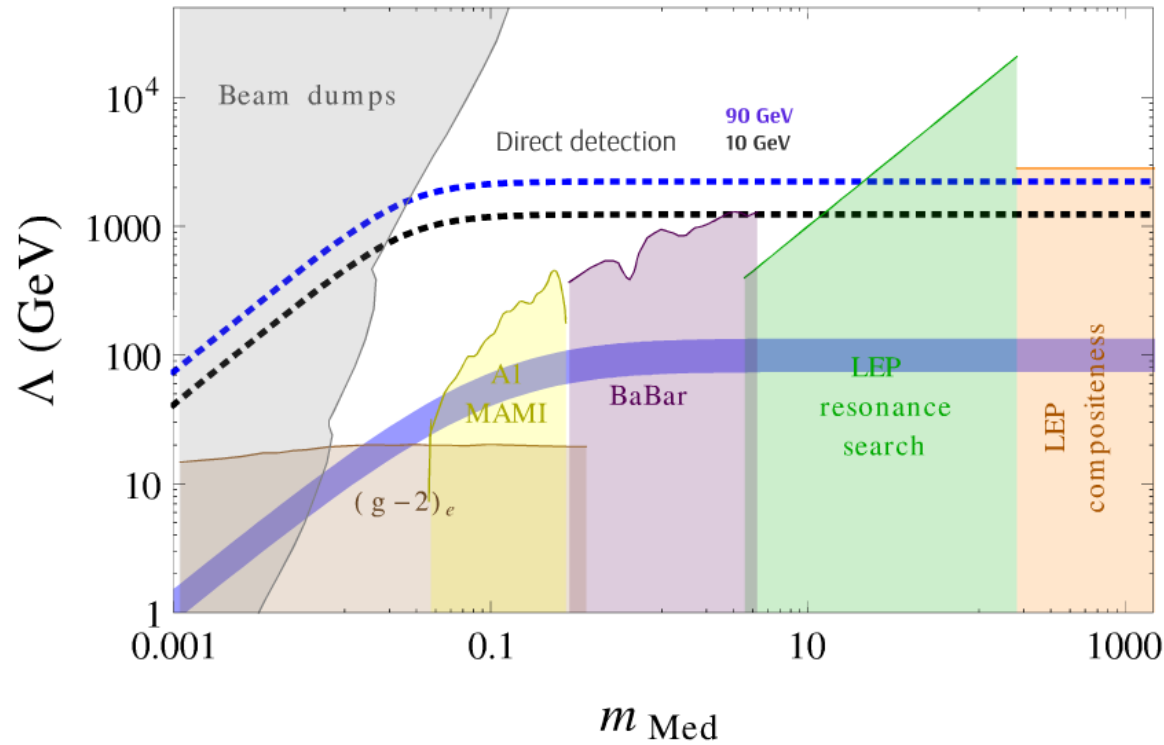
$$\mathcal{L} = \frac{\lambda_\chi}{2} \bar{\chi} \gamma^\nu (\cos \theta + \sin \theta \gamma^5) \chi V_\nu + \frac{\lambda}{2} \bar{\mu} \gamma^\nu (\cos \phi + \sin \phi \gamma^5) \mu V_\nu$$

Dirac

$$\hat{a}_\mu^V (\phi=0)$$



$$\hat{a}_\mu^V (\phi=0)$$

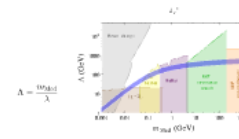


$$\Lambda = \frac{m_{\text{Med}}}{\lambda}$$

$$\mathcal{L} = \frac{\lambda_\chi}{2} \bar{\chi} \gamma^\nu (\cos \theta + \sin \theta \gamma^5) \chi V_\nu + \frac{\lambda}{2} \bar{\mu} \gamma^\nu (\cos \phi + \sin \phi \gamma^5) \mu V_\nu$$

Dirac

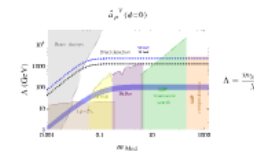
Scalar



Real

$$\mathcal{L}_S = \lambda \bar{\psi} \psi S$$

Complex



Dirac

$$C = \frac{1}{2} (\psi^T (\cos \theta + i \sin \theta \gamma_5) \psi) + \frac{1}{2} (\bar{\psi} (\cos \theta + i \sin \theta \gamma_5) \psi)$$

