

# 1/28 Introduction to Classical Frustrated Magnetism

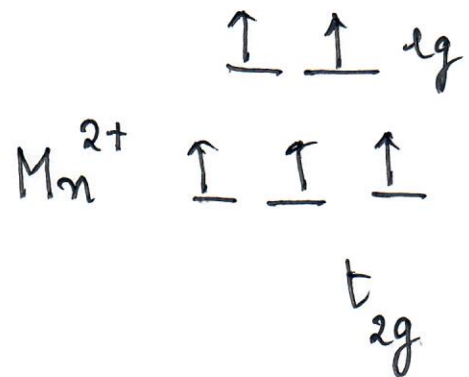
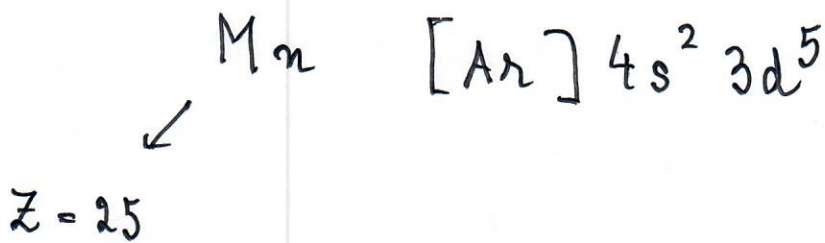
Let's unpack the title

- Magnets

- Focus on systems with well-defined local moments

- Almost isolated ions with partially filled shells.

- Magnetism arises from Hund's rule e.g.



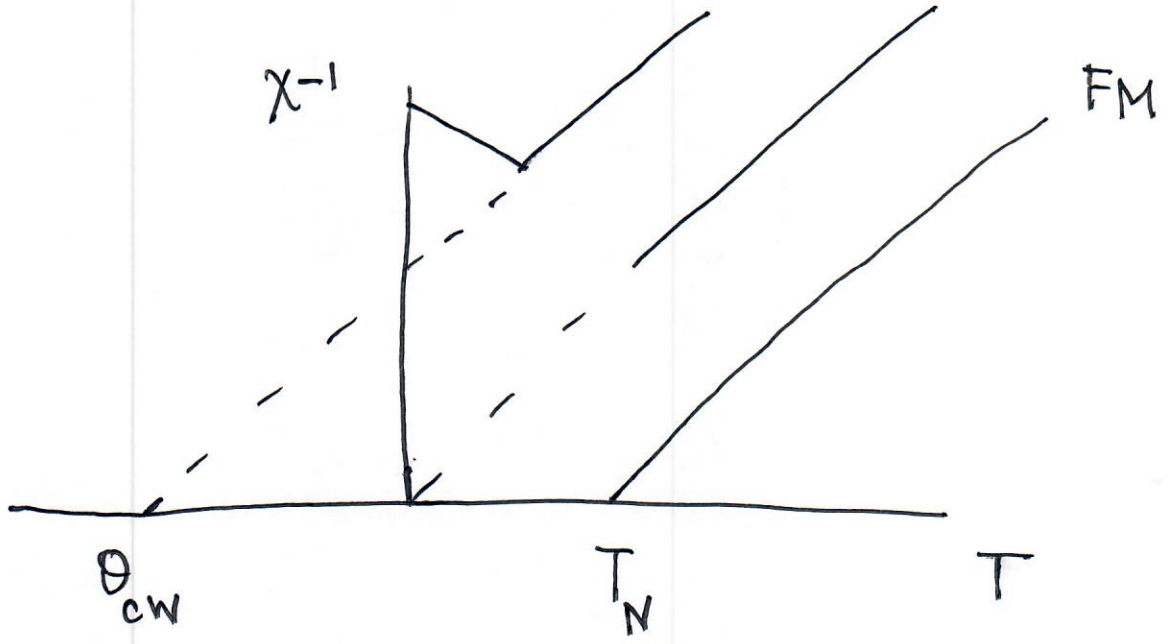
Local atomic physics (e.g. spin orbit, crystal fields) => wide variety of magnetism

# Experimental Identification

- electrically insulating  $\rho \sim e^{E_g/T}$
- Curie susceptibility

$$\chi = \lim_{H \rightarrow 0} \frac{M}{H}$$

Signature of free moments



Ferromagnetism



Let's calculate the net magnetic moment  
of a system of electrons in a solid

3.

Magnetization

$$M = \frac{\vec{\mu}}{V}$$

$$E = -\vec{\mu} \cdot \vec{B}$$

Lorentz force

$$\vec{F} = q (\vec{E} + \vec{v} \times \vec{B})$$

$$\vec{F} \perp \vec{v}$$

No work  
done.

$$E \neq E(B)$$

classically magnetization = 0!

