

2/4 More Frustrated Spin Systems

Spin Models = "Economy" Strongly Correlated Systems

AFM + frustration \rightarrow Degenerate
↓
gd. state
competing
interactions

Large Ground State Degeneracy (extensive)

- spins highly correlated but fluctuating strongly
 - classical and/or quantum fluctuations
- remarkable collective phenomena
- emergent gauge fields
 - fractional particle excitations

2.

Frustration is a short wavelength phenomenon associated with the geometry of the lattice.

Local properties (constraint) (geometry) \Rightarrow Long wavelength behavior

Effect of Temperature on degenerate ground state manifold ?

Order by Disorder

Degenerate ground state is very sensitive to ground-state fluctuations

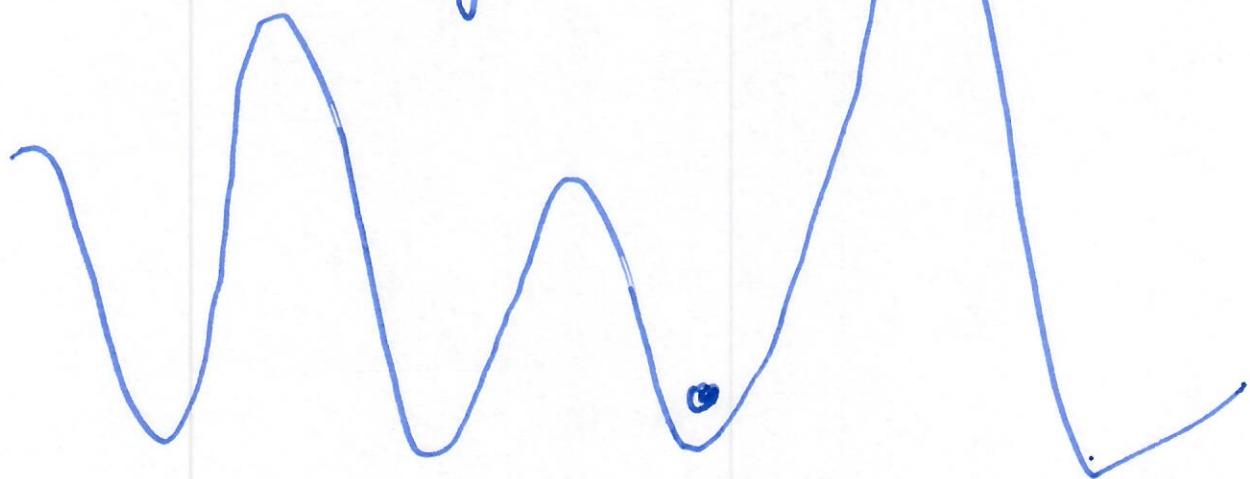
Thermal (and quantum) fluctuations
select an ordered phase from the
ground-state manifold

$$F = E - TS$$



maximize
entropy .

* Another possibility



finite T



spin freezing!
(system gets stuck!)

spin glasses



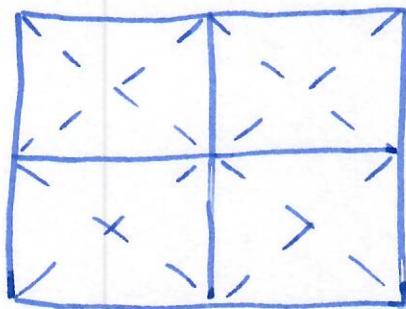
neural networks



machine learning

• A Simple Example of Order by Disorder

$$H = J_1 \sum_{nn} \vec{S}_i \cdot \vec{S}_j + J_2 \sum_{nnn} \vec{S}_i \cdot \vec{S}_j$$



