SSP I 2017
601

Lecture #2

- How to distinguish various phases.

NB. - A concept of an order parameter
(show in class presentation)

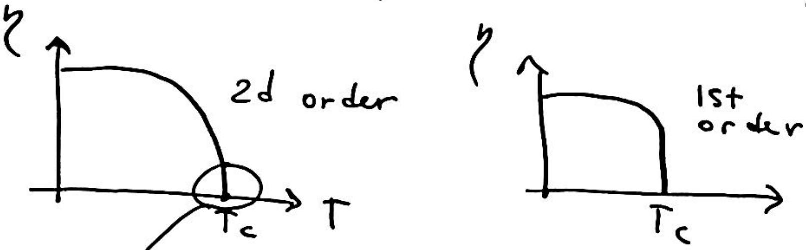
<u>Statement:</u> At finite T, the state of a system is a <u>minimum</u> of of a 'thermodynamic potential Helmholz or Gibs.

e.g. $F = E - T \cdot S$ $\rightarrow$ $T \to 0$ ~~≡~~ $S \to$ ~~0~~ is unimportant for to order
~~⇒~~ T goes up $S \uparrow$

with increasing T, entropy favors <u>disorder</u>.

Order parameter (OP)

OP depends ~~on~~ ~~on~~ T, P, V, external things like strain, mag. field electric fields etc.



$T_c \equiv$ critical temperature

around T_c η is small, lets try tallor expansion

Gibs free energy: $\Phi(P, T, \eta) = \Phi_0 + \alpha\eta + A\eta^2 + C\eta^3 + \ldots$

here α, A, C, B are functions of P, T
and hence we can determine them from the condition that

$\min \Phi$ for $T > T_c$ $\eta = 0$
for $T < T_c$ $\eta \neq 0$

$$\Phi = \Phi_0 + \alpha(P,T)\eta + A(P,T)\eta^2 + C(P,T)\eta^3 + B_{(P,T)}\eta^4 + \ldots$$

$$\frac{\partial \Phi}{\partial P} \text{ or } \frac{\partial \Phi}{\partial T} = 0$$

What can we say about the expansion coeff.s?
since Φ has to have a minimum the term
$\alpha\eta$ $\rightarrow$ requires $\alpha = 0$ otherwise
as a function of η Φ would always grow,

Next, for $A\eta^2$: we know $\eta = 0$ for $T > T_c$
$\eta \neq 0$ for $T < T_c$

① $A\eta^2$ = is present in the expansion.

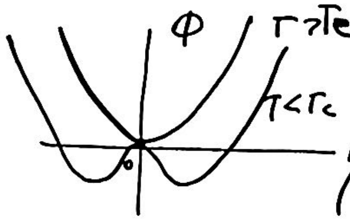
② $A(P,T)\eta^2 \Rightarrow$ $\eta = 0$ for $T > T_c$

$$A(P,T) = \begin{cases} > 0 & T > T_c \\ < 0 & T < T_c \end{cases}$$

the simplest term like this

$$A(P,T) = a(T - T_c)$$

we also assume for the moment that $C = 0$
and $B > 0$



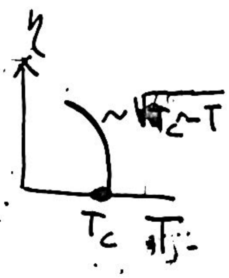
$$\Phi = \phi_0 + A(P,T)\eta^2 + \underset{>0}{B}(P,T)\eta^4 =$$

$$\boxed{\Phi = \Phi_0 + a(T - T_c)\eta^2 + B\eta^4}$$

To find out how the or. parameter depends on T

$$\frac{d\Phi}{d\eta} = 0 \rightarrow \cancel{2}a(T - T_c)\cancel{\eta} + \cancel{4}_{2}B\eta^{\cancel{3}\,2} = 0 \Rightarrow \boxed{\eta^2 = -\frac{\cancel{2}a(T - T_c)}{2B}}$$

Here all the coeff. can be a function of P
or ~~external~~ parameters: $a = a(P)$
$B = B(P)$
$T_c = T_c(P)$



For cond. mat. $T_c = T_c(P)$ is the most important.

