

1 Physics 613: Problem Set 2 (due Friday March 20)

1.1 Spin and the Dirac Equation

1. Verify that $[L_i, P_j] = i\epsilon_{ijk}P_k$ where $\mathbf{L} = \mathbf{r} \times \mathbf{p}$ is the angular momentum operator.
2. Verify that $[L_z, H_{Dirac}] = i(\alpha_x P_y - \alpha_y P_x)$. What are $[L_x, H_{Dirac}]$ and $[L_y, H_{Dirac}]$?
3. Verify that with $\mathbf{S} = \frac{1}{2} \begin{pmatrix} \boldsymbol{\sigma} & 0 \\ 0 & \boldsymbol{\sigma} \end{pmatrix} \equiv \frac{1}{2}\boldsymbol{\Sigma}$, the total angular momentum $\mathbf{J} = \mathbf{L} + \mathbf{S}$ commutes with H_{Dirac} .

1.2 Solutions to the Dirac Equation

In class we introduced the solutions to the Dirac equation $\psi(x) = u_s(k)e^{-ikx}$ and $\psi(x) = v_s(k)e^{ikx}$.

1. Verify by plugging into the Dirac equation that u_s and v_s satisfy

$$(\not{k} - m)u_s(k) = 0 \quad (1)$$

and

$$(\not{k} + m)v_s(k) = 0 \quad (2)$$

2. By considering the eigenvalues of $\not{k} - m$ and $\not{k} + m$, prove that there are exactly two independent u_s and two independent v_s solutions for every k (so $s = 1, 2$).
3. Show that the *helicity operator* $h = \frac{\mathbf{P} \cdot \boldsymbol{\Sigma}}{|\mathbf{P}|}$ commutes with the Dirac Hamiltonian $H_{Dirac} = \boldsymbol{\alpha} \cdot \mathbf{P} + \beta m$ and find the eigenvalues of h .
4. For momentum in the z direction (i.e. $\mathbf{k} = (0, 0, k)$), find explicitly the solutions u_s and v_s classified by eigenvalues of the helicity operator.

1.3 Dirac Hamiltonian and charge

1. Substitute the mode expansion

$$\Psi(x) = \int \frac{d^3\mathbf{k}}{(2\pi)^3 2E_{\mathbf{k}}} \sum_s (b_s(\mathbf{k})u_s(\mathbf{k})e^{-ikx} + d_s^\dagger(\mathbf{k})v_s(\mathbf{k})e^{ikx}) \quad (3)$$

into the Dirac Hamiltonian

$$H = \int d^3x (-i\bar{\Psi}\gamma^i\partial_i\Psi(x) + m\bar{\Psi}\Psi(x)) \quad (4)$$

and, using the canonical commutation relations for b_s and d_s , derive

$$H = \sum_s \int \frac{d^3\mathbf{k}}{(2\pi)^3 2E_{\mathbf{k}}} E_{\mathbf{k}} (b_s^\dagger(\mathbf{k})b_s(\mathbf{k}) + d_s^\dagger(\mathbf{k})d_s(\mathbf{k})) \quad (5)$$

2. Do the same for the charge operator

$$Q = \int d^3x \Psi^\dagger \Psi \quad (6)$$

and derive

$$Q = \sum_s \int \frac{d^3\mathbf{k}}{(2\pi)^3 2E_{\mathbf{k}}} (b_s^\dagger(\mathbf{k})b_s(\mathbf{k}) - d_s^\dagger(\mathbf{k})d_s(\mathbf{k})) \quad (7)$$

1.4 Dirac Matrix Identities

Prove the following identities involving Dirac matrices:

1. $\text{Tr}(\gamma^\mu) = 0$
2. $\text{Tr}(\gamma^\mu \gamma^\nu) = 4\eta^{\mu\nu}$
3. $\text{Tr}(\gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma) = 4\eta^{\mu\nu} \eta^{\rho\sigma} - 4\eta^{\mu\rho} \eta^{\nu\sigma} + 4\eta^{\mu\sigma} \eta^{\nu\rho}$
4. $(\gamma^\mu)^\dagger = \gamma^0 \gamma^\mu \gamma^0$
5. $(\bar{f} \gamma^{\mu_1} \dots \gamma^{\mu_n} f')^* = \bar{f}' \gamma^{\mu_n} \dots \gamma^{\mu_1} f$ where f and f' can be any Dirac spinor (i.e. u_s or v_s).

1.5 Rutherford Scattering

In class we used crossing symmetry to transform $e^+e^- \rightarrow \mu^+\mu^-$ into $e^-\mu^- \rightarrow e^-\mu^-$ (in class we called it p instead of μ^- but it doesn't matter); the answer for the squared and summed/averaged matrix element is

$$\frac{1}{4} |\mathcal{M}|^2 = \frac{2e^4}{t^2} (u^2 + s^2 + 4t(m_e^2 + m_\mu^2) - 2(m_e^2 + m_\mu^2)^2) \quad (8)$$

1. Rederive this directly from the t -channel Feynman diagram for $e^-\mu^- \rightarrow e^-\mu^-$ scattering (thereby verifying explicitly the validity of crossing symmetry in this example).

2. Carefully take the $m_\mu \rightarrow \infty$ limit and derive the Mott formula:

$$\left. \frac{d\sigma}{d\Omega} \right|_{m_\mu \rightarrow \infty} = \frac{e^4}{64\pi^2 v^2 p^2 \sin^4 \frac{\theta}{2}} (1 - v^2 \sin^2 \frac{\theta}{2}) \quad (9)$$

where $v = p/E$ and p and E are the 3-momentum and energy of the incoming electron respectively. [Be careful! I think the treatment in Matt Schwartz's book may not be completely correct!]