— J_1
--- J_2

Classically

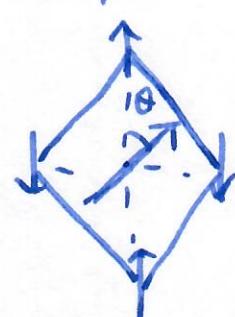
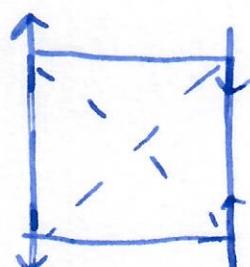
- $J_1 \gg J_2$

$$E = S^2 \{-2J_1 + 2J_2\}$$

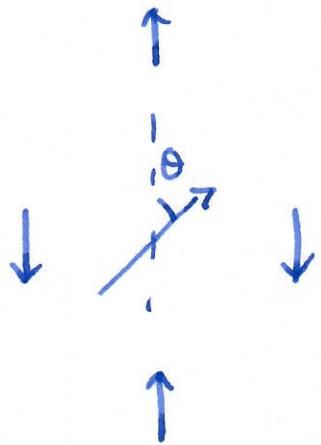
- $J_2 \gg J_1$

$$E = S^2 \left\{ -2J_2 + 2J_1 \cos \theta - 2J_1 \sin \theta \right\}$$

Classically the two sublattices are decoupled.



Spin Wave Theory



$$\propto \left(\frac{1}{\pi} S \right)$$

$$\omega(q, \theta) = (4J_2 S) f(q_x, q_y, \theta)$$

$$\delta F(\theta) = F(\theta) - F$$



$$F(\theta) = T \sum \ln \frac{\omega(q, \theta)}{T}$$

$$\frac{J_2}{J_1} \gg 1$$

Favors $\theta = \pm 1$

$\uparrow \uparrow$
 $\downarrow \downarrow$
 $\uparrow \uparrow$

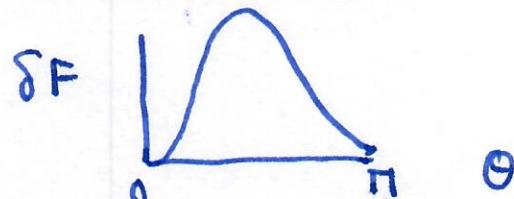
$\uparrow \downarrow$
 $\uparrow \downarrow$
 $\uparrow \downarrow$

FM $\omega \sim q^2$

AFM $\omega \sim q$

$$\boxed{\delta F(\theta) = -E(T) \{ 1 + \alpha \cos^2 \theta \}}$$

2 states selected



False zero mode

associated w/

"accidental soft" mode

Discrete degree of freedom

$$\vec{M}_1, \vec{M}_2 = \sigma = \pm 1.$$

Fluctuation - Selected Order (2d)

$$\xi \sim e^{\frac{J_2 S^2}{T}} \underset{\text{finite}}{\underline{}} \quad \text{at } T \neq 0$$

Despite finite correlation length

(no long-range Heisenberg AFM)



emergent \mathbb{Z}_2 long-range order



emergent \mathbb{Z}_2 phase transition in
a disordered 2D Heisenberg system.

- \mathbb{Z}_2 order parameter

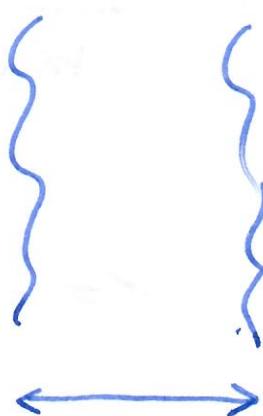
$$\sigma = \frac{(\vec{s}_1 - \vec{s}_3) \cdot (\vec{s}_2 \cdot \vec{s}_4)}{|\vec{s}|^2}$$

order in relative spins
 $(\xi \text{ finite at } T=0)$

Example of composite order

- Phase Diagram?