and $a = B$ = const in P

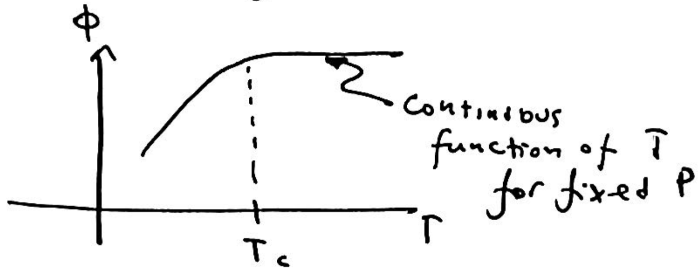
At the equilibrium the free energy:

$$\Phi_{min}\Big|_{T<T_c} = \Phi_0 + A\eta^2 + B\eta^4 = \Phi_0 + A\cdot\frac{a(T - T_c)}{2B} + \frac{\cancel{B}\,a^2(T - T_c)^2}{4B^{\cancel{2}}}$$

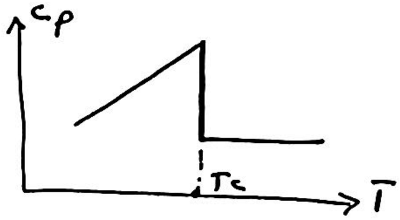
$$= \Phi_0 + (-)\frac{a^2(T - T_c)^2}{2B} + \cancel{B}\,\frac{a^2(T - c)^2}{4B} = \Phi_0 - \frac{a^2}{4B}(T - T_c)^2$$

and $\Phi = \Phi_0$ for $T > T_c$

$$\Phi_{min} = \begin{cases} \Phi_0 - \frac{a^2}{4B}(T-T_c)^2 & T<T_c \\ \Phi_0 & T>T_c \end{cases}$$



but $\frac{\partial \Phi}{\partial T}$ has a kink at T_c!



★ Homework: 1) Derive c_p behavior at the 2nd order phase transition temperature.

Given $c_p = T\left(\frac{\partial S}{\partial T}\right)_p$; $S = -\left(\frac{\partial \Phi}{\partial T}\right)_p$

$\eta^2 = \frac{a}{2B}(T_c - T)$, $\Phi_{min} = \Phi_0 - \frac{a^2}{4B}(T_c - T)^2$

~~Comment on~~

Assume $S = 1/2$, calculate total entropy in the ordered phase:

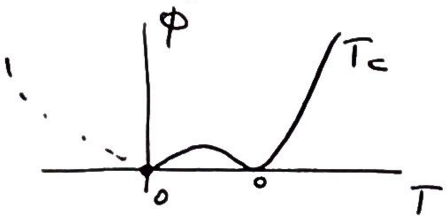
2). e.g. $S_{ord} = \int^{T_c} S\,dT$

(SEE SEPARATE PDF)

General Rules:

① if you include the cubic term $\sim \eta^3$ you always end up with 1^{st} order transitions.

② in this case T_c is not the singular minimum of Φ but the point where



where two minima have equal energy.

③ T^* and T^{**} are the singular points

④ Thermodynamic quantities e.g. C_p or C_v will not diverge at T_c but will diverge at T^* and T^{**}

limit of overcooling

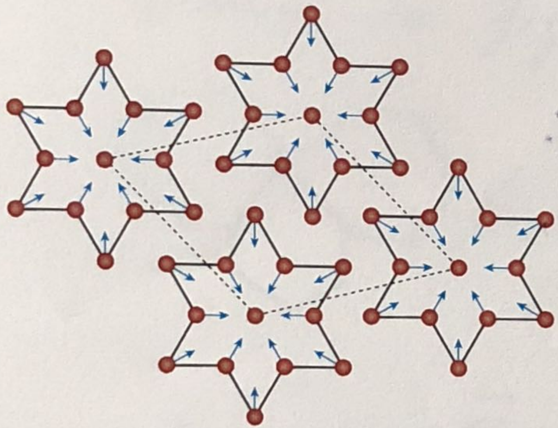
limit of ~~overheating~~ overheating

A Cool example of 1st order ph. transition.

