

Summary of C, P, S, T:

	$\bar{\psi}\psi$	$i\bar{\psi}\gamma^5\psi$	$\bar{\psi}\gamma^\mu\psi$	$\bar{\psi}\gamma^\mu\gamma^5\psi$	$\bar{\psi}\gamma^{\mu\nu}\psi$	∂_μ
P	+1	-1	$(-1)^\mu$	$-(-1)^\mu$	$(-1)^\mu(-1)^\nu$	$(-1)^\mu$
T	+1	-1	$(-1)^\mu$	$(-1)^\mu$	$-(-1)^\mu(-1)^\nu$	$-(-1)^\mu$
C	+1	+1	-1	+1	-1	+1
CPT	+1	+1	-1	-1	+1	-1

$$(-1)^\mu = 1 \text{ for } \mu = 0$$

$$(-1)^\mu = -1 \text{ for } \mu = 1, 2, 3$$

$$\bar{\psi}\gamma^{\mu\nu}\psi + \bar{\psi}\gamma_\mu\psi + \bar{\psi}\gamma_\nu\psi$$

$$P=1, T=-1, C=-1$$

$$CPT=1$$

Dirac Lagrangian:

$$\mathcal{L} = \bar{\psi} (\gamma^\mu \partial_\mu - m) \psi$$

invariant under C, P, S, T separately

$$\text{1st term} \sim \bar{\psi}\gamma^\mu\psi \cdot i\partial_\mu$$

$$P: (-1)^\mu (-1)^\mu = 1$$

$$T: (-1)^\mu [-(-1)^\mu] (1) = 1$$

$$C: -1 \rightarrow +1 \text{ (since } \partial_\mu^T = -\partial_\mu)$$

$$CPT = +1$$

$$\text{Recall: } C \bar{\psi} \psi C = \psi^T \bar{\psi}^T = (\bar{\psi} \psi)^T = \bar{\psi} \psi$$

$$C \bar{\psi} C \partial_\mu C \psi C = \psi^T \partial_\mu \bar{\psi}^T = -(\bar{\psi} \partial_\mu \psi)^T = -\bar{\psi} \partial_\mu \psi$$

$$\partial_\mu^T = -\partial_\mu \quad \quad \quad -\partial_\mu^T$$

$$\langle \alpha | \circ \beta \rangle = \langle \circ^\dagger \alpha | \beta \rangle \quad \circ^\dagger = (\circ^T)^* =$$

$$\partial_\mu^+ = -\partial_\mu \quad \quad \quad \partial_\mu^* = \partial_\mu \quad \quad \quad = (\circ^*)^T$$

$$\int \alpha^* \partial_\mu \beta = - \int \beta \partial_\mu \alpha^* = \quad (i \partial_\mu)^+ = i \partial_\mu$$

Can add \mathcal{L} violating P or C or T

e.g. P S C $\mathcal{L} = \bar{\psi} \not{\partial} \psi + \bar{\psi} \not{\partial}_\mu \not{\partial}^\mu \psi$

$$P = -1, C = -1, T = +1 \quad \text{CPT} = +1$$

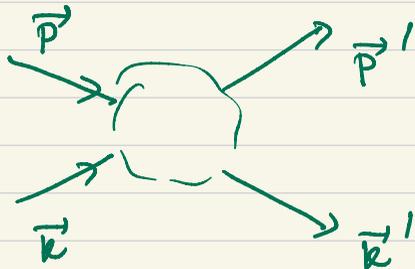
T: $\bar{\psi} \not{\partial}^{\mu\nu} \psi + \bar{\psi} \not{\partial}_\mu \psi + \bar{\psi} \not{\partial}_\nu \psi$

Thm: \nexists Lorentz-invariant QFT with $H^\dagger = H$ that violates CPT

Interacting field theories

$k \in$ theory:
$$\phi(x) = \int \frac{1}{\sqrt{2\epsilon_{\vec{p}}}} \left(a_{\vec{p}} e^{-i\vec{p}\cdot x} + a_{\vec{p}}^{\dagger} e^{i\vec{p}\cdot x} \right)$$

2-body int:



$$a_{\vec{p}'}^{\dagger} a_{\vec{k}'}^{\dagger} a_{\vec{p}} a_{\vec{k}} \leftarrow \phi^4$$

similarly:

3-body

$$\phi^6$$

4-body

$$\phi^8$$

⋮

n -body

$$\phi^{2n}$$

$$\phi^2(x) \phi^2(y)$$

$\phi^2(x) \phi^2(y)$ - non local

Dimensional analysis (~~poor man's RG~~)

$$S = \int dt \underbrace{d^3x}_{M^{-4}} \mathcal{L} \quad \Rightarrow \quad \mathcal{L} \sim M^4$$