Quantum Mechanics Necessary for  
Magnetism

(Bohr - van Leeuwen Theorem)

Niels Bohr 1911

Hendrika Johanna van Leeuwen 1919

No thermal equilibrium magnetization  
in classical systems.

Basic Idea

$Z$  for  $N$  particles of charge  $q$

$$Z \sim \int \int \int \exp - \beta E (r_i, p_i)$$

$dr_1 \dots dr_N$        $\downarrow$  position  
 $dp_1 \dots dp_N$        $\nearrow$  momenta

$6N$  dimensional phase space

( $3N$  positions,  $3N$  particles)

$$B \neq 0$$

$$p \rightarrow p - qA$$

Limits of integral  $-\infty \rightarrow +\infty$

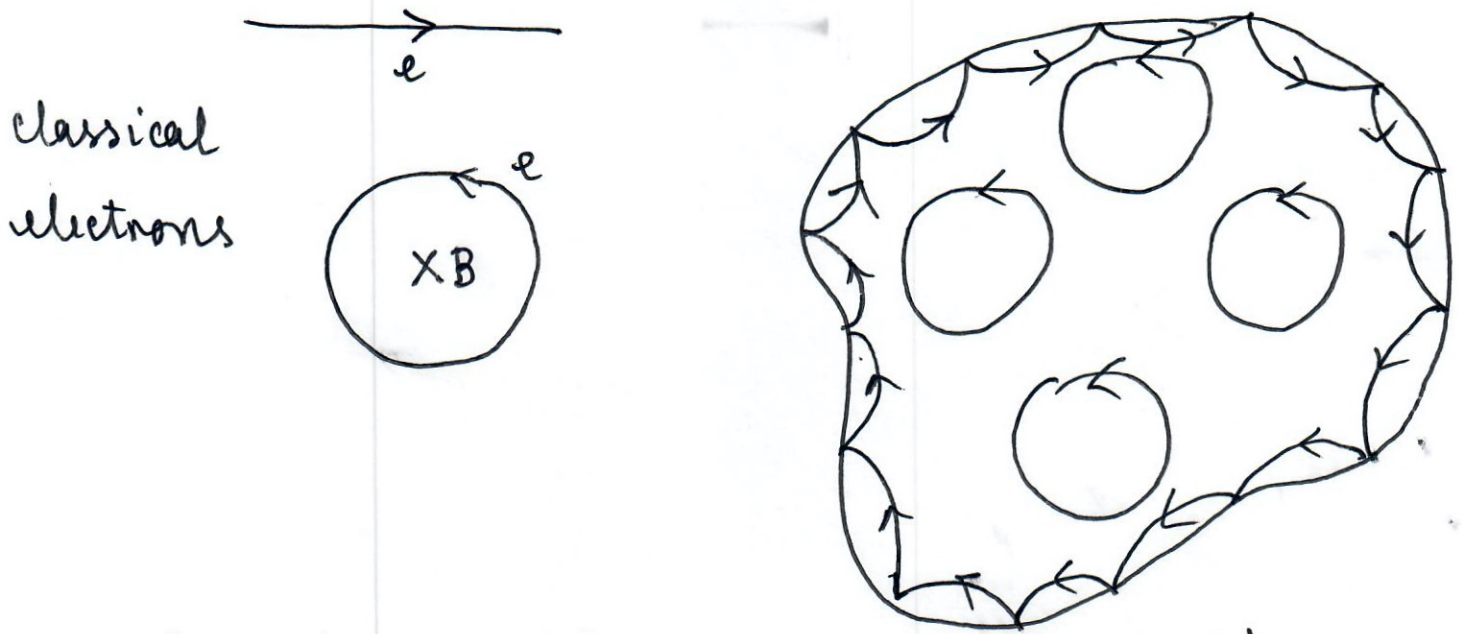
shift absorbed by shifting

origin of momentum integration



$$Z \neq Z(B), \quad F \neq F(B).$$

# Result Intuitively Surprising



Quantum Mechanics Needed!