SEE NEXT PAGE

L2

TaS_2



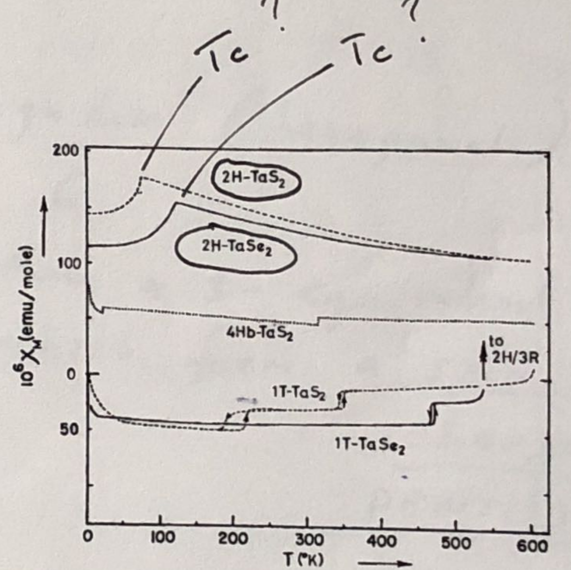


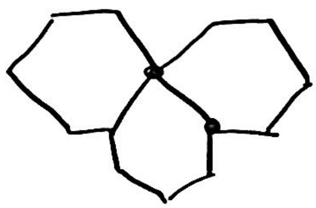
Fig. 2. The molar magnetic susceptibility (χ_M) versus temperature (T) for TaS_2 and T_c with different lattice structures. The background diamagnetic term has not been subtracted. The data are taken from ref. 3.

So how do we describe the order parameter in this system?

① we include every possible term allowed by symmetry

i.e. on triangular (hexagonal) lattice shown in fig. 6

there a 3-equivalent vectors which form a spin or SDW / charge CDW DENSITY WAVE



$$\eta_n = \eta\, e^{i \bar{Q}_n \cdot \bar{r}}$$

$n = 1, 2, 3$

so that $\bar{Q}_1 + \bar{Q}_2 + \bar{Q}_3 = 0$

Then we can form a new invariant in the Gibs free energy, i.e.

$$C\,\eta_1 \cdot \eta_2 \cdot \eta_3 = C\,\eta^3 e^{i\underbrace{(\bar{Q}_1 + \bar{Q}_2 + \bar{Q}_3)}_{=0} \cdot \bar{r}} =$$

$$= C\,\eta^3$$ ← this term will cause the 1st order phase transition found in TaS_2 or TaSe as shown in fig. on page 6 L2

Q: what do you think about crystallization or melting? is it a 1st or 2d order phase transition?

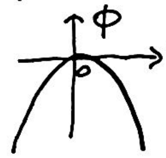
Another possibility of getting 1^{st} order.

Assume that there are no odd terms η^3, η^5 etc

$$\Phi = A(P,T)\eta^2 + B(P,T)\eta^4 =$$
$$= a(T - T_c)\eta^2 + \underset{>0}{B}(P,T)\eta^4$$

Assume that at some value of P the coeff. B is <0

so $T<T_c$ both η^2 and $\eta^4 < 0$

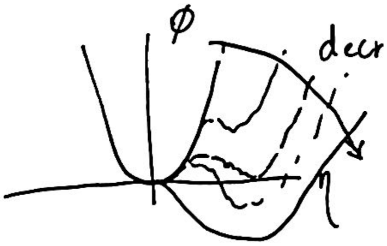


, so this system is unstable.

We need to stabilize the system, we add up

$$\Phi = \dots + \underset{>0}{D}\eta^6$$

For $D>0$ and $T>T_c$ where $A>0$ and very small, $B<0$



near T_c $a(T-T_c)$ is small

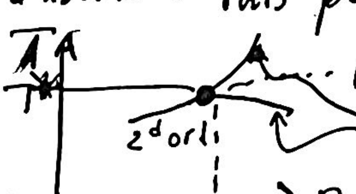
$-B\eta^4$ is large enough

so we may end up with 2 minima

or the 1^{st} order phase transition.

Note however if $B(P,T)>0$ then we will have a "regular" 2^d order phase tr.

Thus if at some pressure P^* $B(P^*,T)$ changes sing we will have 2^d to 1^{st} order transition. This point is called the tricritical point



lines of overheating and overcooling.

L2 (two)

Interaction with other degrees of freedom

Suppose we want to study how strain ε or pressure P affects a ~~the~~ magnetic phase transition.