Direct detection

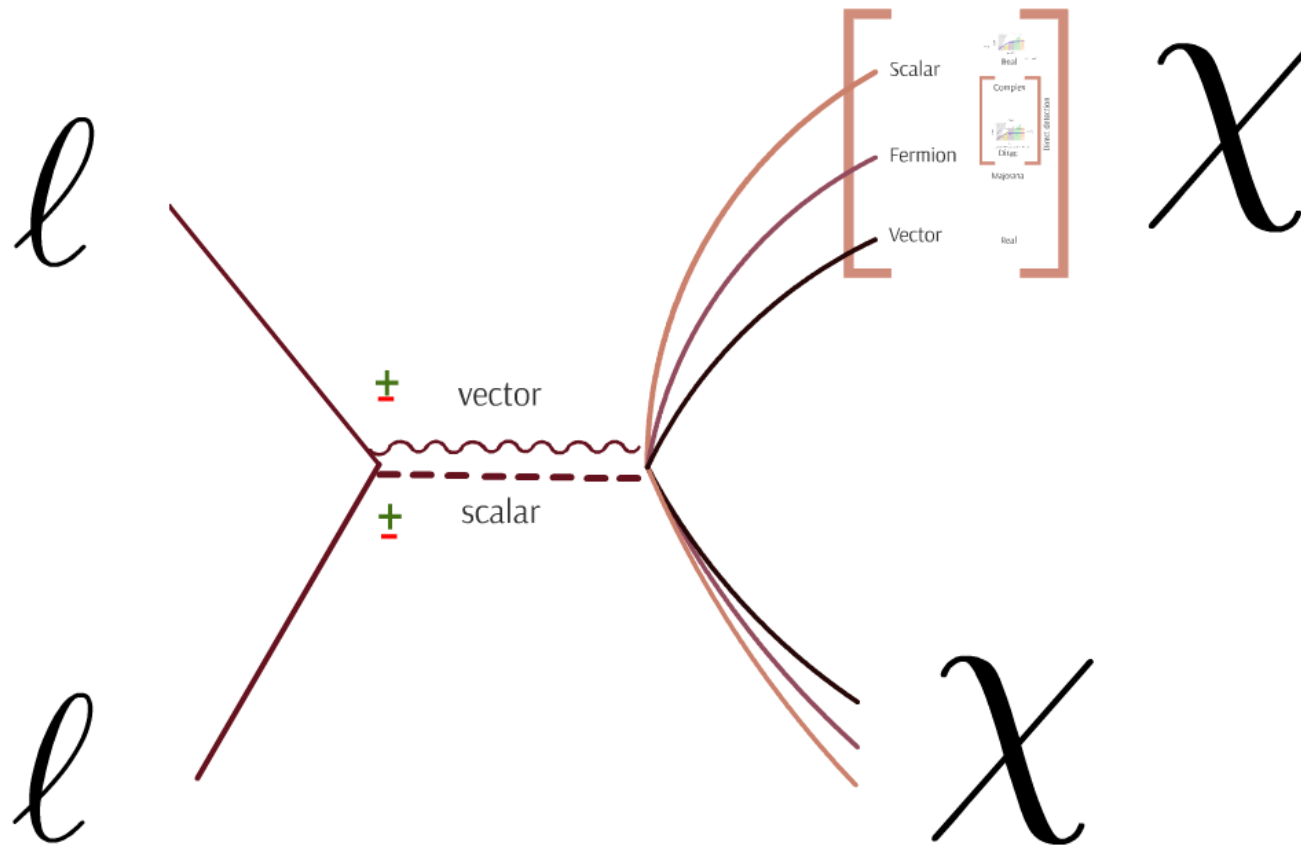
Majorana

Vector

Real

Simplified Models

Neutral Mediator



Conclusions

Leptophilic dark matter has novel phenomenology

Can potentially explain the anomalous magnetic moment of the muon

For chiral couplings, new physics scale is $O(100)$ GeV

Under tension from direct detection and collider constraints