\mathbb{Z}_2 domains OK if domain walls are defined



$$\xi \sim w_{DW}$$

estimate of stability

W_{DW} ?

$$S \sim \int_{\text{z}} J_2 \sum_i (\partial_{\mu} n_i)^2 - \frac{\gamma}{a^2} \int_{\text{x}} (n_1 \cdot n_2)^2$$

$$\Delta E = - \frac{\chi}{2} \langle B_{\perp}^2 \rangle$$

$$\begin{matrix} 2 \\ \uparrow & \uparrow \\ \sim \frac{1}{J_2} & J_1^2 \end{matrix}$$

$$\frac{E_{DW}}{a} \sim \left\{ \frac{J_2}{w} + \left(\frac{J_1^2}{J_2} \right) \frac{w}{a^2} \right\}$$

$$\frac{dE_{DW}}{dw} = 0 \Rightarrow - \frac{J_2}{w^2} + \frac{J_1^2}{J_2 a^2} = 0$$

⇓

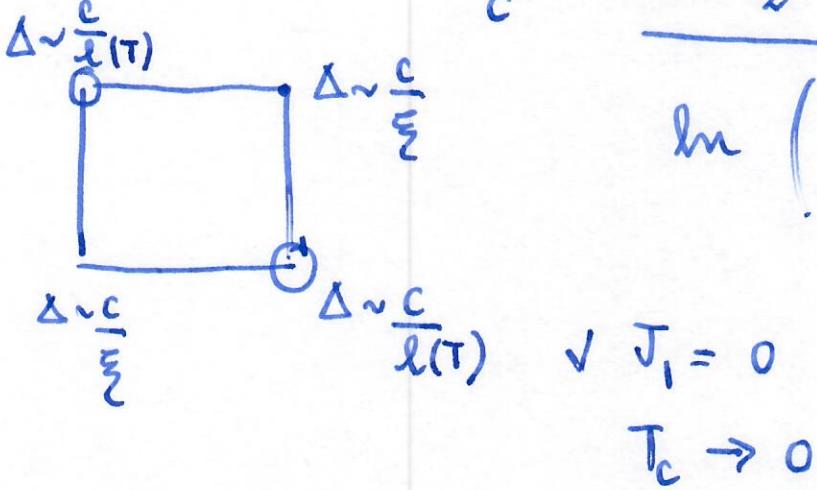
$$w = \left(\frac{J_2}{J_1} \right) a$$

Condition for stability of \mathbb{Z}_2 order

$$e^{2\pi J_2 S^2 / T_c} \sim \frac{J_2}{J_1}$$



$$T_c = \frac{2\pi J_2 S^2}{\ln(J_2/J_1)}$$



- Link to experiment \rightarrow pnictides

Nematic and structural (orthorhombic)

transitions coincide

(and both appear to be driven
by magnetic fluctuations)

How to probe nematic fluctuations?

shear

$$H_{\text{mag-el}} = \lambda \Psi \epsilon_s$$

$$\epsilon_s = \frac{\epsilon_{aa} - \epsilon_{bb}}{\sqrt{2}}$$

$$c_s^{-1} = c_{so}^{-1} + \lambda^2 c_{so}^{-2} \chi_{\text{nem.}}$$

Connection to Superconductivity ? !

Fluctuation-Selected \mathbb{Z}_2 order

Emergent Composite Order

(why it's so hard to identify)
 a classical spin liquid

Order by Disorder

Thermal

$$F = E - TS$$

entropy weight \Rightarrow
 gd. state selection
 from each degenerate
 manifold

Fluctuations \Rightarrow Ordered State

$$S = S_0 + S_D$$

↓

entunious
fluctuations
of each
spin

usually

$$S_0 > S_D$$



difficult to
determine
analytically.

Here above analogy of statistical properties

in disordered systems

and different from regular

Exact enumeration allowed in

recursive structures (e.g. Husimi cactus)



$S_D \sim$ "Delocalization Bandwidth" w

$S_G \sim$ RMS of random part of potential $\sqrt{\delta V_r^2}$
inhibiting diffusion

Here using analogy w/ electronic propagation
in disordered medium

Delocalization criterion

$$w > \sqrt{\delta v_r^2}$$



Breakdown of order by disorder

$$S_D > S_0$$

S_D must scale w/ N

then prefactors crucial

entropy (unlike electronic energy)
no time scale!

Usually

Fluctuations \rightarrow Discrete Order
(frustrated spin system)

short wave-length \rightarrow breaks lattice symmetry

fluctuations in
frustrated system \downarrow Discrete
Long-Range Order