bulk modulus
$b = \Delta P / \Delta V/V$
$b_{steel} = 160 \times 10^9 \ \frac{N}{m^2}$
$b_{H_2O} = 2.2 \times 10^9 \ \frac{N}{m^2}$
$\frac{1}{b}$ = compressibility

$$\Phi = a(T-T_c)\eta^2 + B\eta^4 + \underbrace{\frac{bu^2}{2}}_{\text{elastic energy}} + \lambda\eta^2 \cdot u$$

elastic energy

coupling between distortion u and magnetic order parameter

Findin min:

$$\frac{d\Phi}{du} = \cancelto{0}{a(T-T_c)\eta^2} + \cancelto{0}{B\eta^4} + bu + \lambda\eta^2 = 0$$

$$\Rightarrow \quad u_{min} = -\frac{\lambda\eta^2}{b}$$

$$\Phi(u_{min}) = a(T-T_c)\eta^2 + B\eta^4 + \frac{\cancel{b}}{2}\frac{\lambda^2\eta^4}{\cancel{b^2}} + \lambda\eta^2\left(-\frac{\lambda\eta^2}{b}\right) = a(T-T_c)\eta^2 + \left(B - \frac{\lambda^2}{2B}\right)\eta^4$$

now if coupling to the lattice, e.g. magnetoelastic coupling is strong, or lattice compressibility is large, i.e. bulk modulus b is small $\Rightarrow$

$$\underbrace{\left(B - \frac{\lambda^2}{2b}\right)}_{<0}\eta^4 \Rightarrow$$

2^d order turns into 1^{st} order phase transition

NB. thing which are difficult to compress has a large bulk modulus, b but small compressibility $1/b$

For those who care!

What if the order parameter is complex?

Superconductivity for pedestrians.

In transitioning to superconductivity the brooken symmetry is the gage symmetry.

① in the superconducting state the macroscopic wafe function is the complex order parameter, e.g.

$$\eta = \eta_0 e^{i\theta} \quad , \eta_0 \text{ and } \theta \text{ are real}$$

Now we can write down Φ!

$$\Phi = \Phi_0 + \underset{a(T-T_c)}{A}|\eta|^2 + \underset{>0}{B}|\eta|^4 + \dots =$$

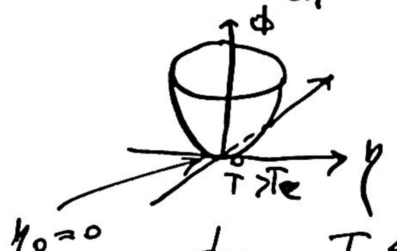
$$= \Phi_0 + a(T-T_c)\eta_0^2 + B\eta_0^4 \Rightarrow$$

$$\frac{d\Phi}{d\eta_0} = a(T-T_c)\cdot 2\eta_0 + 4B\eta^3 = 0 \Rightarrow \left(a(T-T_c) + 2B\eta_0^2\right)\cdot \eta_0 = 0$$

or $\eta_0 = 0$ for $T > T_c$

$$\eta_0 = \left[\frac{a(T-T_c)}{2B}\right]^{1/2}, T < T_c$$

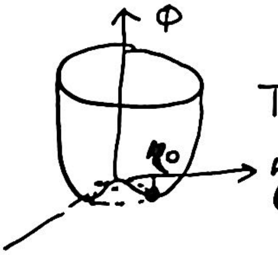
For the normal state of a superconductor when $T > T_c$ $\Phi = \Phi_0$ and gage symmetry is preserved



or $\begin{cases} \eta_0 = 0 \\ \text{and thus } \theta \text{ can take any value} \end{cases}$

for $T < T_0$ $\eta_0 = \left(\frac{a(T-T_c)}{2B}\right)^{1/2}$ and θ is fixed $\Rightarrow$

$\Rightarrow$ the symmetry is broken!

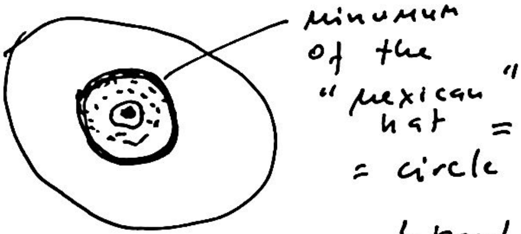


$\eta = \eta_0 e^{i\theta}$, however the system can moved inside the "mexican hat"

i.e. $\eta \to \eta' = e^{i\theta'}\cdot\eta$ the set of all rotations in the complex plane $U(1)$

So it seems that even in the superconducting phase we have gage symmetry in the ordered phase. i.e $\theta \to \theta'$ but with the same energy.
All those orderer states are degenerate in energy.

Topview:



every point on the circk is a possible state.

When the symmetry is spontaneously broken the system choses a specific η^* and θ^*.

Physical picture:

Just below T_c only a small number of electons codense into Cooper pairs with specific θ. When $T = 0K$ max number of N condens.
so N and θ are conjugate: or:
$\Delta N \cdot \Delta\theta \sim 1$,