$$\vec{M} = 0.$$

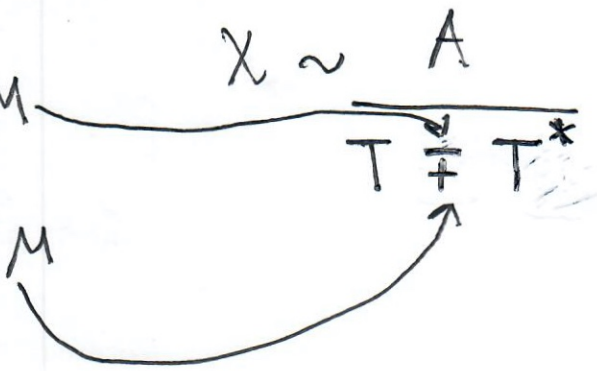
## Interaction between Spins

### Heisenberg Model

$$H = J \sum_{ij} \vec{S}_i \cdot \vec{S}_j$$

$J < 0$  FM

$J > 0$  AFM



# AFM: A Brief History



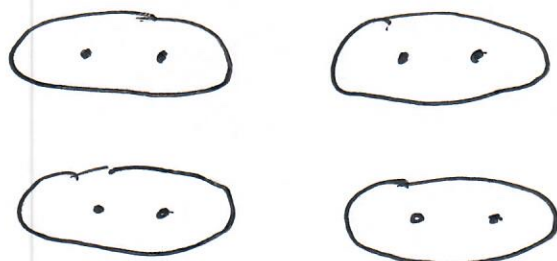
(Néel).

Problem: ↑ ↓ ↑ ↓ not ground state

( $\vec{m}$  is not a conserved order parameter)

Landau

singlet  $\frac{1}{\sqrt{2}} \{ |\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle \}$



Bethe's exact solution of  
power-law (critical)

$S = 1/2$  AFM



philosophically extended to higher  $d$ ?  
(No LRO).

Magnetic Moments  
of Neutrons $\Rightarrow$ Magnetic Moments  
of Electrons.

MnO (NaCl)

X-rays chemical  
unit cell

$$a = 4.43 \text{ \AA}$$

Neutrons 235 K

$$a = 4.43 \text{ \AA}$$

80 K

$$a = 8.85 \text{ \AA}$$

Spins parallel in [111] plane /  
antiparallel in adjacent ones.

$$\langle \vec{S}_i \cdot \vec{S}_j \rangle \propto M_i^2 \cos \vec{Q} \cdot (\vec{R}_i - \vec{R}_j)$$

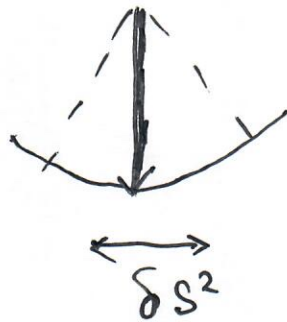
$$\vec{Q} = \left( \frac{\pi}{a}, \frac{\pi}{a}, \frac{\pi}{a} \right)$$

1952 Anderson / Kubo Semiclassical Theory 8.

Analogy with quantum harmonic oscillators

(zero-point energy neglected previously)

AFM = Semiclassical Spins + Zero Point Fluctuations



Fast, slow degrees of freedom

Reduction in sublattice magnetization  $\propto \left(\frac{1}{\sqrt{S}}\right)$

$$\delta M \sim \int d^d q \left\{ n(\omega_q) + \frac{1}{2} \right\} \frac{1}{\omega_q}$$

$T \gg \omega$

$$\delta M \sim \int d^d q \frac{T}{\omega_q^2}$$

$T \ll \omega$

$$\delta M \sim \int \frac{d^d q}{\omega_q}$$



We can get a feeling for this w/ SHO.

### Equipartition Theorem

$$\frac{1}{2} m \omega^2 \langle x^2 \rangle \sim kT \quad \text{classical}$$

$$\langle x^2 \rangle \sim \frac{kT}{\omega^2}$$

$$\frac{1}{2} m \omega^2 \langle x^2 \rangle \sim \hbar \omega$$

$$\langle x^2 \rangle \sim \frac{1}{\omega} \quad \text{quantum}$$

AFM

d-dimensional  
hypercubic  
lattice

$$\omega_q \sim q$$

$T=0$

$d=1$      X

$d \geq 2$      ✓

Neel order OK

$T \neq 0$

$d \geq 3$

Neel order OK

Neutrons - excellent agreement with  
zero point fluctuations in  $\vec{M}$

2D Heisenberg  
model

$T = 0$  Néel  
 $T \neq 0$  Disordered  
(PM)

Aside: 1973 Anderson Distinct absence  
of 2D  $s = 1/2$  AFMs  
low d, low spin frustrated  
systems

could quantum fluctuations  $\Rightarrow$   
destroy LRO?