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$$\hbar = c = 1 \quad L \equiv T \equiv M^{-1} \quad m \equiv E \equiv \frac{1}{T} \equiv \frac{1}{L}$$

$$m \equiv E \equiv M$$

$$\text{KE: } \mathcal{L}_0 = \frac{(\partial_\mu \phi)^2}{2} - \frac{m^2 \phi^2}{2}$$

$$\text{Dirac: } \mathcal{L}_0 = \bar{\psi} (i \partial^\mu \partial_\mu - m) \psi$$

$$KR: \mathcal{L}_0 = \frac{(\partial_\mu \phi)^2}{2} - \frac{m^2 \phi^2}{2}$$

$$\frac{\partial \phi}{\partial x^\mu} \equiv \frac{\phi}{L}$$

$$Dirac: \mathcal{L}_0 = \bar{\psi} (i \gamma^\mu \partial_\mu - m) \psi$$

$$\frac{\phi^2}{L^2} - m^2 \phi^2 \sim M^4 \Rightarrow \phi \sim M$$

ϕ dimension 1

$$\frac{\psi^2}{L} - m \psi^2 \sim M^4 \Rightarrow \psi^2 \sim M^3, \psi \sim M^{3/2}$$

ψ - dim 3/2

What int can we have?

scalar theory: $\alpha \phi^4$

Axiom: Theory must be renormalizable

Consider e.g. $\langle \phi(x) \phi(y) \rangle = \int \frac{e^{-i p \cdot (x-y)}}{2 E_{\vec{p}}}$

Let $x=y$



$$\langle \phi^2(x) \rangle \propto \int_0^{\infty} \frac{p^2 dp}{\sqrt{p^2 + m^2}} \rightarrow \int_0^{\Lambda} \frac{p^2 dp}{p} = \frac{\Lambda^2}{2}$$

Need to regularize: cutoff regularization $\infty \rightarrow \Lambda$

In the reg, $\Lambda \rightarrow \infty$. Physical quantities - should remain finite, i.e., Λ -independent

$$\alpha \phi^u \sim M^4 \Rightarrow \alpha \sim M^{4-u}$$

$$\phi \sim M$$

$$\text{scat. ampl.} \sim \alpha \Lambda^k = M^{4-u} \Lambda^{u-4} = \left(\frac{M}{\Lambda} \right)^{4-u}$$

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$$u \leq 4 \Rightarrow \mathcal{L}_{int} \sim \mu \phi^3, \lambda \phi^4$$

Must have coupling $\sim M^k$ with $k \geq 0$

$$4-u \geq 0$$

Dirac field ψ :

1) ~~self int~~ $\psi^3 \sim M^{3/2}$ $4 - 4 = 4 - \frac{3}{2} < 0$

Only possible int: $g \bar{\psi} \psi \phi$ ϕ^3 ϕ^4

$$\frac{3}{2} + \frac{3}{2} + 1 = 4 \quad \text{OK}$$

Yokawa theory

$$\mathcal{L} = \mathcal{L}_{\text{Dirac}} + \mathcal{L}_{\text{KG}} - g \bar{\psi} \psi \phi$$

Scalar $\phi \in \mathbb{D}$

$$\mathcal{L} = |\partial_\mu \phi|^2 - m^2 |\phi|^2$$

$D_\mu \equiv \partial_\mu + i e A_\mu(x)$ gauge covariant derivative

$$e^2 |\phi|^2 A^2 \quad e A^\mu \phi \partial_\mu \phi \quad e \bar{\psi} \gamma^\mu \psi A_\mu$$

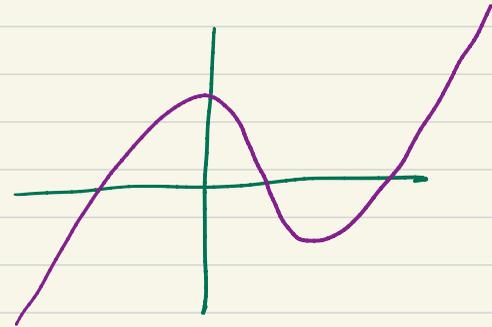
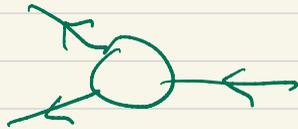
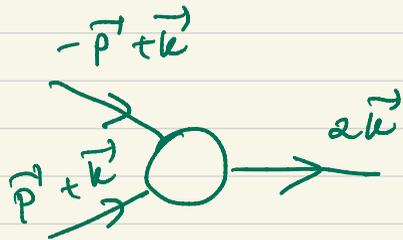
Back to $k f$:

$$1) \mu \phi^3$$

$$\tilde{\mathcal{L}} = \frac{\pi^2}{2} + \frac{u^2 \phi^2}{2} + \mu \phi^3$$

h.o.

$$\frac{p^2}{2} + \frac{\omega^2 x^2}{2} + \mu x^3$$



$$E_{\text{gr. st}} \rightarrow -\infty$$

unbounded
spectrum

$$\mathcal{V}_e = \frac{\pi^2}{2} + \frac{\mu^2 \phi^2}{2} + \frac{\lambda}{4!} \phi^4$$

$$\mathcal{L} = \frac{(\partial_\mu \phi)^2}{2} - \frac{\mu^2 \phi^2}{2} - \frac{\lambda}{4!} \phi^4$$

phi-fourth theory.