1987 LaCuO4 2D  $s = 1/2$  AFM

Doped semimetallic  $\rightarrow$  superconducting state

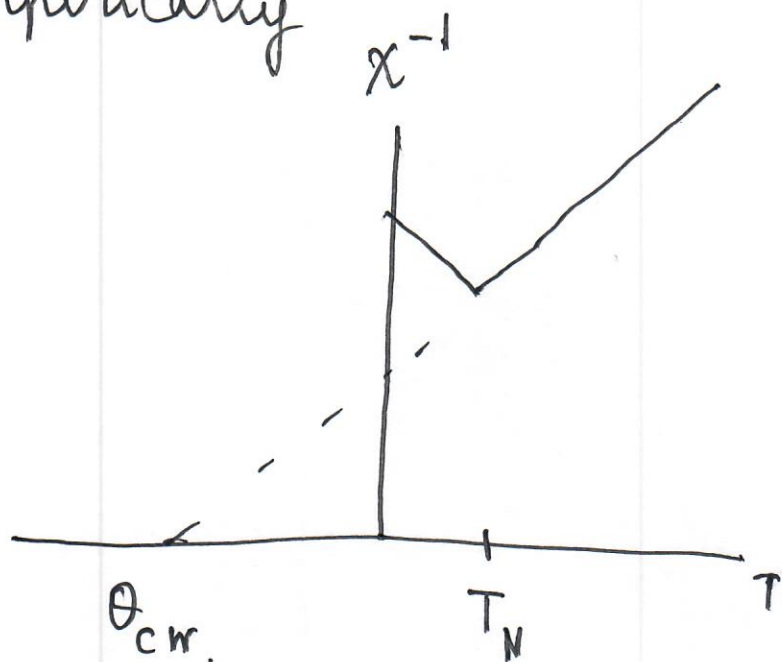
Suggestion  
(Anderson)

Large q. fluctuations  $\Rightarrow$  destroy magnetism

"Spin Liquid" + charge  
new type of Superconductor?

# Frustration

Empirically

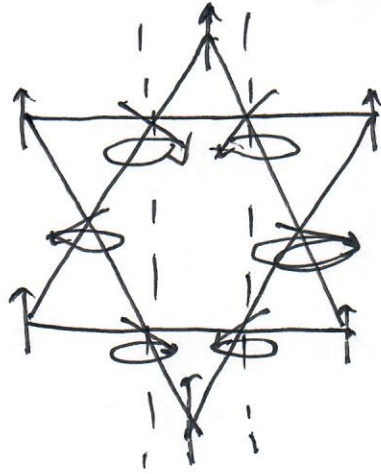


$$f = \frac{|\theta_{cw}|}{T_N} \gg 1.$$

(e.g. 5-10)

Heisenberg systems?

e.g. Kagome



local degeneracy  
(soft mode)

Frustration  $\Rightarrow$  large ground-state  
degeneracy

(why corner-sharing?)

Continuous spins  $\Rightarrow$  simple counting  
argument ("Maxwellian")

$$F = D - k$$

"

"

"

# degrees of  
freedom in  
the ground  
state

total #

DOF

constraints

linear algebra  $\Rightarrow$  system of  $K$  equations  
 for  $D$  variables  
 $\downarrow$   
 solution space  $F = D - K$

Heisenberg Hamiltonian

$$H = J \sum_{\langle ij \rangle} \vec{S}_i \cdot \vec{S}_j \quad J < 0 \quad (\text{AFM})$$

$\downarrow$  cluster of  $q$  spins

$$H = \frac{J}{2} \left( \sum_{i=1}^q \vec{S}_i \right)^2$$

$$= \frac{J}{2} \vec{L}^2$$

$$L = \sum_{i=1}^q \vec{S}_i$$

ground-  
states

$\Leftrightarrow$  states w  
 $\vec{L} = 0$ .

total spin  
of unit,

Heisenberg  
 $n=3$

$$L_1 + L_2 + \dots + L_n = 0$$

independent of  $q$

Spins w/  $n$  components

$\left( \begin{array}{l} n=2 \text{ } x, y \\ n=3 \text{ Heisenberg} \end{array} \right)$

$\Downarrow$

$n$  constraints / unit

$q$  spins/unit

$n$  components

$$F = \frac{\# \text{ degrees of freedom}}{\text{unit share conners}}$$

$$F = (n-1)q - n$$

$\downarrow$

$\#$  global rotations

$$K = n$$

$\downarrow$

gd. state constraint

$$D = F - K = \frac{N}{2} \left[ (n-1)q - 2n \right]$$

$\downarrow$

gs. degeneracy

$$= \frac{N}{2} \left[ n(q-2) - q \right]$$

Heisenberg model  
on  
pyrochlore  
lattice

$$D = \frac{N}{2} [6-4]$$

$\Downarrow$

$$\underline{D = N}$$

Spin Ice

Historically first frustrated system  $\Rightarrow$   
ice.

Entropy measured by Giauque et al.

1956

Anderson

Ice rules work for specific types  
of pyrochlore

systems w/ Ising DOFs.

1997

Holmium Titanate

Ising on pyrochlore.